

Hadronization & Underlying Event

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Lecture Notes:

[P. Skands, arXiv:1207.2389](https://arxiv.org/abs/1207.2389)

From Partons to Pions

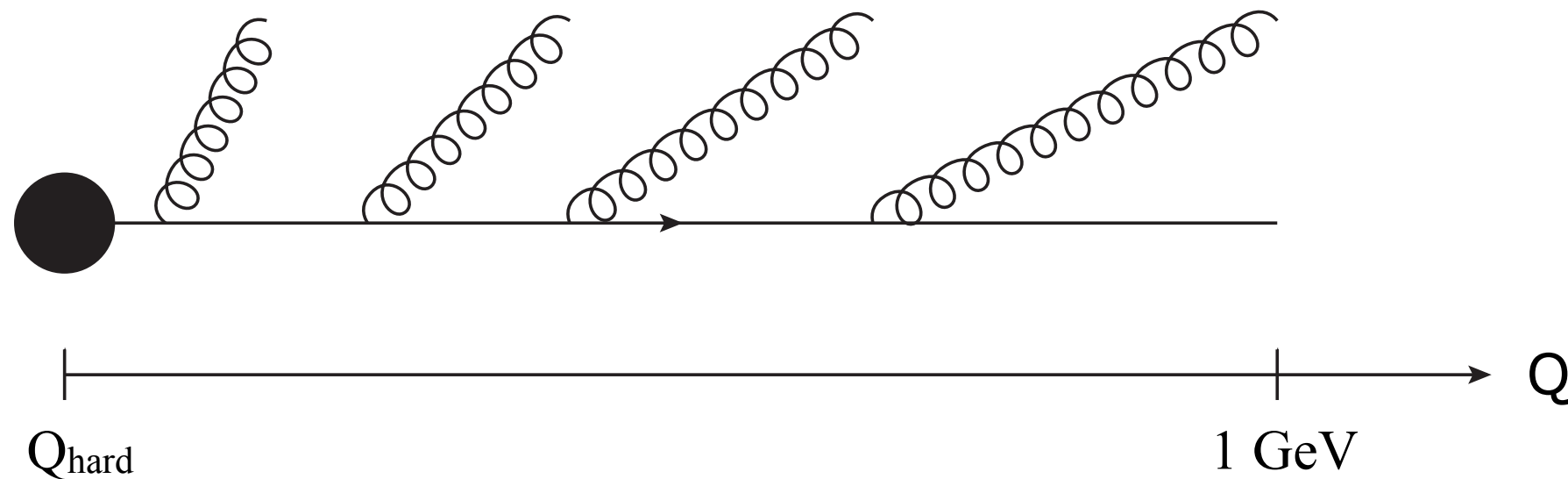
Here's a fast parton

Fast: It starts at a high factorization scale

$$Q = Q_F = Q_{\text{hard}}$$

It showers
(perturbative
bremsstrahlung)

It ends up
at a low effective
factorization scale
 $Q \sim m_\rho \sim 1 \text{ GeV}$



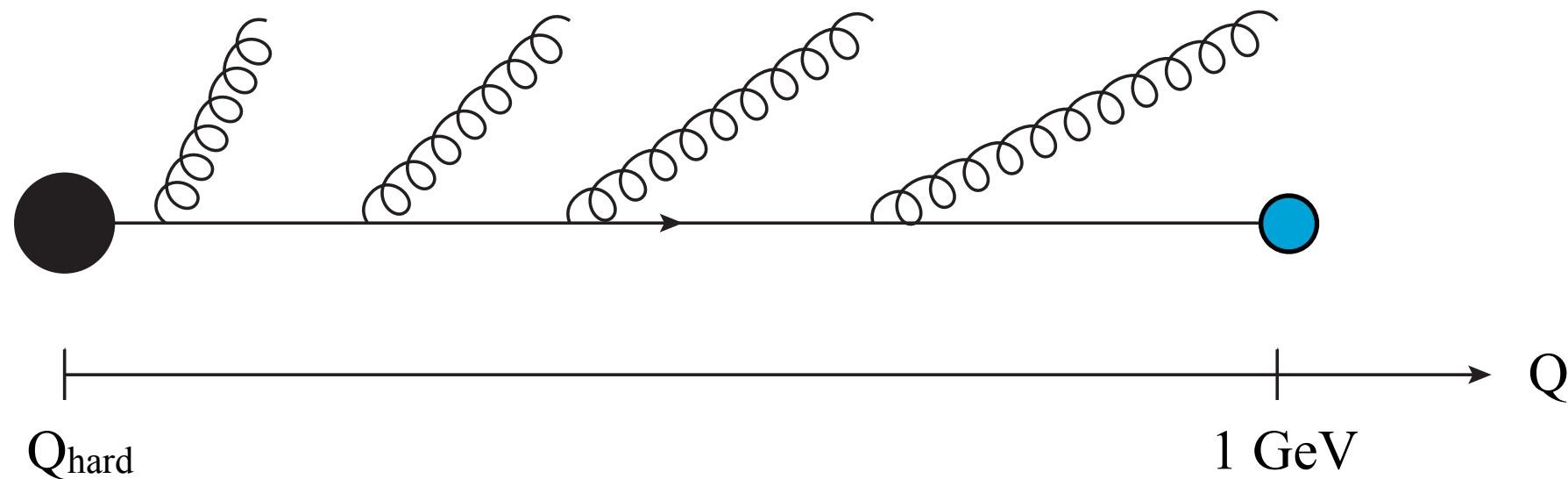
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How about I just call it a hadron?

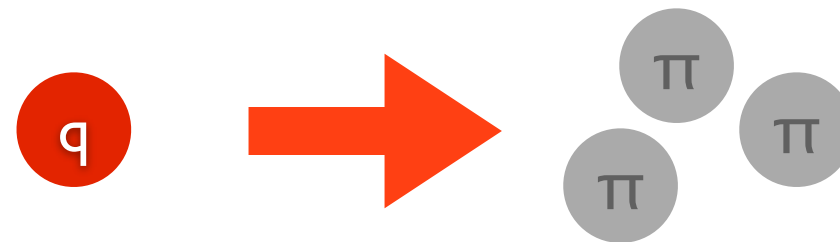
→ "Local Parton-Hadron Duality"

Parton \rightarrow Hadrons?

Early models: “Independent Fragmentation”

Local Parton Hadron Duality (LPHD) can give useful results for **inclusive** quantities in collinear fragmentation

Motivates a simple model:



But ...

The point of confinement is that partons are coloured

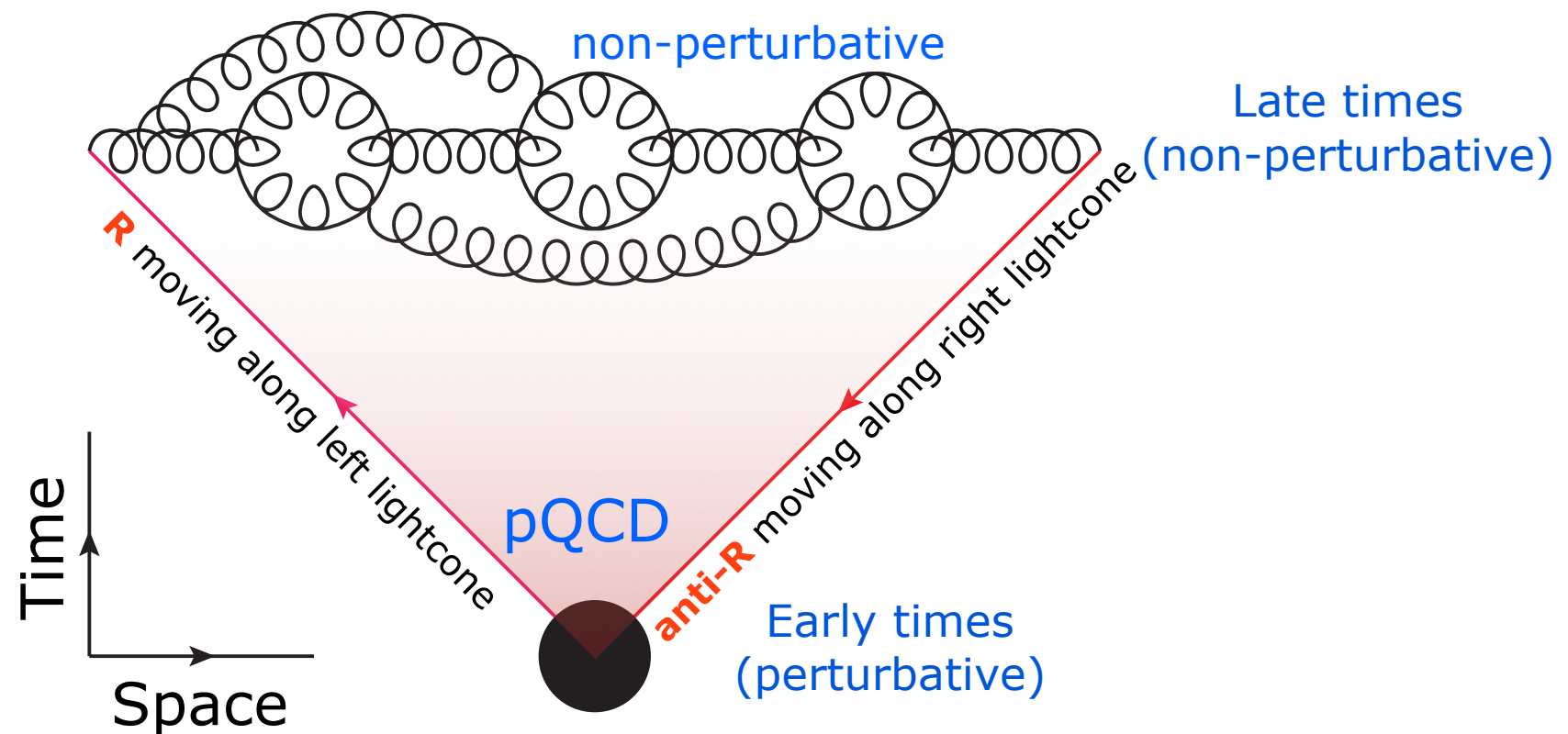
Hadronization = the process of colour neutralization

- \rightarrow Unphysical to think about independent fragmentation of a single parton into hadrons
- \rightarrow Too naive to see LPHD (inclusive) as a justification for Independent Fragmentation (exclusive)
- \rightarrow More physics needed

Colour Neutralization

A physical hadronization model

Should involve at least TWO partons, with opposite color charges (e.g., **R** and **anti-R**)



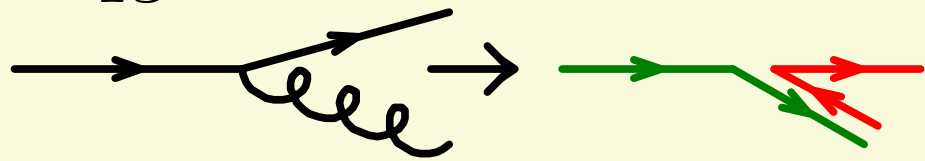
Strong “confining” field emerges between the two charges when their separation $> \sim 1\text{fm}$

Color Flow

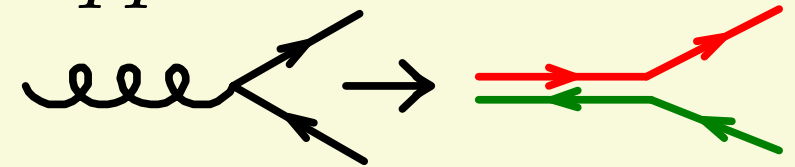
Between which partons do confining potentials arise?

Set of simple rules for color flow, based on large- N_c limit

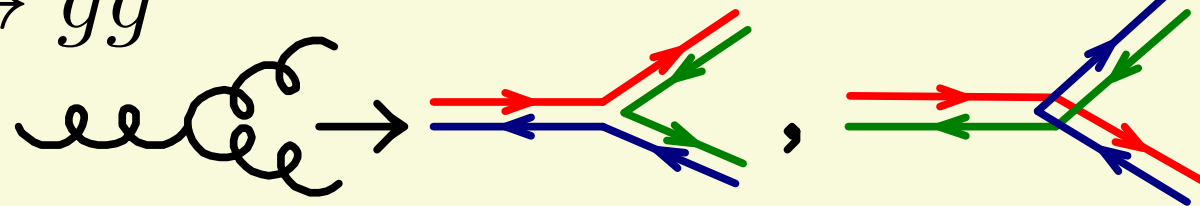
$q \rightarrow qg$



$g \rightarrow q\bar{q}$



$g \rightarrow gg$

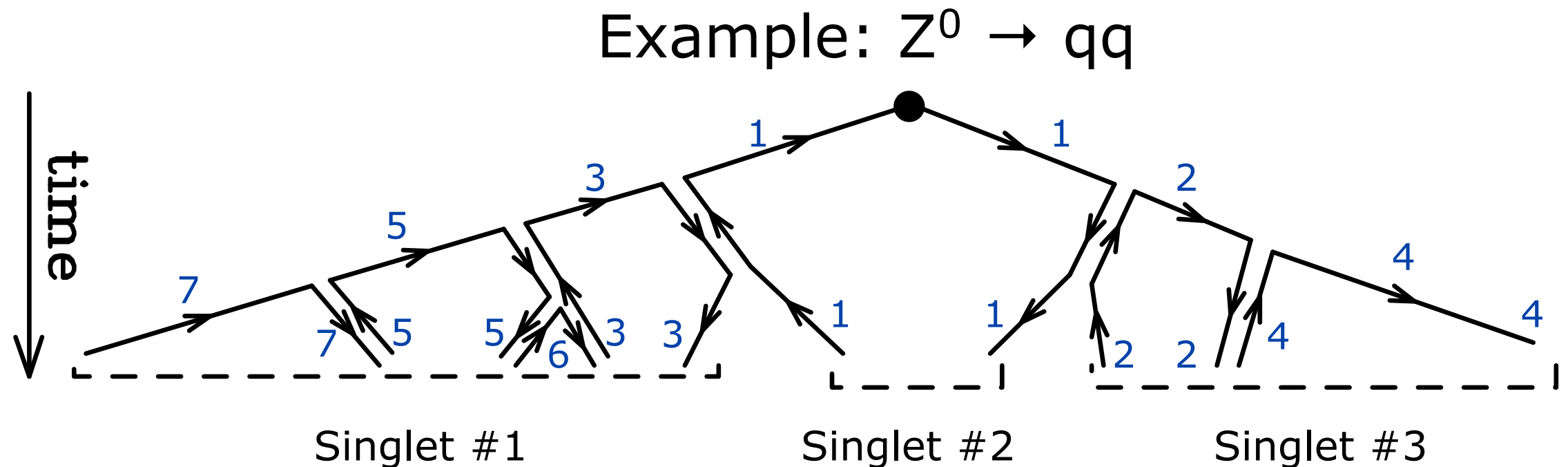


(Never Twice Same Color: true up to $O(1/N_c^2)$)

Illustrations from: P.Nason & P.S.,
PDG Review on MC Event Generators, 2012

Color Flow

For an entire Cascade



Coherence of pQCD cascades \rightarrow not much “overlap” between singlet subsystems
 \rightarrow Leading-colour approximation pretty good

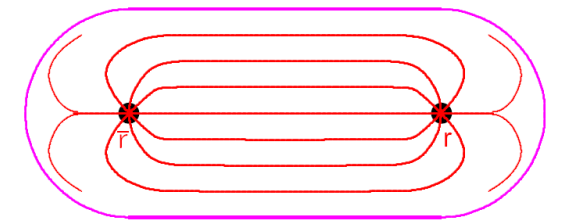
LEP measurements in WW confirm this (at least to order 10% $\sim 1/N_c^2$)

Note: (much) more color getting kicked around in hadron collisions \rightarrow more later

Confinement

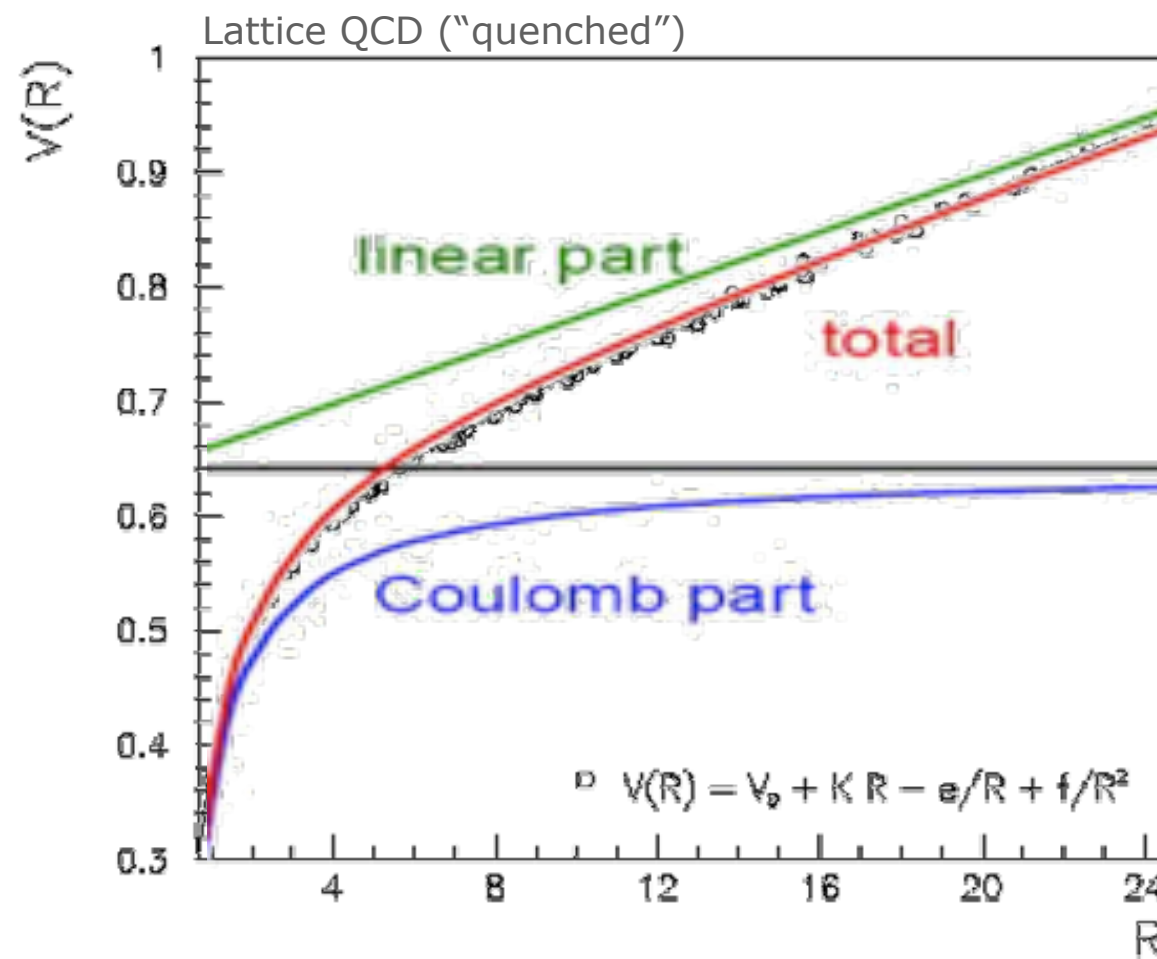
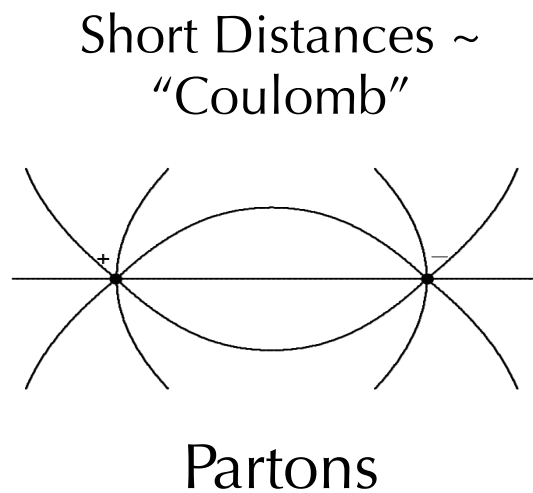
Potential between a quark and an antiquark as function of distance, R

Long Distances ~ Linear Potential



Quarks (and gluons) confined inside hadrons

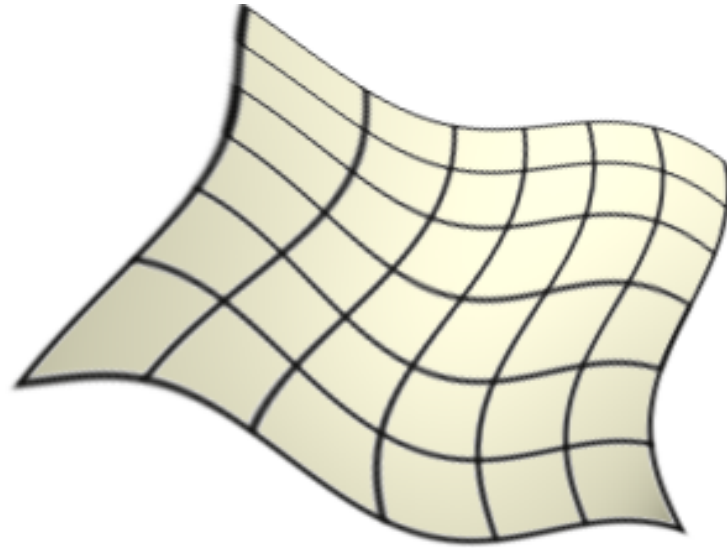
What physical system has a linear potential?



$$F(r) \approx \text{const} = \kappa \approx 1 \text{ GeV/fm} \iff V(r) \approx \kappa r$$

~ Force required to lift a 16-ton truck

From Partons to Strings



Motivates a model:

Let color field collapse into a (infinitely) narrow flux tube of uniform energy density $\kappa \sim 1 \text{ GeV} / \text{fm}$

→ Relativistic 1+1 dimensional worldsheet – string

Pedagogical Review: B. Andersson, *The Lund model*.
Camb. Monogr. Part. Phys. Nucl. Phys. Cosmol., 1997.

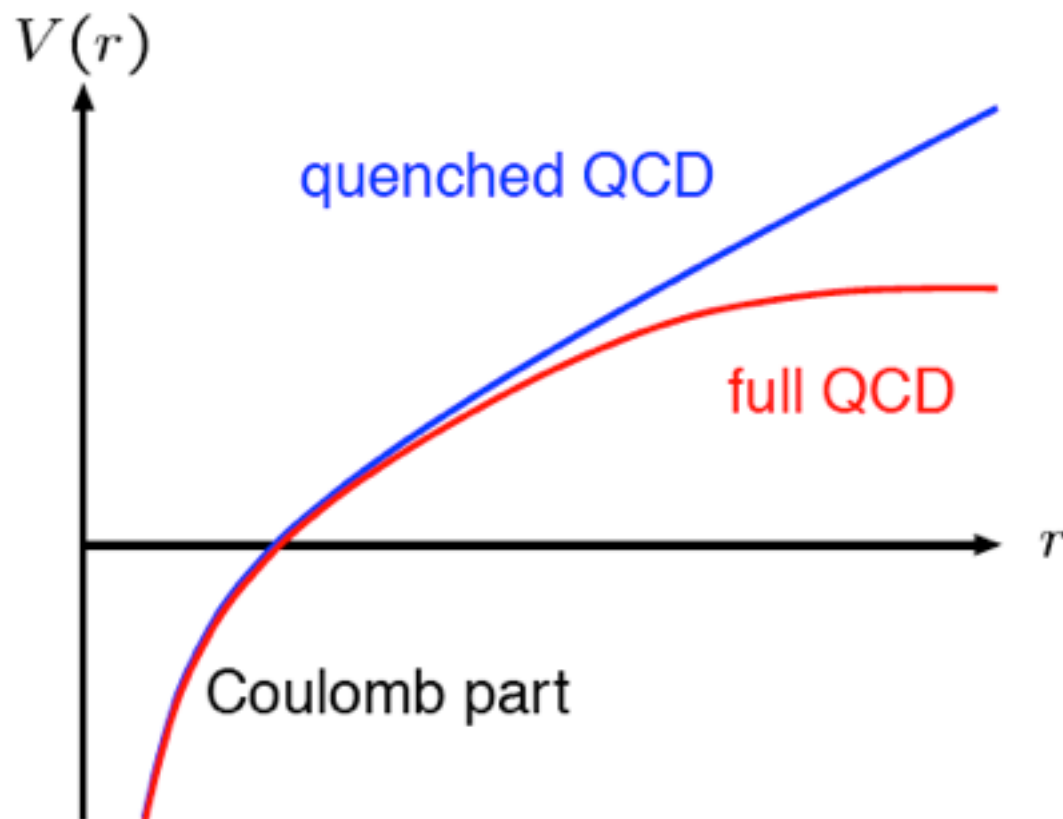
String Breaks



String Breaks

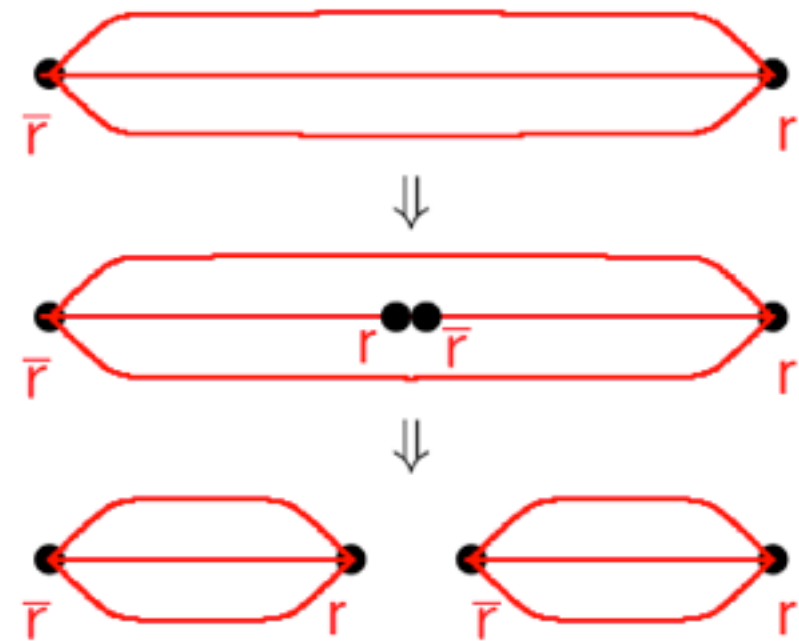
In “unquenched” QCD

$g \rightarrow qq \rightarrow$ The strings would break



- Gaussian p_T spectrum
- Heavier quarks suppressed. $\text{Prob}(q=d,u,s,c) \approx 1 : 1 : 0.2 : 10^{-11}$

String Breaks: via Quantum Tunneling



(simplified colour representation)

$$\mathcal{P} \propto \exp \left(\frac{-m_q^2 - p_{\perp}^2}{\kappa/\pi} \right)$$

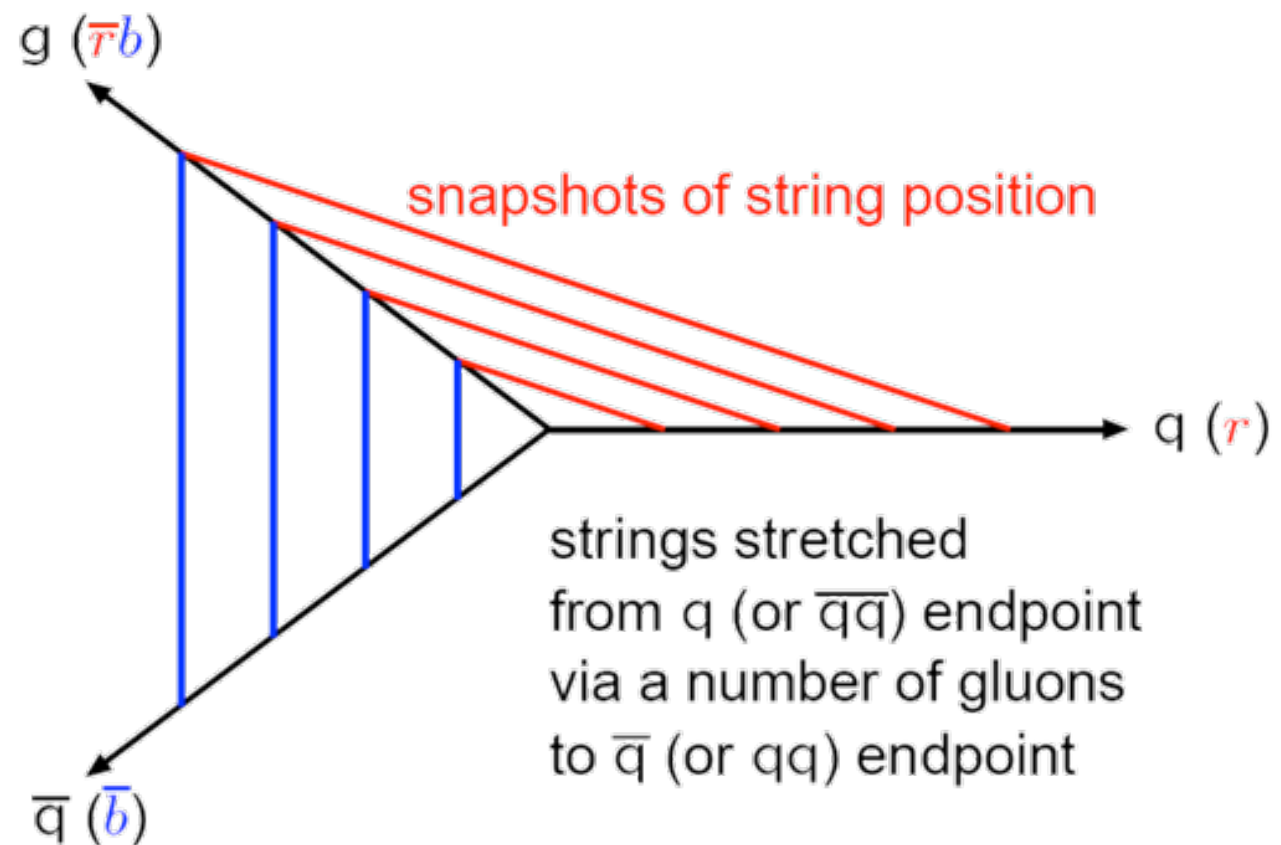
Illustrations by T. Sjöstrand

The (Lund) String Model

Map:

- **Quarks** → String Endpoints
- **Gluons** → Transverse Excitations (kinks)
- Physics then in terms of string worldsheet evolving in spacetime
- Probability of string break (by quantum tunneling) constant per unit area → **AREA LAW**

See also Yuri's 2nd lecture



Gluon = kink on string, carrying energy and momentum

→ **STRING EFFECT**

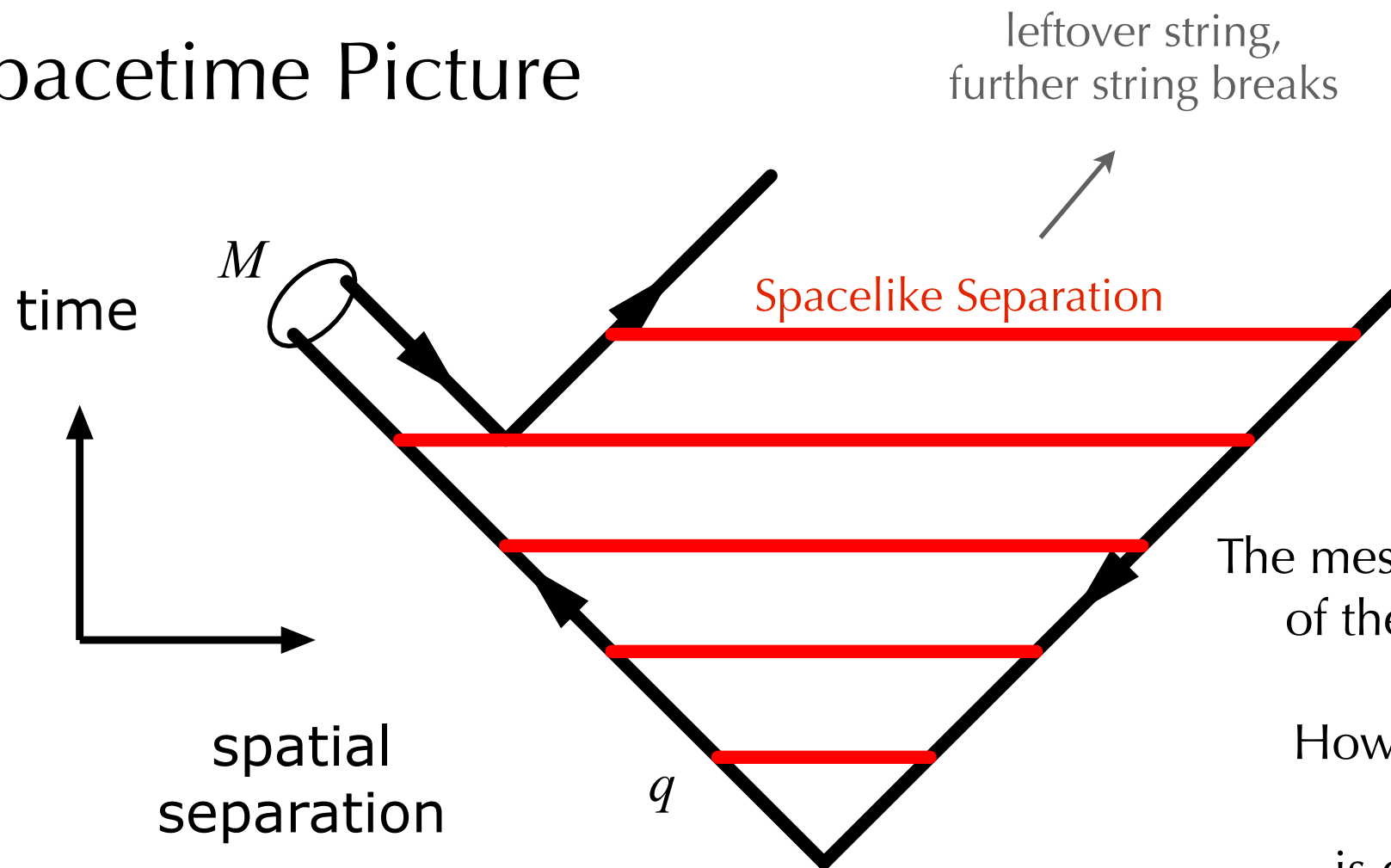
Physics now in terms of strings, with kinks, evolving in spacetime
Very simple space-time picture, few parameters at this point

Fragmentation Function

Having selected a hadron flavor

How much momentum does it take?

Spacetime Picture



The meson M takes a fraction z of the quark momentum,

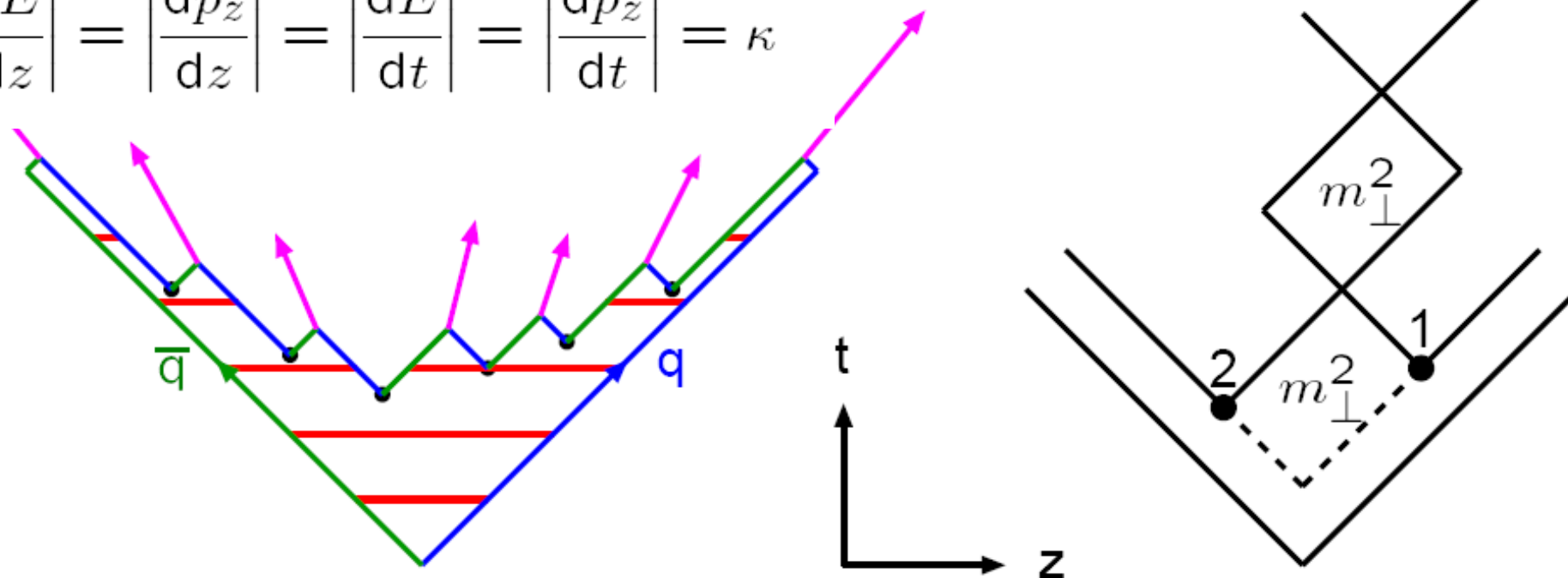
How big that fraction is, $z \in [0,1]$, is determined by the fragmentation function, $f(z, Q_0^2)$

Large System

Illustrations by T. Sjöstrand

Repeat for large system \rightarrow Lund Model

$$\left| \frac{dE}{dz} \right| = \left| \frac{dp_z}{dz} \right| = \left| \frac{dE}{dt} \right| = \left| \frac{dp_z}{dt} \right| = \kappa$$



String breaks are causally disconnected

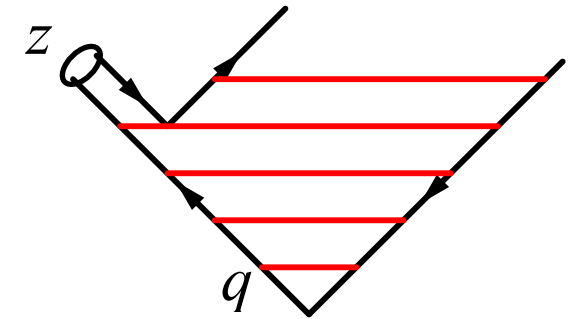
- \rightarrow can proceed in arbitrary order (left-right, right-left, in-out, ...)
- \rightarrow constrains possible form of fragmentation function
- \rightarrow Justifies iterative ansatz (useful for MC implementation)

Left-Right Symmetry

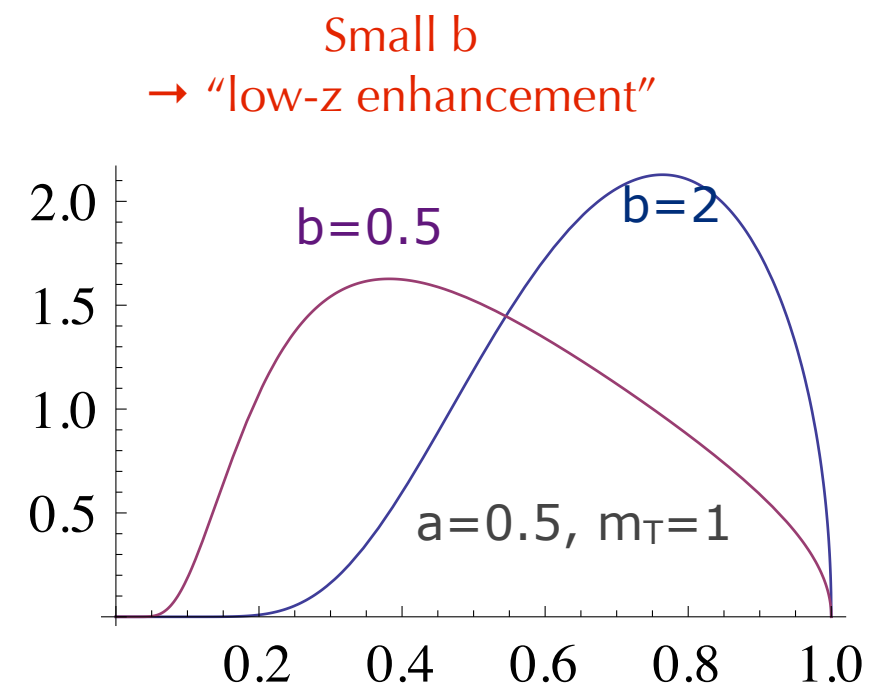
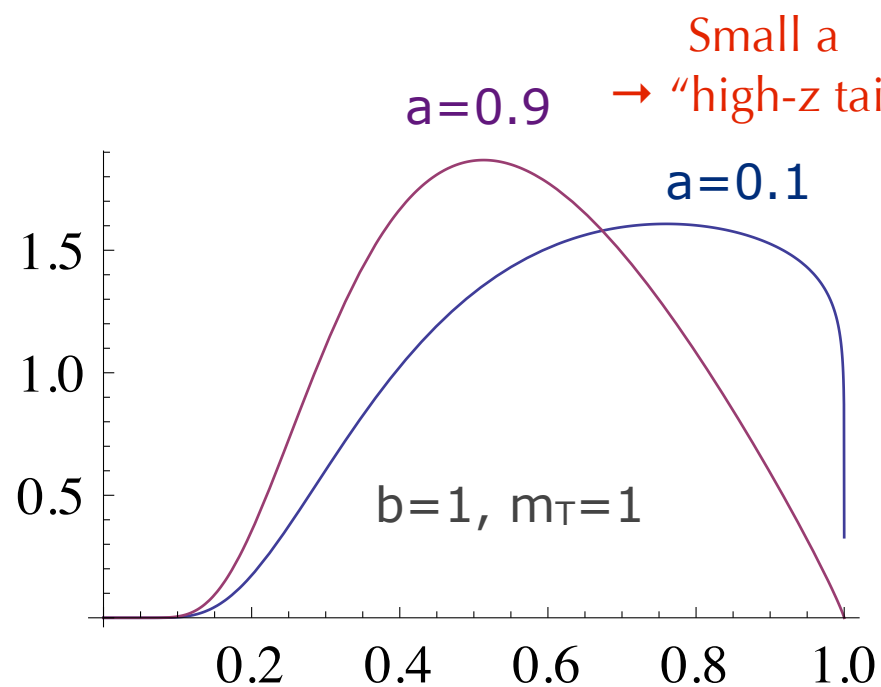
Causality → Left-Right Symmetry

→ Constrains form of fragmentation function!

→ Lund Symmetric Fragmentation Function



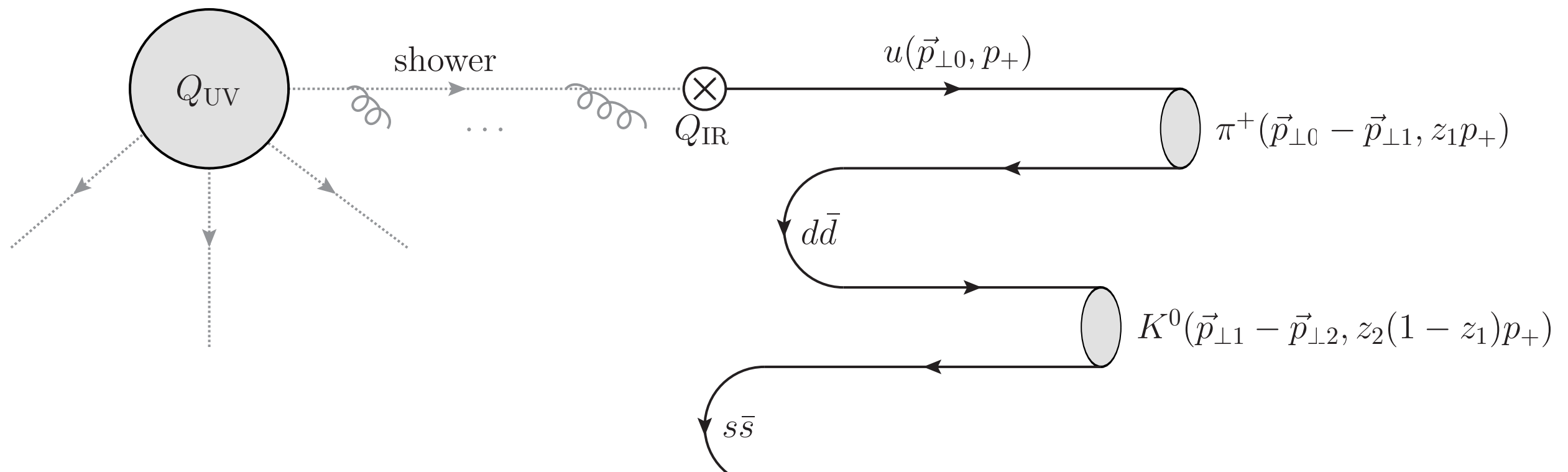
$$f(z) \propto \frac{1}{z} (1-z)^a \exp\left(-\frac{b(m_h^2 + p_{\perp h}^2)}{z}\right)$$



Note: In principle, a can be flavour-dependent. In practice, we only distinguish between baryons and mesons

Iterative String Breaks

Causality → May iterate from outside-in



The Length of Strings

In Space:

String tension ≈ 1 GeV/fm \rightarrow a 5-GeV quark can travel 5 fm before all its kinetic energy is transformed to potential energy in the string.

Then it must start moving the other way. String breaks will have happened behind it \rightarrow yo-yo model of mesons

In Rapidity :

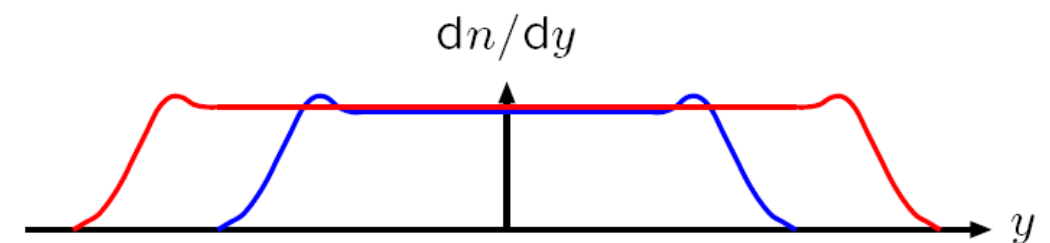
$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) = \frac{1}{2} \ln \left(\frac{(E + p_z)^2}{E^2 - p_z^2} \right)$$

For a pion with $z=1$ along string direction
(For beam remnants, use a proton mass):

$$y_{\max} \sim \ln \left(\frac{2E_q}{m_\pi} \right)$$

Note: Constant average hadron multiplicity per unit $y \rightarrow$ logarithmic growth of total multiplicity

Scaling in lightcone $p_\pm = E \pm p_z$ (for $q\bar{q}$ system along z axis) implies flat central rapidity plateau + some endpoint effects:



$$\langle n_{\text{ch}} \rangle \approx c_0 + c_1 \ln E_{\text{cm}}, \sim \text{Poissonian multiplicity distribution}$$

Alternative: The Cluster Model

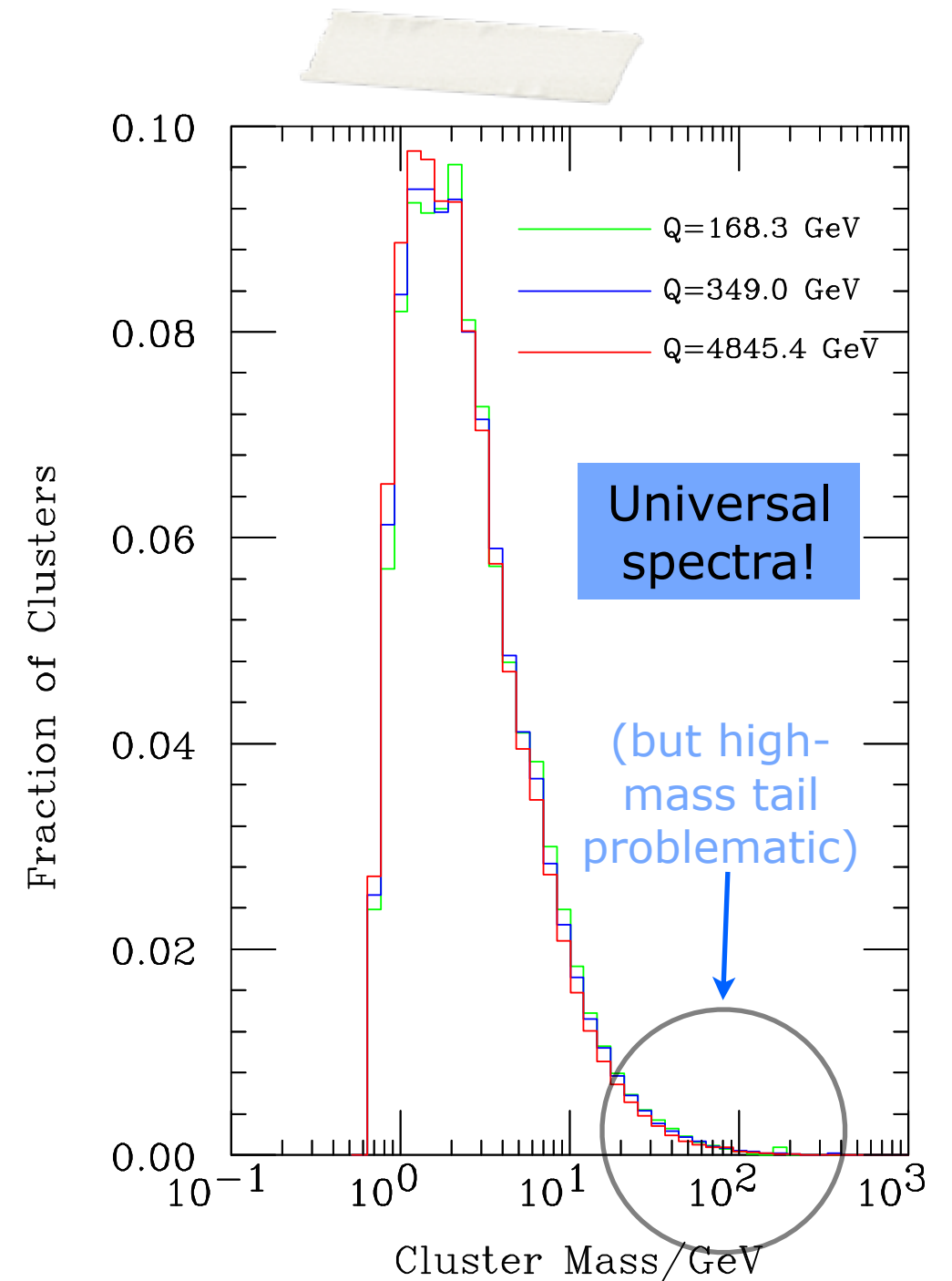
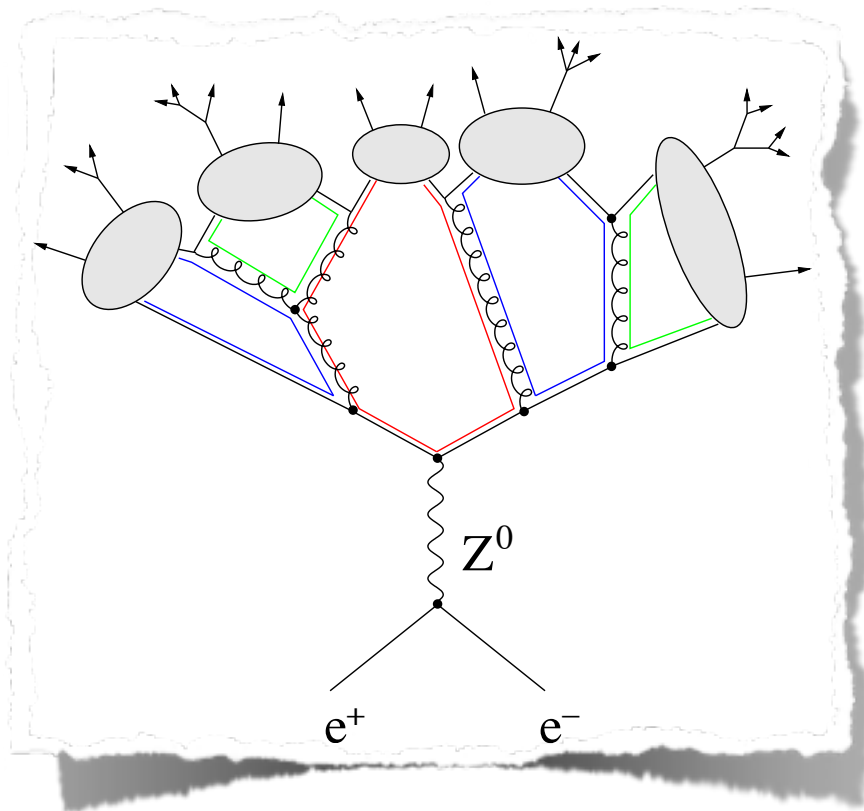
“Preconfinement”

+ Force $g \rightarrow qq$ splittings at Q_0

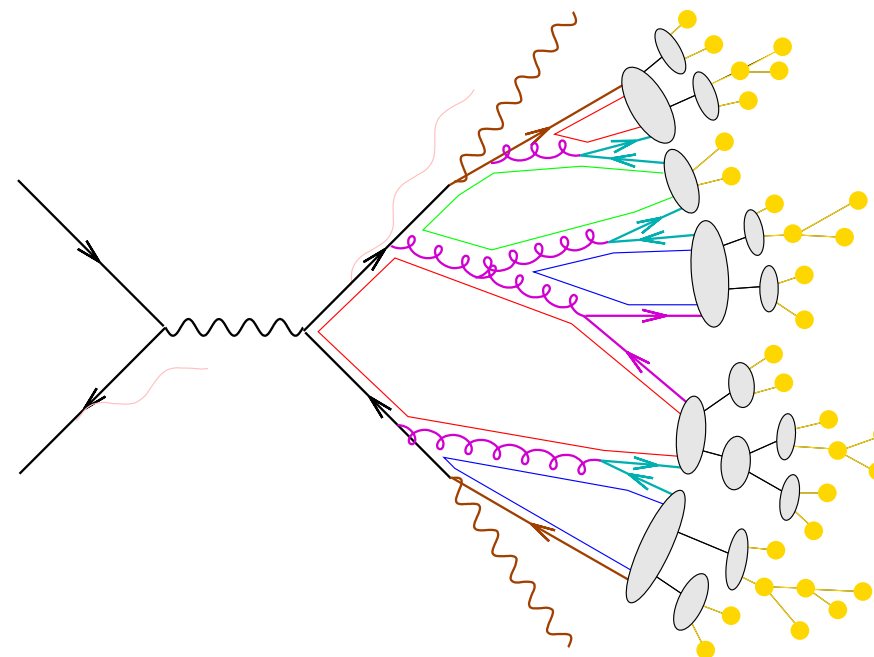
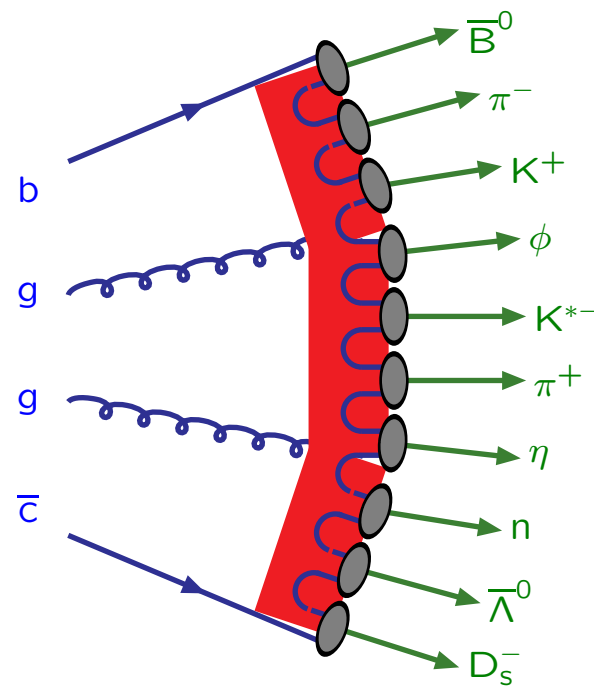
→ high-mass q - q bar “clusters”

Isotropic 2-body decays to hadrons

according to PS $\approx (2s_1+1)(2s_2+1)(p^*/m)$



Strings and Clusters



program	PYTHIA	HERWIG (&SHERPA)
model	string	cluster
energy-momentum picture	powerful predictive	simple unpredictive
parameters	few	many
flavour composition	messy unpredictive	simple in-between
parameters	many	few

Hadron Collisions

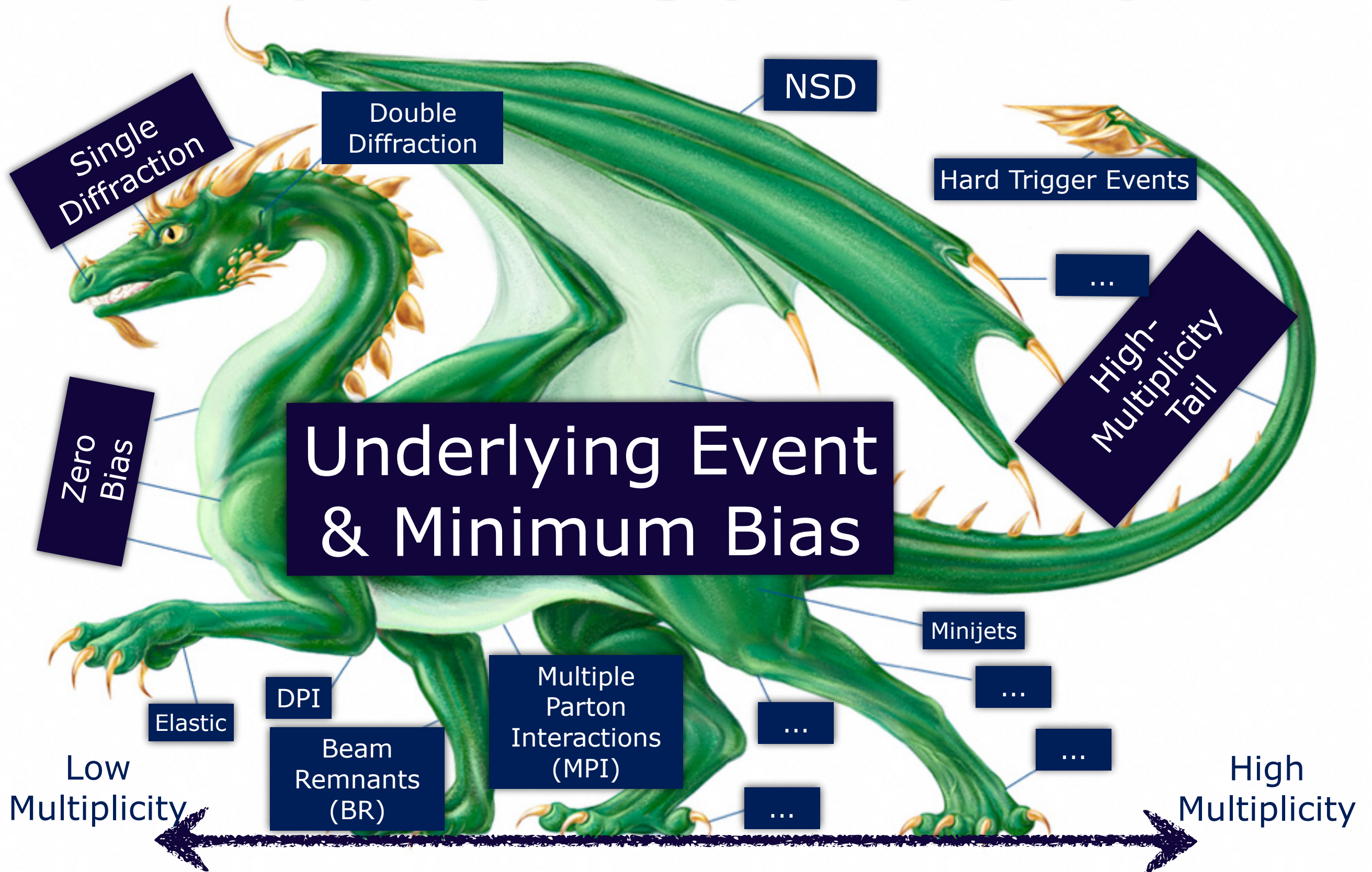
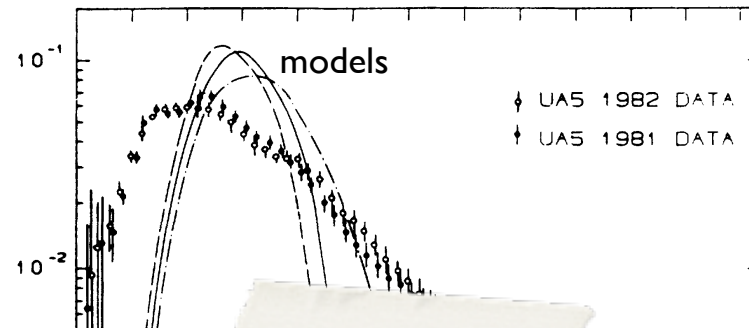


Image credits: E. Arenhaus & J. Walker

Hadron Collisions



Do not be scared of the failure of physical models
(typically points to more interesting physics)

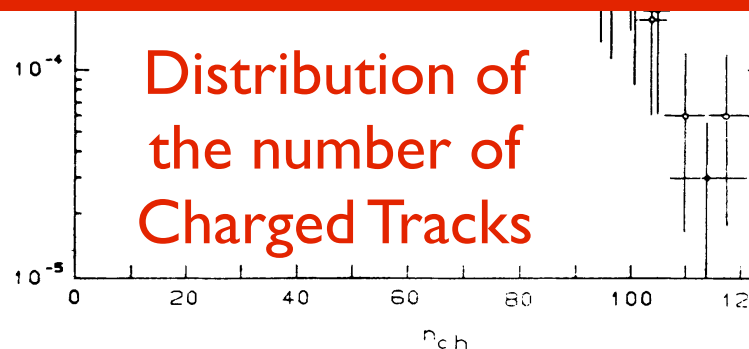


FIG. 3. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs simple models: dashed low p_T only, full including hard scatterings, dash-dotted also including initial- and final-state radiation.

Hadron Collisions

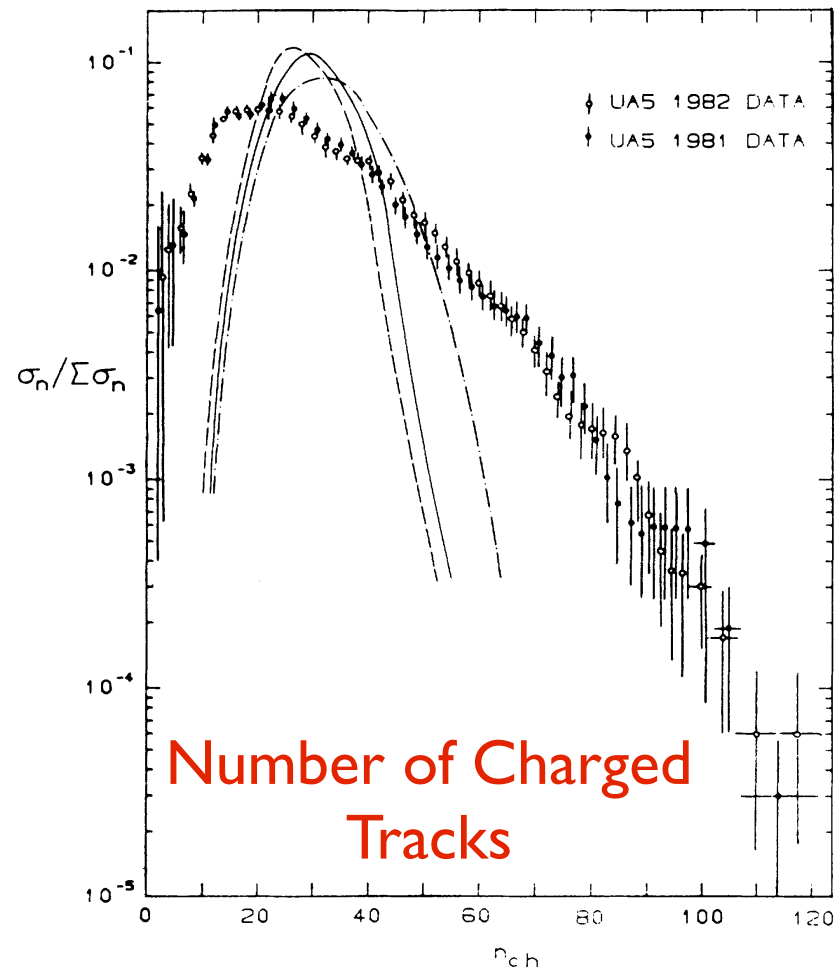


FIG. 3. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs simple models: dashed low p_T only, full including hard scatterings, dash-dotted also including initial- and final-state radiation.

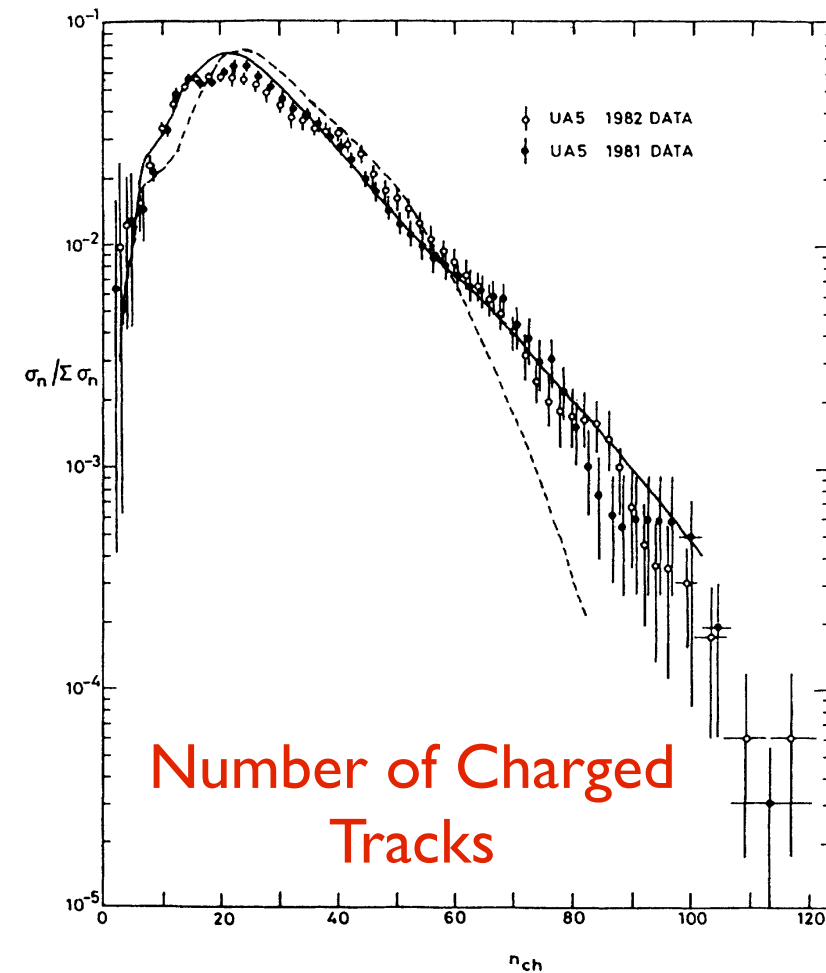


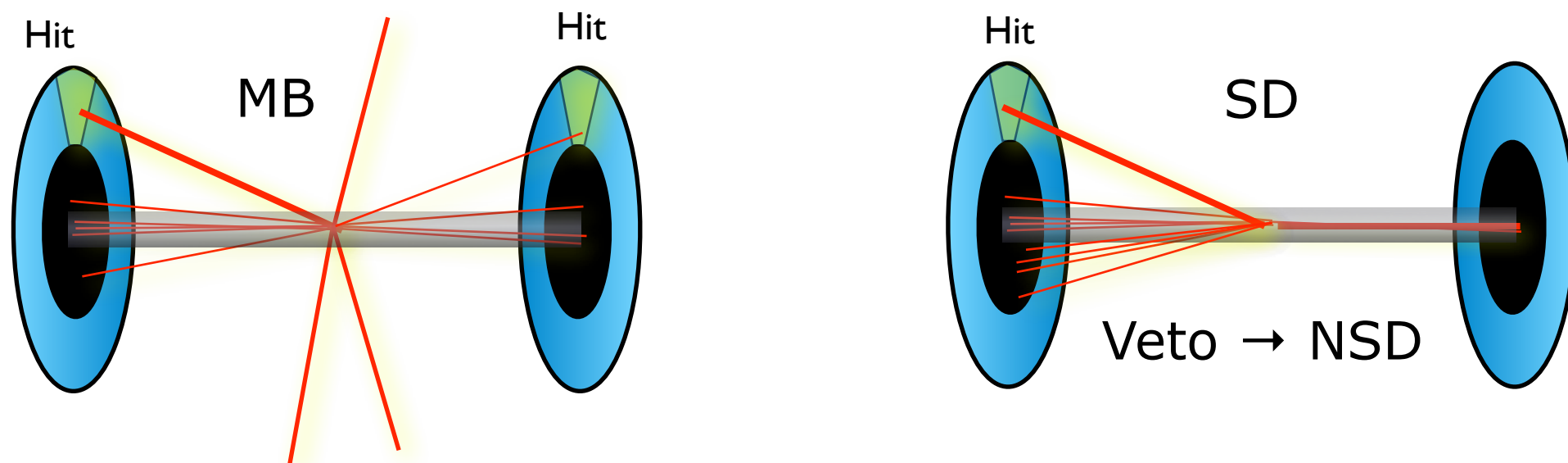
FIG. 12. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs multiple-interaction model with variable impact parameter: solid line, double-Gaussian matter distribution; dashed line, with fix impact parameter [i.e., $\bar{O}_0(b)$].

What is Pileup / Min-Bias?

We use Minimum-Bias (MB) data to test soft-QCD models

Pileup = “Zero-bias”

“Minimum-Bias” typically suppresses diffraction by requiring two-armed coincidence, and/or $\geq n$ particle(s) in central region



→ Pileup contains more diffraction than Min-Bias

Total diffractive cross section $\sim 1/3 \sigma_{\text{inel}}$

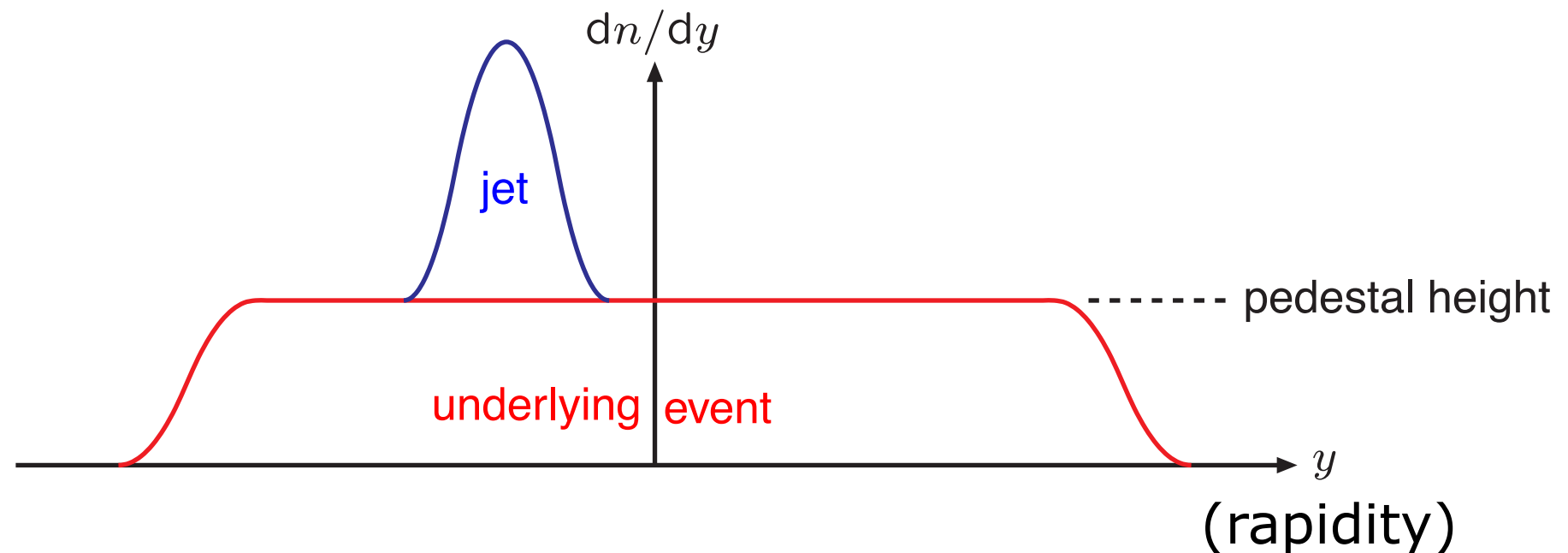
Most diffraction is low-mass → no contribution in central regions

High-mass tails could be relevant in FWD region

→ direct constraints on diffractive components (→ later)

What is Underlying Event ?

“Pedestal Effect”



Useful variable in hadron collisions: **Rapidity** (now along beam axis)

Designed to be additive
under Lorentz Boosts along
beam (z) direction

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right)$$

$$y \rightarrow -\infty \text{ for } p_z \rightarrow -E \quad y \rightarrow 0 \text{ for } p_z \rightarrow 0 \quad y \rightarrow \infty \text{ for } p_z \rightarrow E$$

Illustrations by T. Sjöstrand

Questions

Pileup

How much? In central & fwd acceptance?

Structure: averages + fluctuations, particle composition, lumpiness,
...

Scaling to 13 TeV and beyond

Underlying Event ~ “A handful of pileup” ?

Hadronizes with Main Event → “Color reconnections”

Additional “minijets” from multiple parton interactions

Hadronization

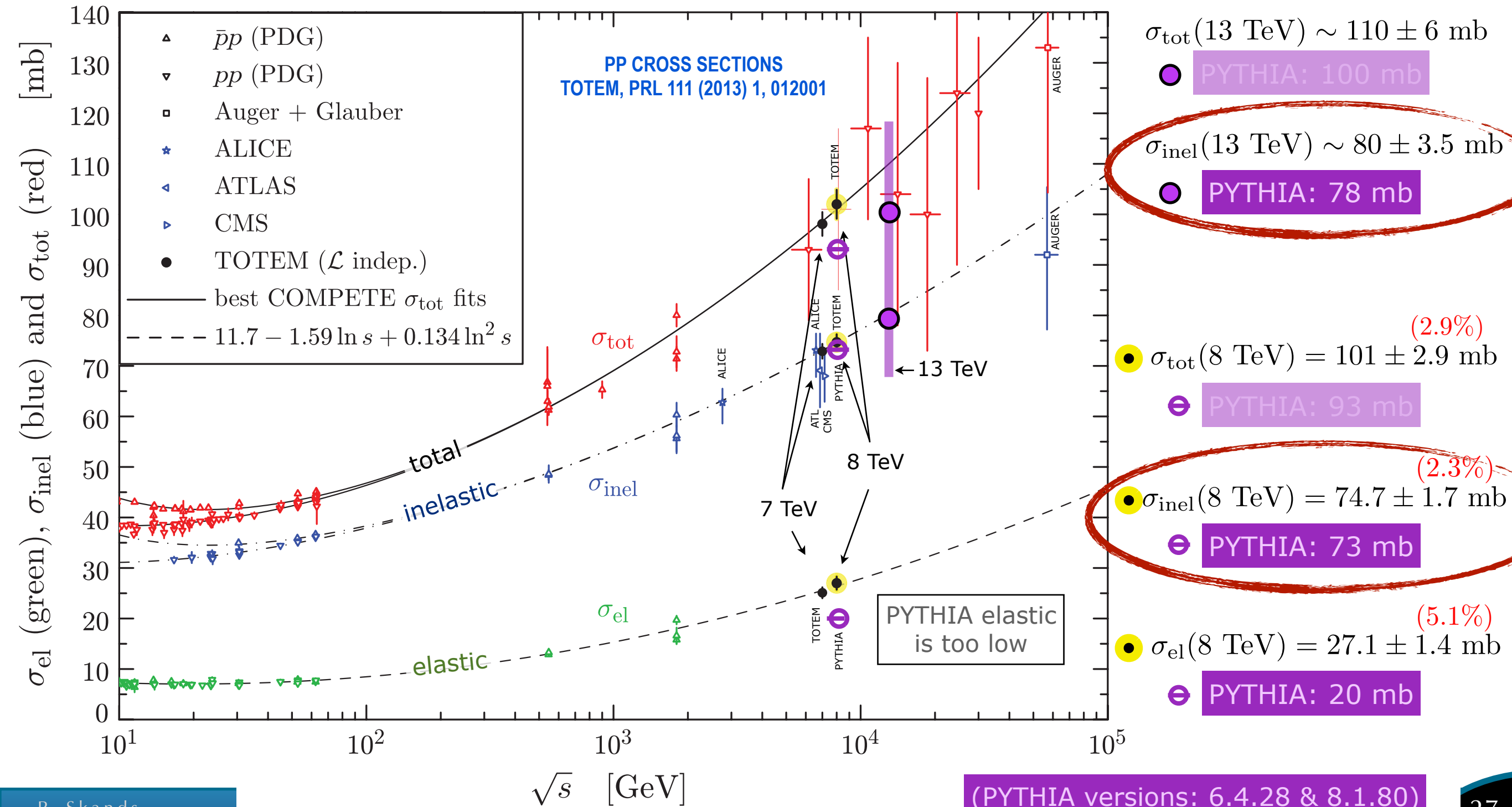
Models from the 80ies, mainly constrained in 90ies

Meanwhile, perturbative models have evolved

The Total Cross Section

Pileup rate $\propto \sigma_{\text{tot}}(s) = \sigma_{\text{el}}(s) + \sigma_{\text{inel}}(s) \propto s^{0.08}$ or $\ln^2(s)$?

Donnachie-Landshoff Froissart-Martin Bound



The Inelastic Cross Section

First try: decompose

$$\sigma_{\text{inel}} = \sigma_{\text{sd}} + \sigma_{\text{dd}} + \sigma_{\text{cd}} + \sigma_{\text{nd}}$$

+ Parametrizations of diffractive components: dM^2/M^2

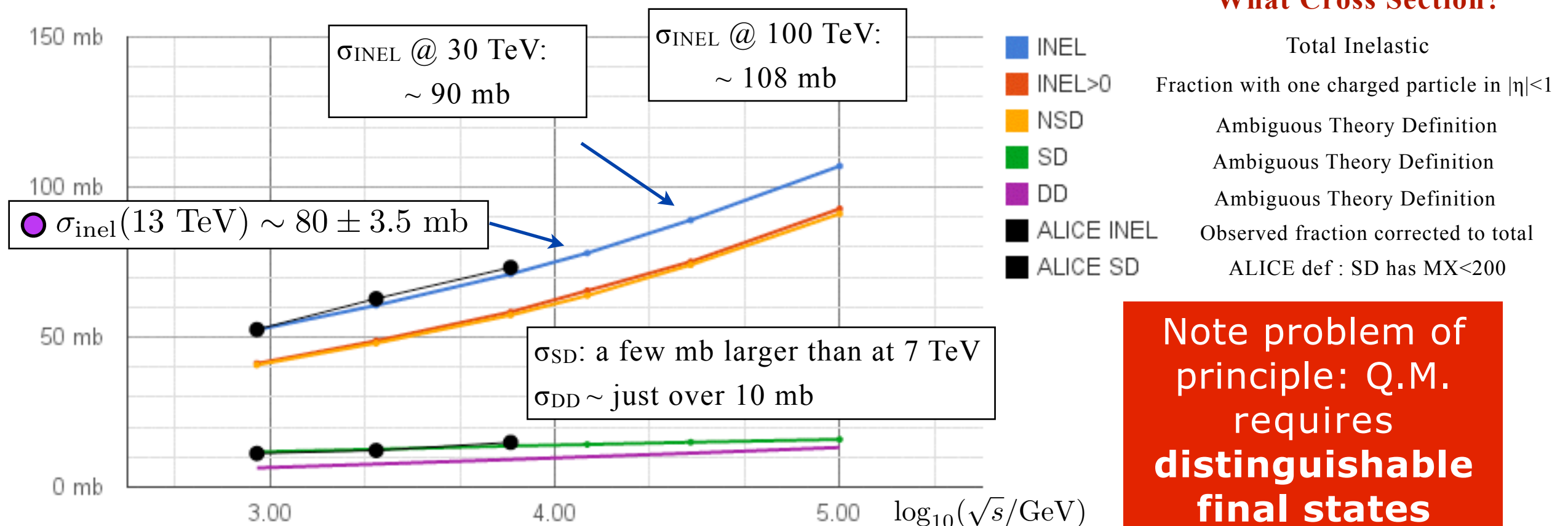
PYTHIA:

$$\frac{d\sigma_{\text{sd}(AX)}(s)}{dt dM^2} = \frac{g_{3\text{IP}}}{16\pi} \beta_{\text{AIP}}^2 \beta_{\text{BIP}} \frac{1}{M^2} \exp(B_{\text{sd}(AX)}t) F_{\text{sd}},$$

$$\frac{d\sigma_{\text{dd}}(s)}{dt dM_1^2 dM_2^2} = \frac{g_{3\text{IP}}^2}{16\pi} \beta_{\text{AIP}} \beta_{\text{BIP}} \frac{1}{M_1^2} \frac{1}{M_2^2} \exp(B_{\text{dd}}t) F_{\text{dd}}.$$

+ Integrate and solve for σ_{nd}

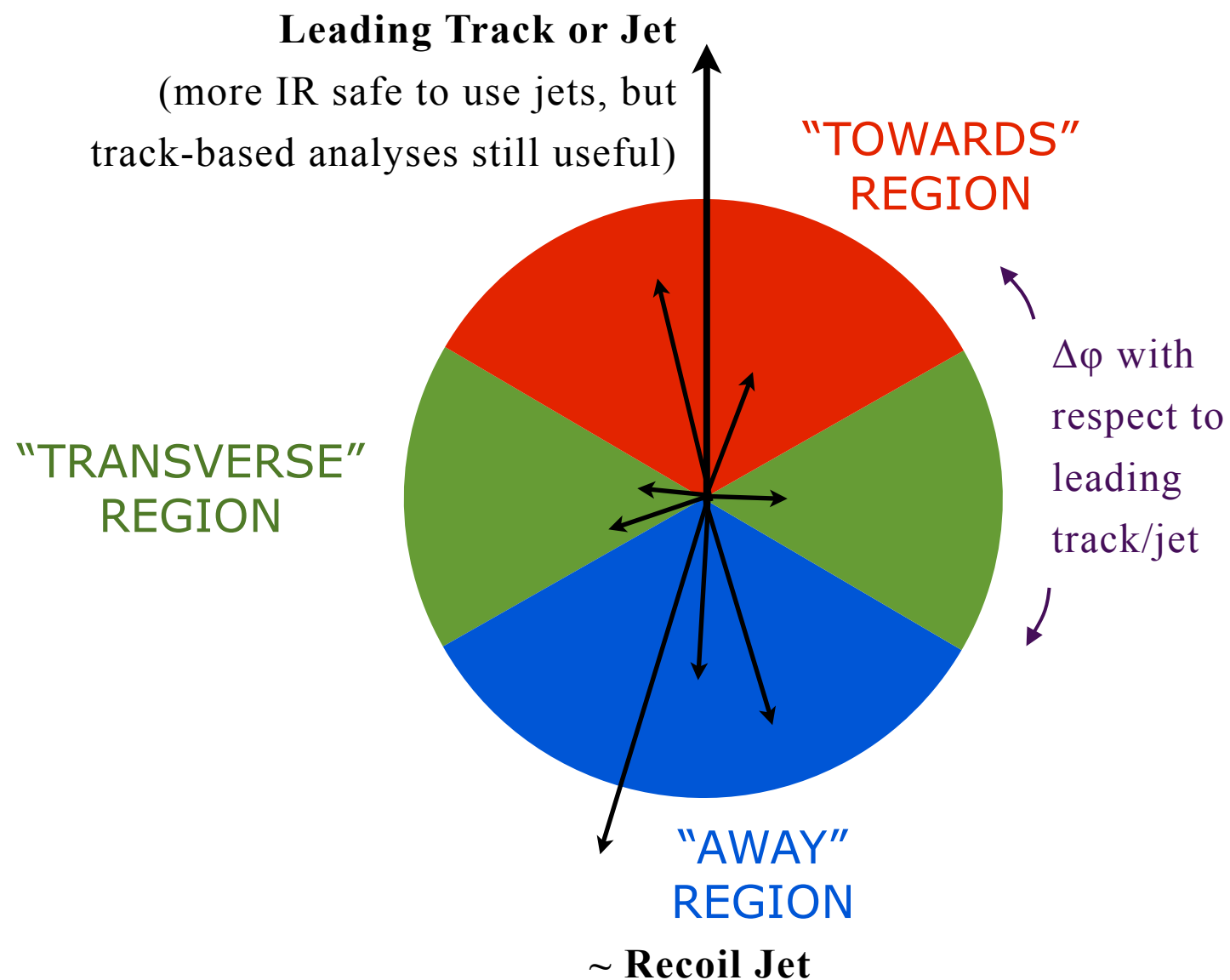
What Cross Section?



The “Rick Field” UE Plots

(the same Field as in Field-Feynman)

There are many UE variables.
The most important is $\langle \Sigma p_T \rangle$ in the “Transverse Region”



Transverse Region (TRNS)

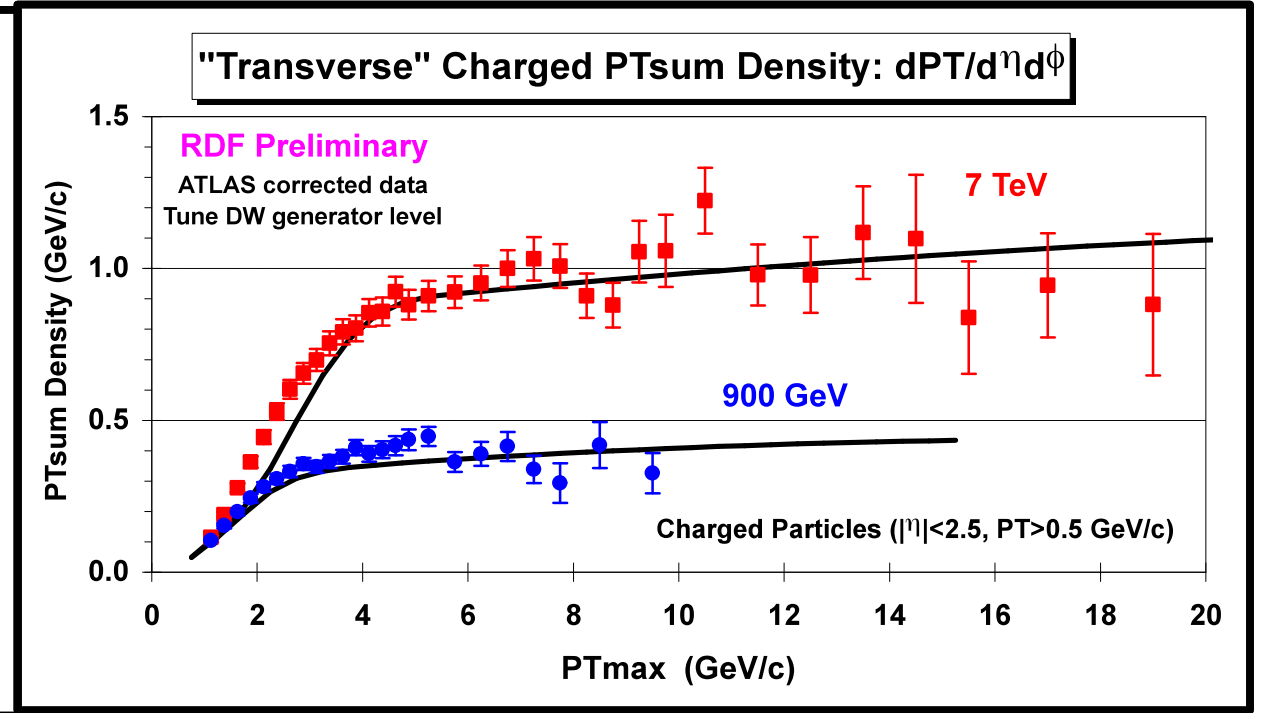
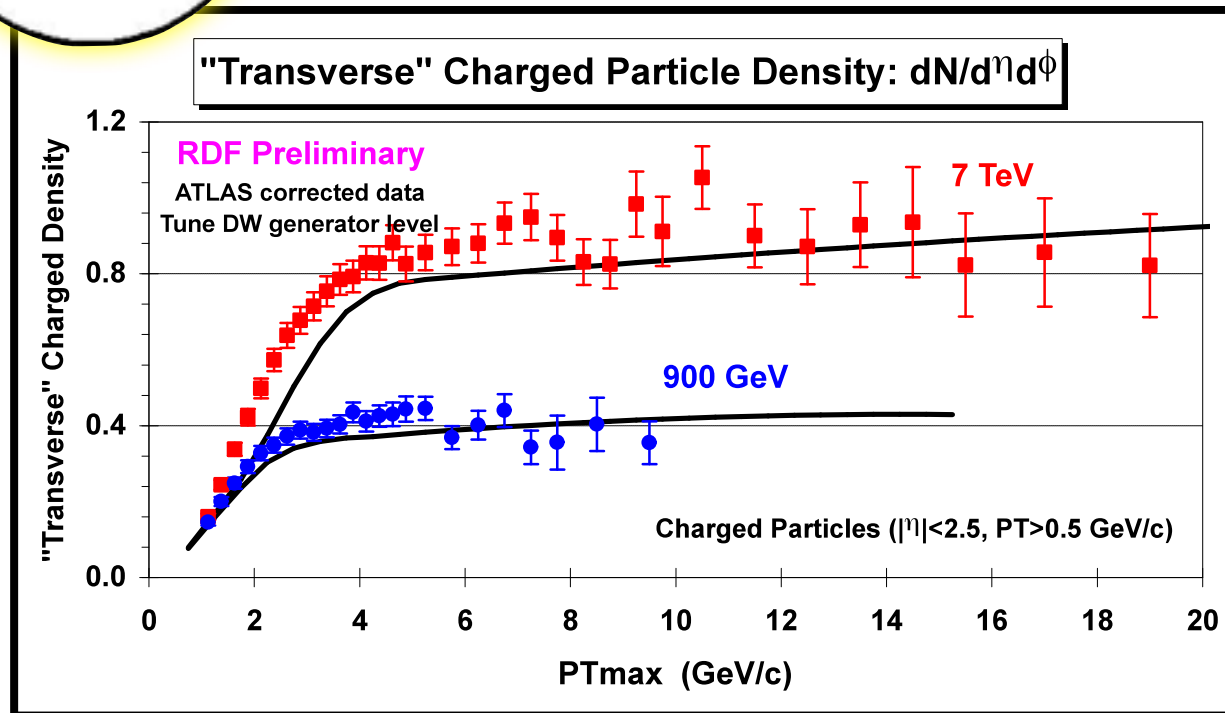
Sensitive to activity at right angles to the hardest jets

Useful definition of Underlying Event

The Pedestal

(now called the Underlying Event)

LHC from 900 to 7000 GeV - ATLAS



Track Density (TRANS)

Not Infrared Safe

Large Non-factorizable Corrections

Prediction off by $\approx 10\%$

Sum(pT) Density (TRANS)

(more) Infrared Safe

Large Non-factorizable Corrections

Prediction off by $< 10\%$

Truth is in the eye of the beholder:

R. Field: "See, I told you!"

Y. Gehrstein: "they have to fudge it again"

From Hard to Soft

Main tools for high- p_T calculations

Factorization and IR safety

Corrections suppressed by powers of $\Lambda_{\text{QCD}}/Q_{\text{Hard}}$

Soft QCD / Min-Bias / Pileup

NO HARD SCALE

Typical Q scales $\sim \Lambda_{\text{QCD}}$

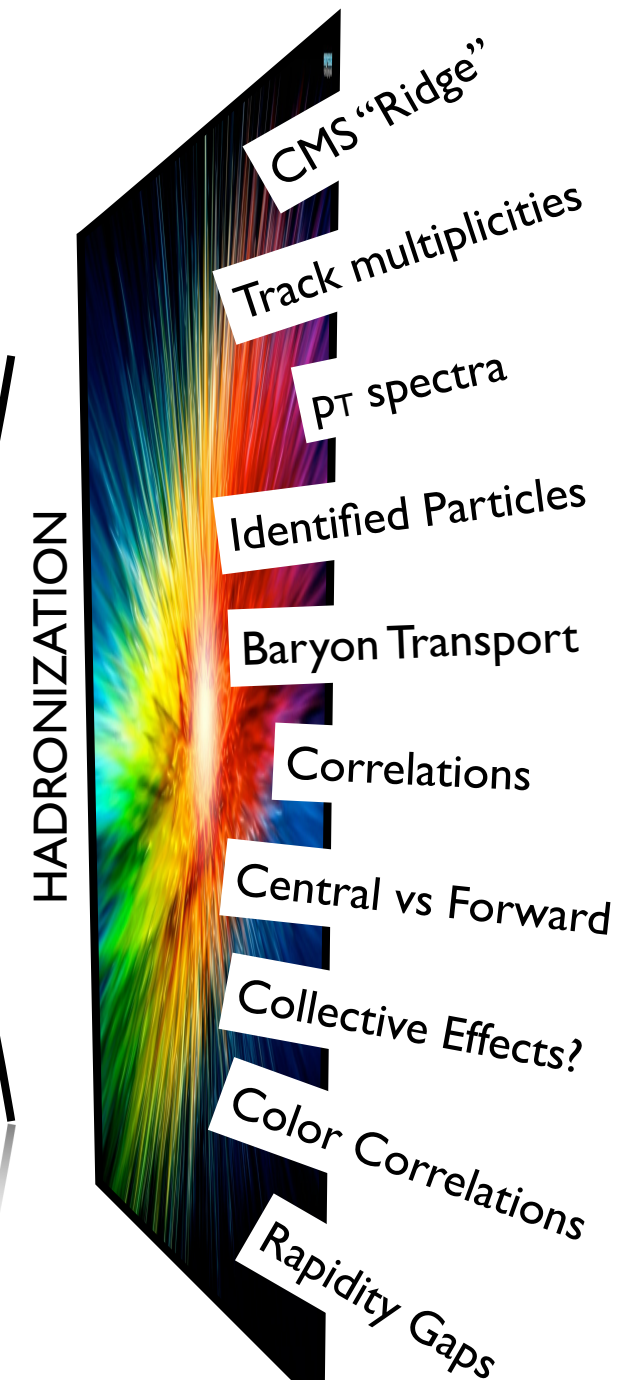
Extremely sensitive to IR effects

→ Excellent LAB for studying IR effects

$\sim \infty$ statistics for min-bias

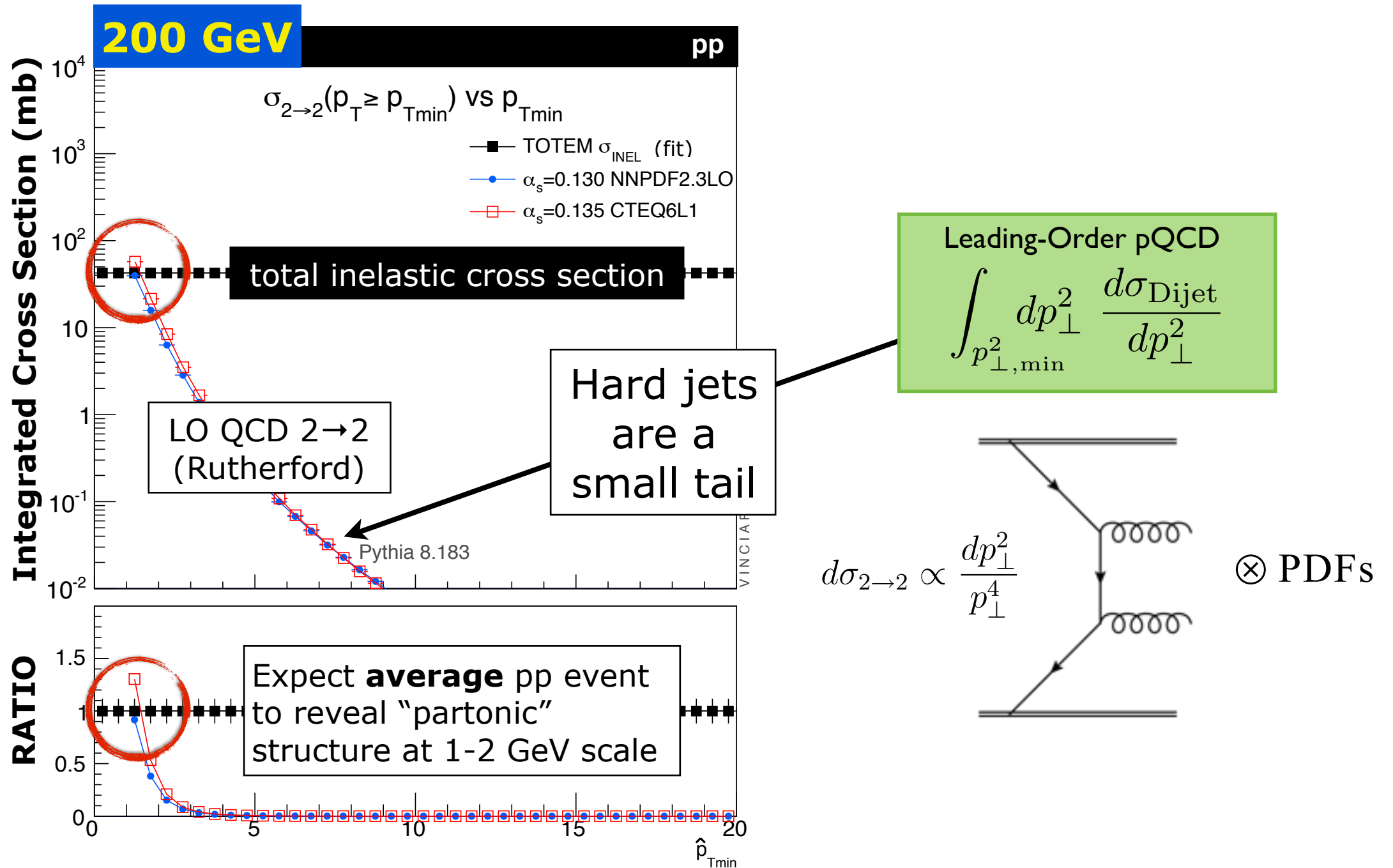
→ Access tails, limits

Universality: Recycling PU ↔ MB ↔ UE



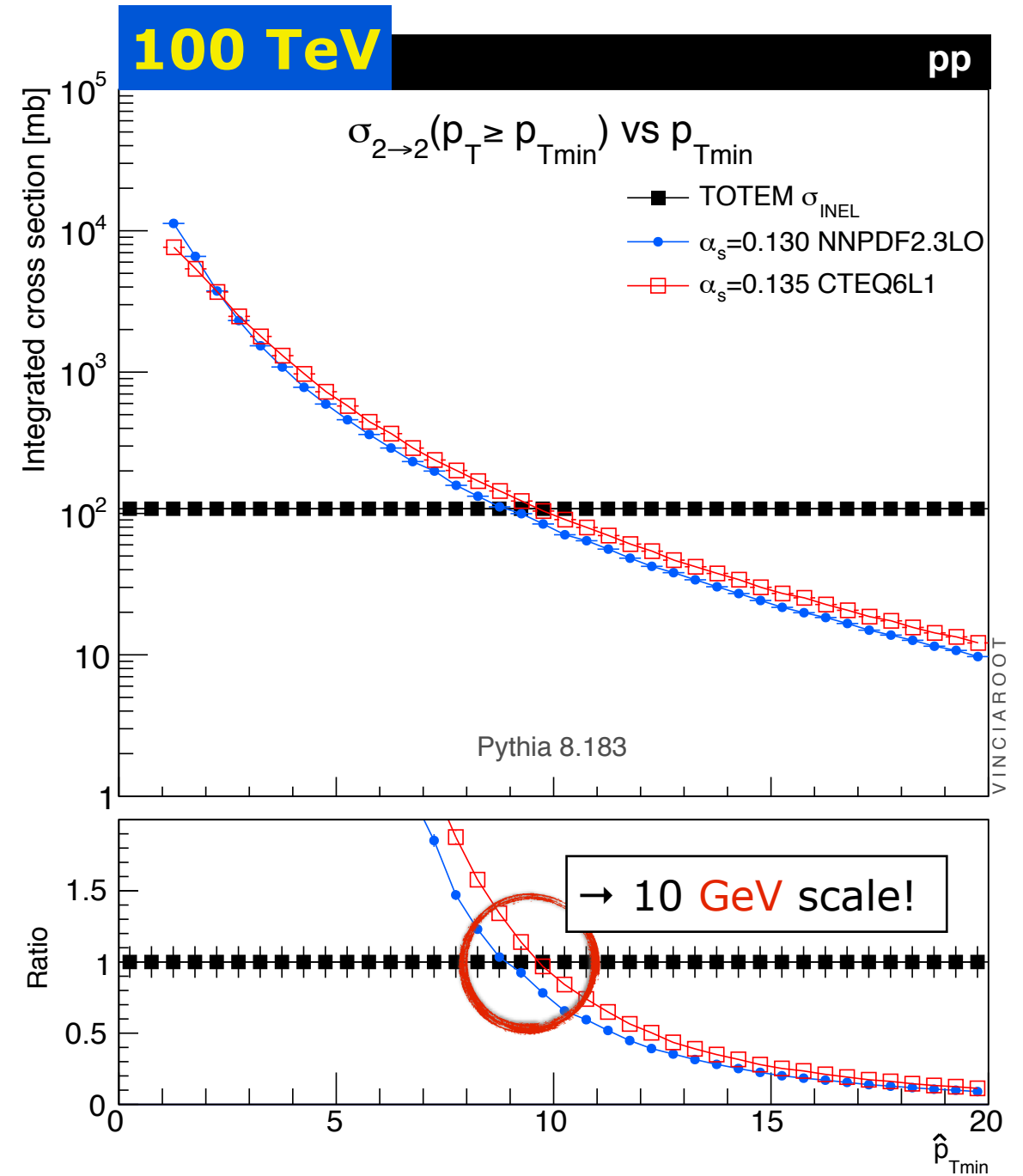
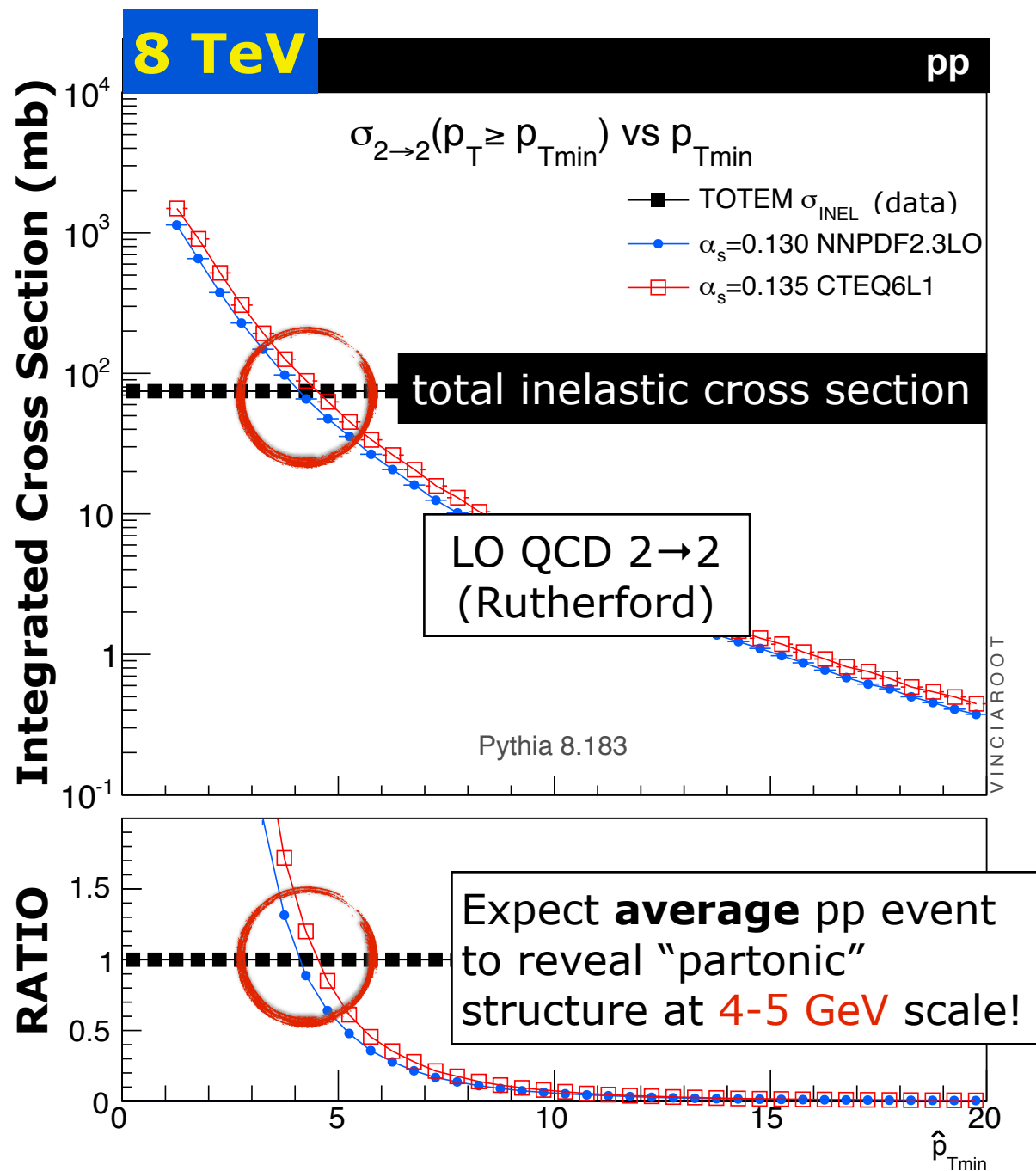
Is there no hard scale?

Compare total (inelastic) hadron-hadron cross section to calculated parton-parton (LO QCD 2→2) cross section



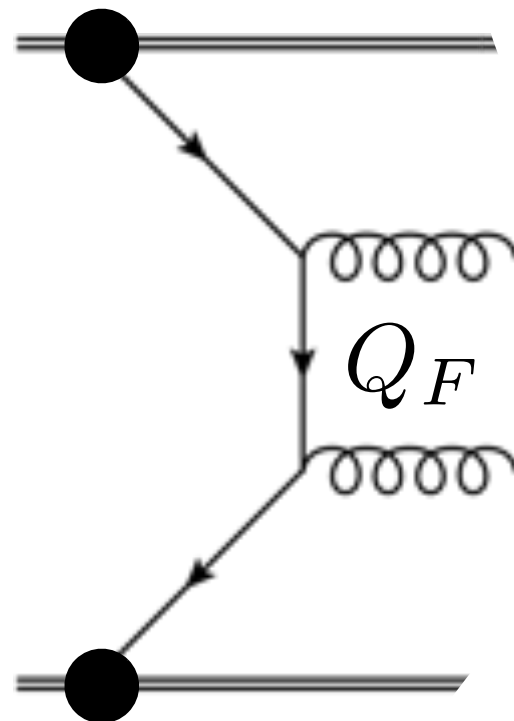
→ 8 TeV → 100 TeV

→ Trivial calculation indicates hard scales in min-bias



Physics of the Pedestal

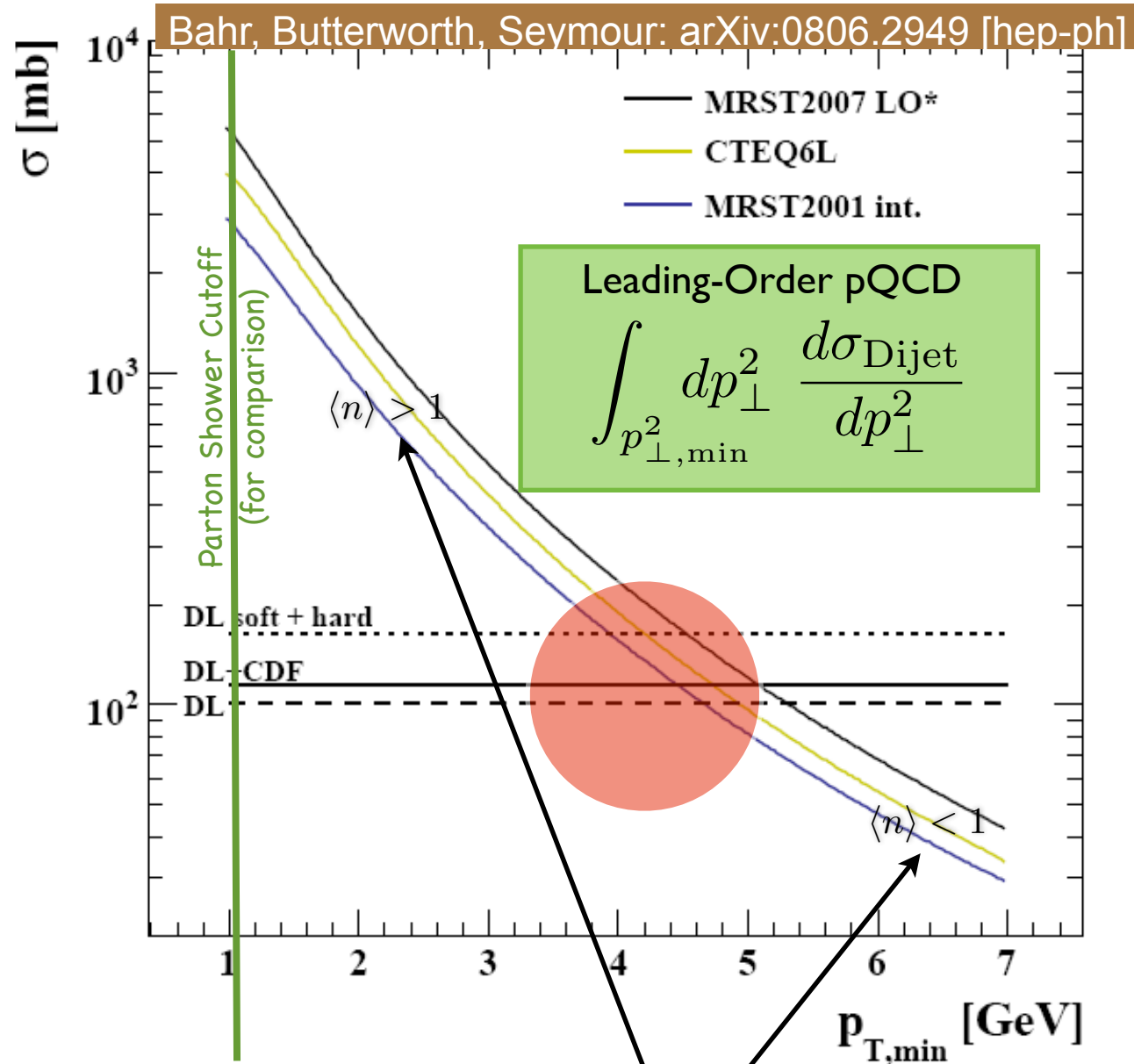
Factorization: Subdivide Calculation



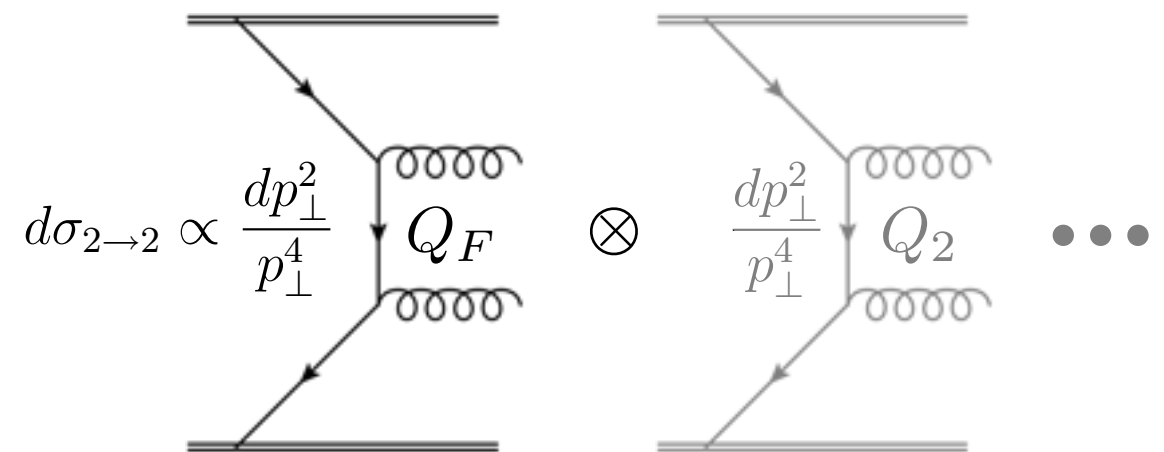
- Multiple Parton Interactions* go beyond existing theorems
- perturbative short-distance physics in Underlying Event
 - Need to generalize factorization to MPI

Multiple Parton Interactions

= Allow several parton-parton interactions per hadron-hadron collision. Requires extended factorization ansatz.



Earliest MC model ("old" PYTHIA 6 model)
Sjöstrand, van Zijl PRD36 (1987) 2019



Lesson from bremsstrahlung in pQCD:
divergences \rightarrow fixed-order breaks down
Perturbation theory still ok, with
resummation (unitarity)

\rightarrow Resum dijets?
Yes \rightarrow MPI!

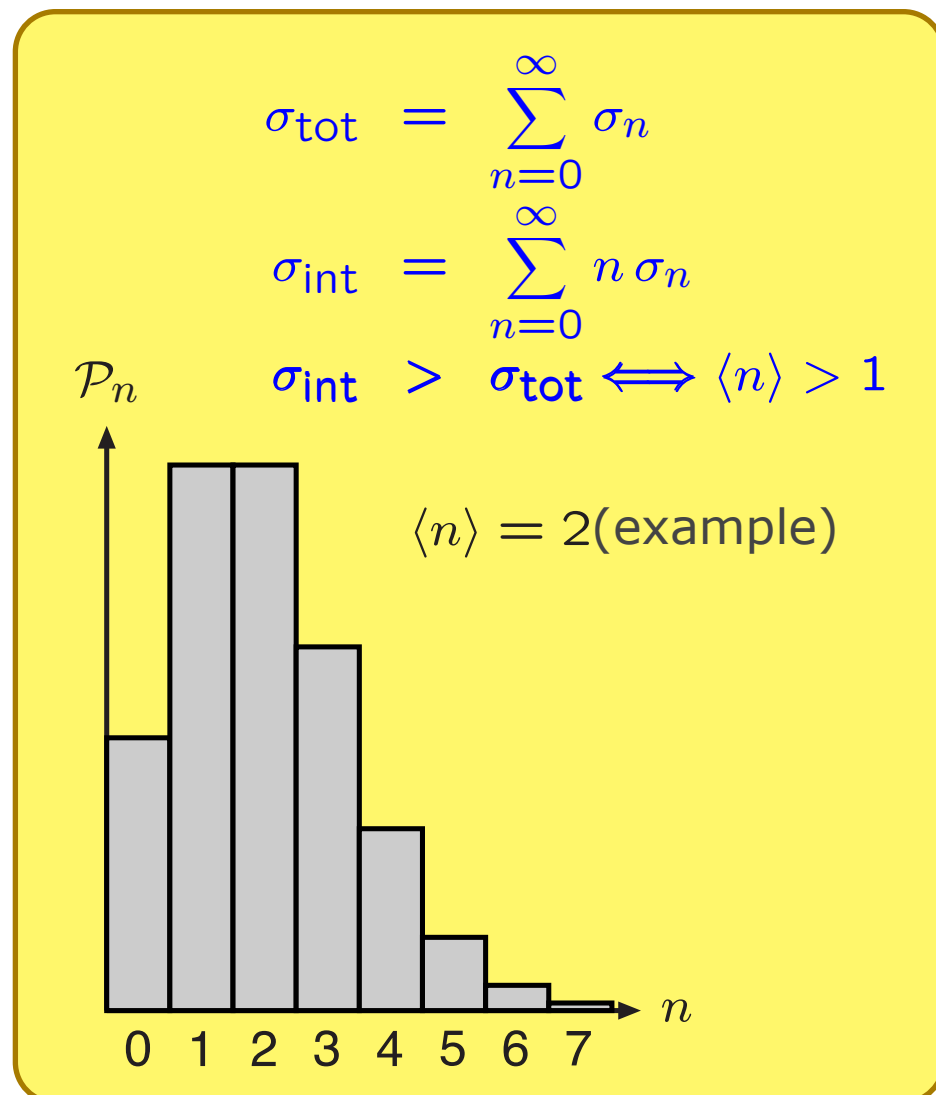
$$\sigma_{2 \rightarrow 2}(p_{\perp \min}) = \langle n \rangle(p_{\perp \min}) \sigma_{\text{tot}}$$

Parton-Parton Cross Section

Hadron-Hadron Cross Section

How many?

Naively $\langle n_{2 \rightarrow 2}(p_{\perp \min}) \rangle = \frac{\sigma_{2 \rightarrow 2}(p_{\perp \min})}{\sigma_{\text{tot}}}$
 Interactions independent (naive factorization) \rightarrow Poisson



$$\mathcal{P}_n = \frac{\langle n \rangle^n}{n!} e^{-\langle n \rangle}$$

Real Life

Color screening: $\sigma_{2 \rightarrow 2} \rightarrow 0$ for $p_{\perp} \rightarrow 0$

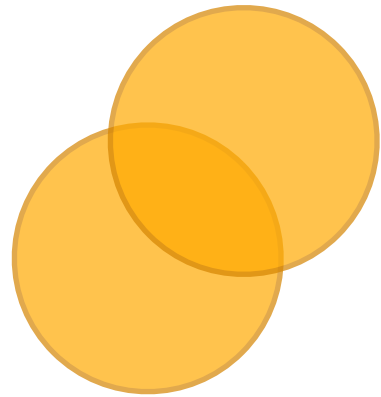
Momentum conservation suppresses high-n tail

Impact-parameter dependence

+ physical correlations

\rightarrow not simple product

Impact Parameter

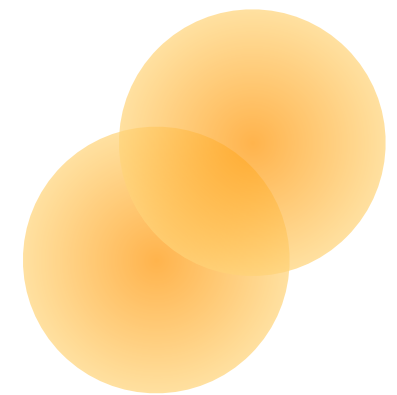


1. **Simple Geometry** (in impact-parameter plane)

Simplest idea: smear PDFs across a uniform disk of size πr_p^2
→ simple geometric overlap factor ≤ 1 in dijet cross section
Some collisions have the full overlap, others only partial
→ Poisson distribution with different mean $\langle n \rangle$ at each b

2. More realistic **Proton b-shape**

Smear PDFs across a non-uniform disk
MC models use Gaussians or **more**/less peaked
Overlap factor = convolution of two such distributions



→ Poisson distribution with different mean $\langle n \rangle$ at each b
“Lumpy Peaks” → large matter overlap enhancements, higher $\langle n \rangle$

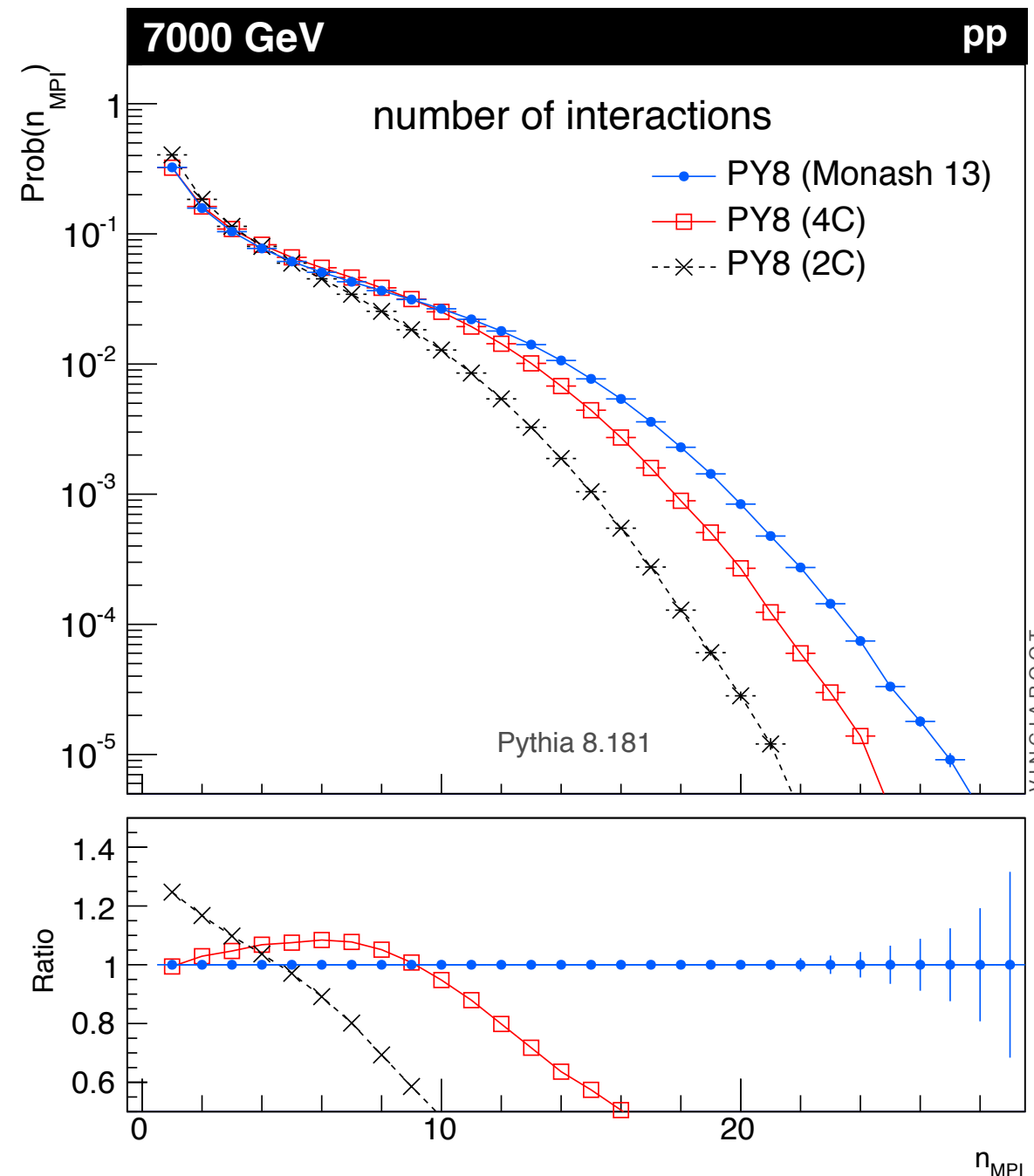
Note: this is an *effective* description. Not the actual proton mass density.
E.g., peak in overlap function ($\gg 1$) can represent unlikely configurations with huge overlap enhancement. Typically use total σ_{inel} as normalization.

Number of MPI *

Minimum-Bias pp collisions at 7 TeV

Averaged over all
pp impact
parameters

(Really:
averaged over all
pp overlap
enhancement
factors)

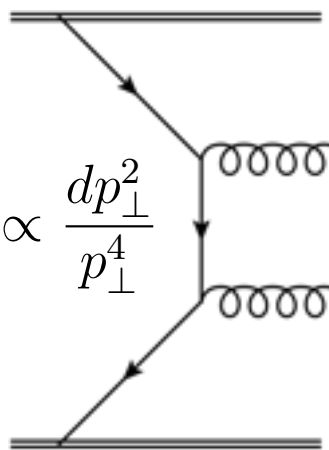


*note: can be
arbitrarily soft

Caveats of MPI-Based Models

Main applications:

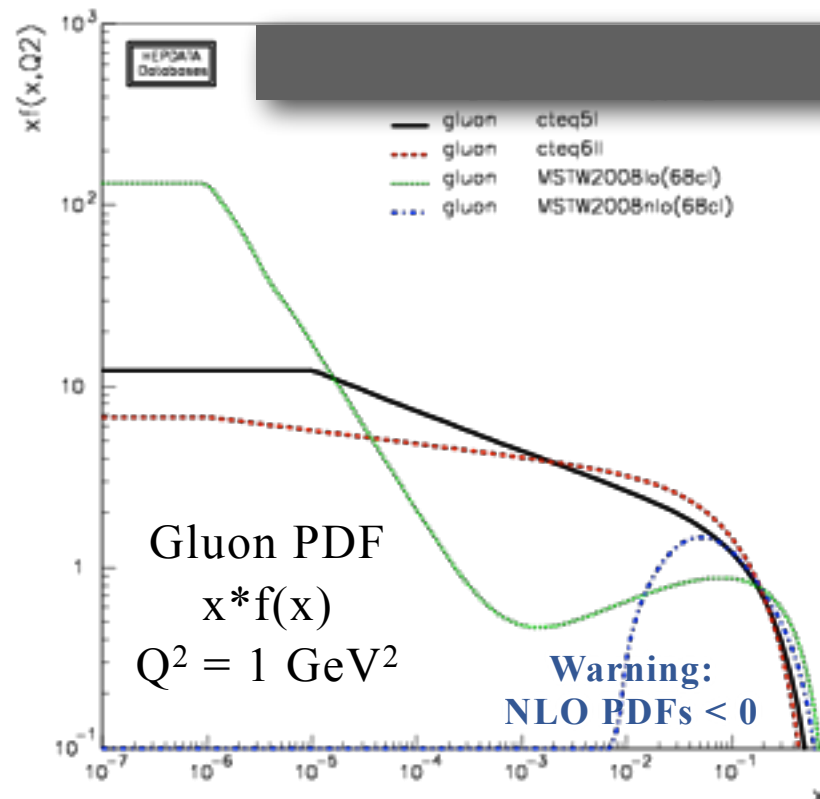
Central Jets/EWK/top/
Higgs/New Physics

$$d\sigma_{2 \rightarrow 2} \propto \frac{dp_{\perp}^2}{p_{\perp}^4} \otimes \text{PDFs}$$


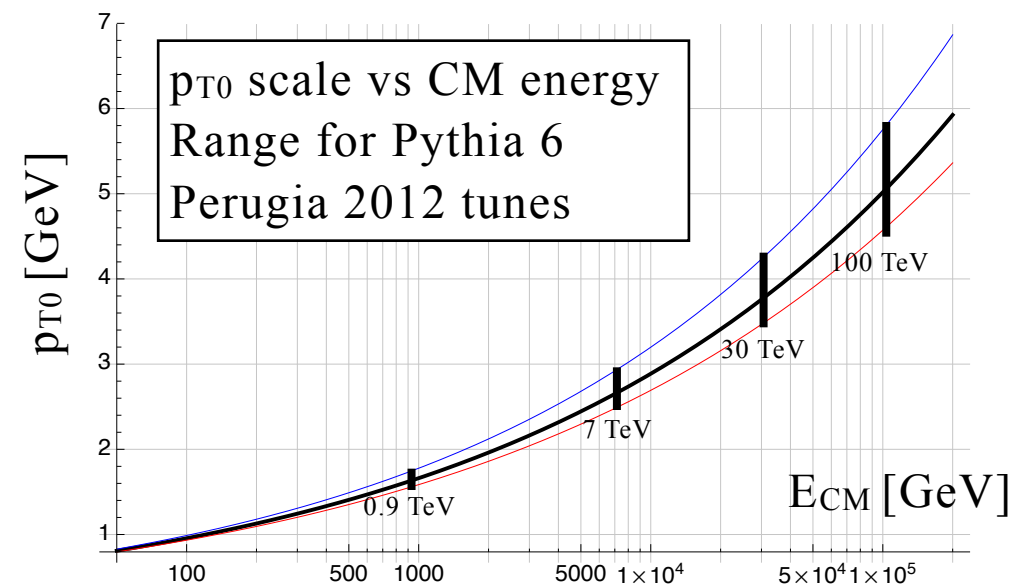
High Q^2
and
finite x

Extrapolation to soft scales delicate.
Impressive successes with MPI-based models but still far from a solved problem

- Form of PDFs at small x and Q^2 ↔ Saturation
- Form and E_{cm} dependence of p_{T0} regulator
- Modeling of the diffractive component
- Proton transverse mass distribution
- Colour Reconnections, Collective Effects



Poor Man's Saturation



See talk on UE
by W. Waalewijn

See also [Connecting hard to soft: KMR, EPJ C71 \(2011\) 1617](#) + [PYTHIA "Perugia Tunes": PS, PRD82 \(2010\) 074018](#) + [arXiv:1308.2813](#)

1: A Simple Model

The minimal model incorporating single-parton factorization, perturbative unitarity, and energy-and-momentum conservation

$$\underbrace{\sigma_{2 \rightarrow 2}(p_{\perp \min})}_{\text{Parton-Parton Cross Section}} = \underbrace{\langle n \rangle(p_{\perp \min})}_{\text{Hadron-Hadron Cross Section}} \sigma_{\text{tot}}$$

1. Choose $p_{T\min}$ cutoff

= main tuning parameter

2. Interpret $\langle n \rangle(p_{T\min})$ as mean of Poisson distribution

Equivalent to assuming all parton-parton interactions equivalent and independent ~ each take an instantaneous “snapshot” of the proton

3. Generate n parton-parton interactions (pQCD $2 \rightarrow 2$)

Veto if total beam momentum exceeded \rightarrow overall (E,p) cons

4. Add impact-parameter dependence $\rightarrow \langle n \rangle = \langle n \rangle(b)$ Ordinary CTEQ, MSTW, NNPDF, ...

Assume factorization of transverse and longitudinal d.o.f., \rightarrow PDFs : $f(x,b) = f(x)g(b)$

b distribution \propto EM form factor \rightarrow **JIMMY model** Butterworth, Forshaw, Seymour Z.Phys. C72 (1996) 637

Constant of proportionality = second main tuning parameter

5. Add separate class of “soft” (zero- p_T) interactions representing

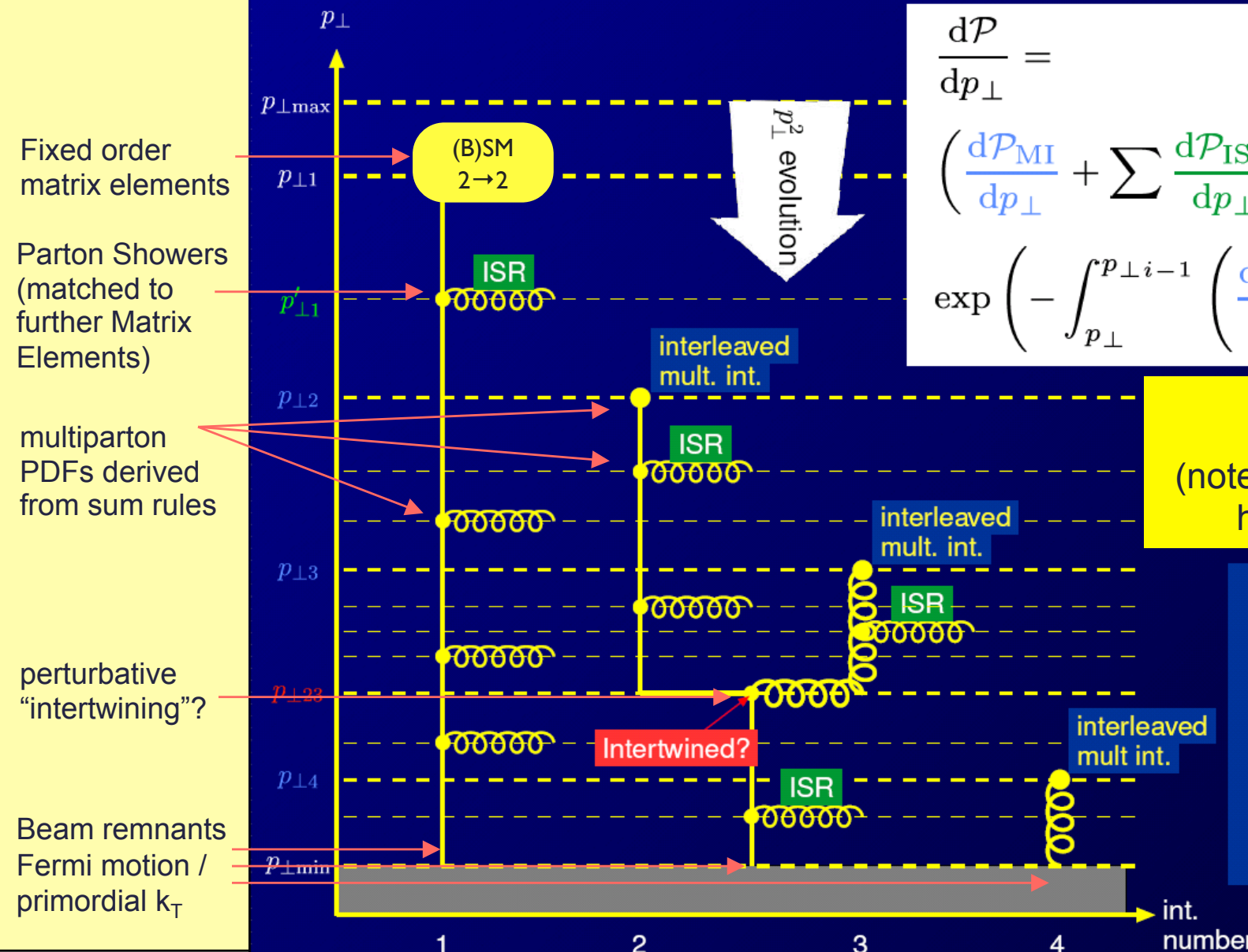
interactions with $p_T < p_{T\min}$ and require $\sigma_{\text{soft}} + \sigma_{\text{hard}} = \sigma_{\text{tot}}$

\rightarrow **Herwig++ model** Bähr et al, arXiv:0905.4671

2: Interleaved Evolution

Sjöstrand, P.S., JHEP 0403 (2004) 053; EPJ C39 (2005) 129

Add exclusivity progressively by evolving *everything* downwards.

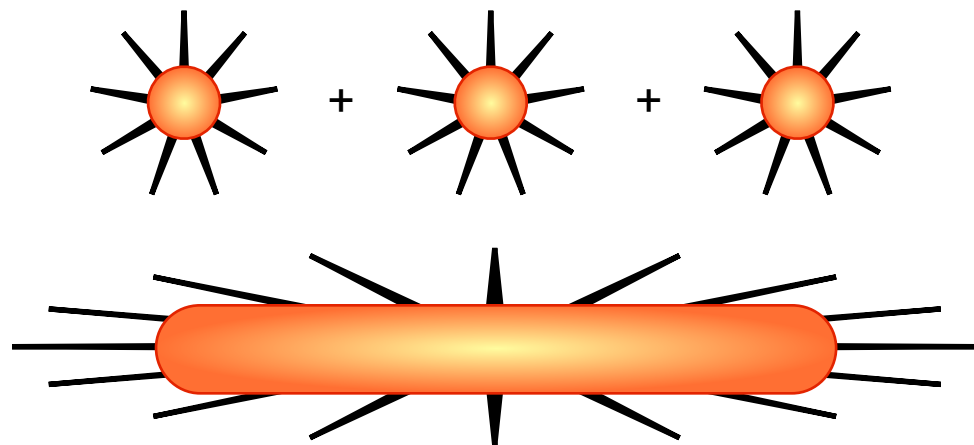


$$\frac{d\mathcal{P}}{dp_{\perp}} = 1 \left(\frac{d\mathcal{P}_{\text{MI}}}{dp_{\perp}} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp_{\perp}} + \sum \frac{d\mathcal{P}_{\text{JI}}}{dp_{\perp}} \right) \times \exp \left(- \int_{p_{\perp}}^{p_{\perp}^{i-1}} \left(\frac{d\mathcal{P}_{\text{MI}}}{dp'_{\perp}} + \sum \frac{d\mathcal{P}_{\text{ISR}}}{dp'_{\perp}} + \sum \frac{d\mathcal{P}_{\text{JI}}}{dp'_{\perp}} \right) dp'_{\perp} \right)$$

→ Underlying Event
(note: interactions correlated in colour: hadronization not independent)

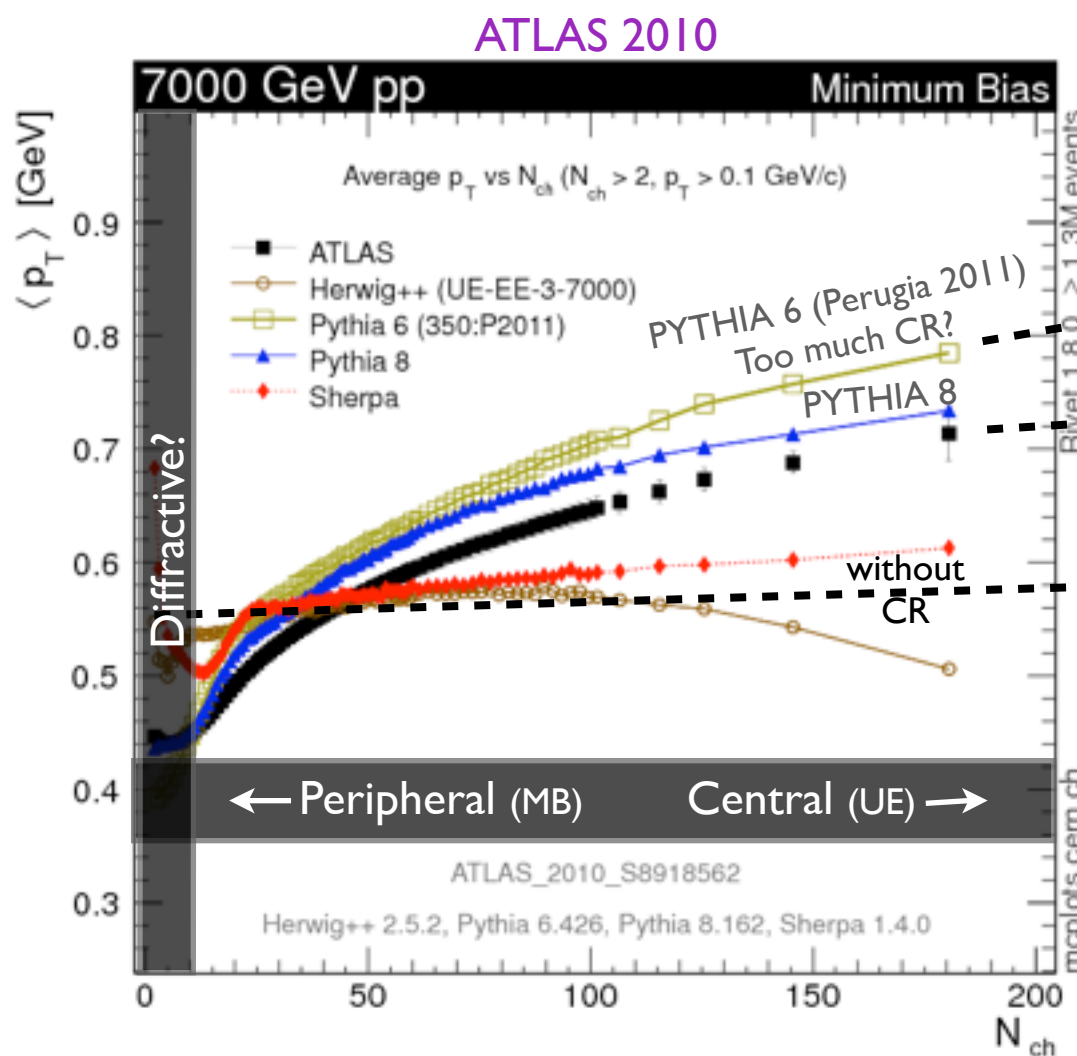
~ "Finegraining"
→ correlations between all perturbative activity at successively smaller scales

$\langle p_T \rangle$ vs N_{ch}



Independent Particle Production:
→ averages stay the same

Correlations / Collective effects:
→ average rises



Extrapolation to high multiplicity ~ UE

Average particles slightly too hard

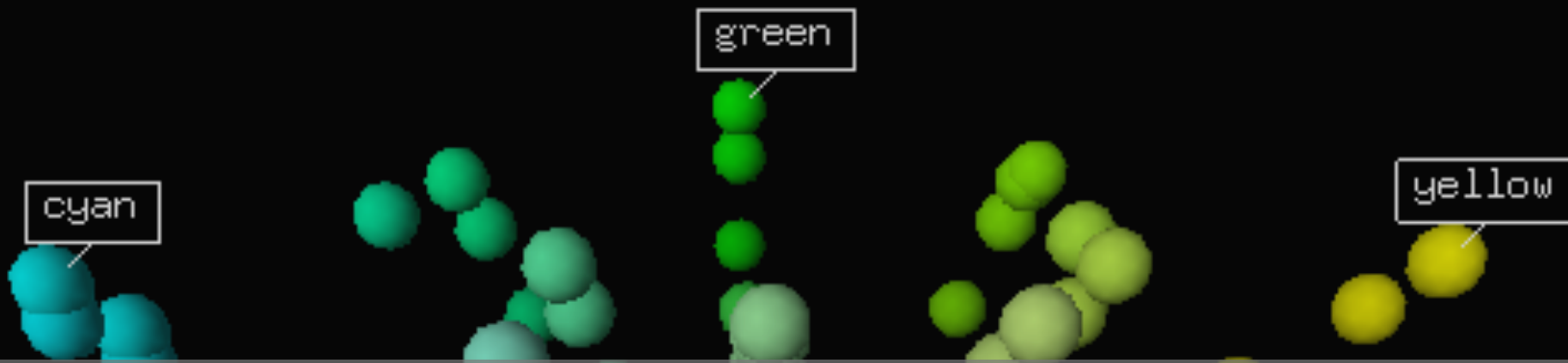
→ Too much energy, or energy distributed on too few particles

~ OK?

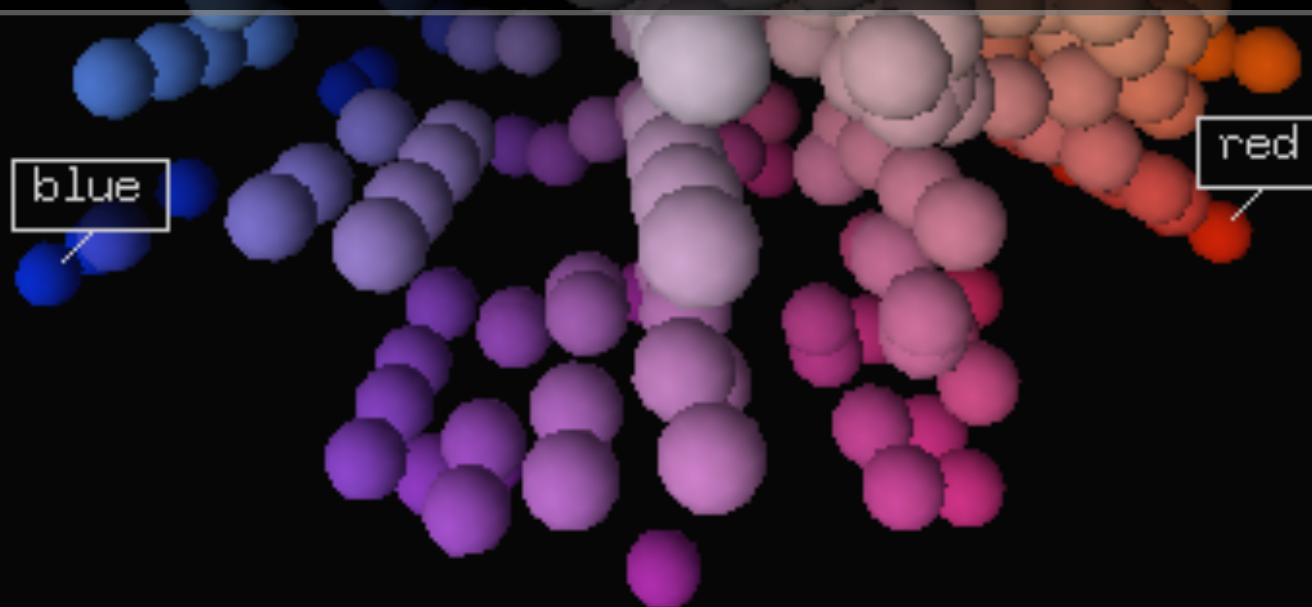
Average particles slightly too soft

→ Too little energy, or energy distributed on too many particles

Evolution of other distributions with N_{ch} also interesting: e.g., $\langle p_T \rangle(N_{ch})$ for identified particles, strangeness & baryon ratios, 2P correlations, ...



Color Space in hadron collisions



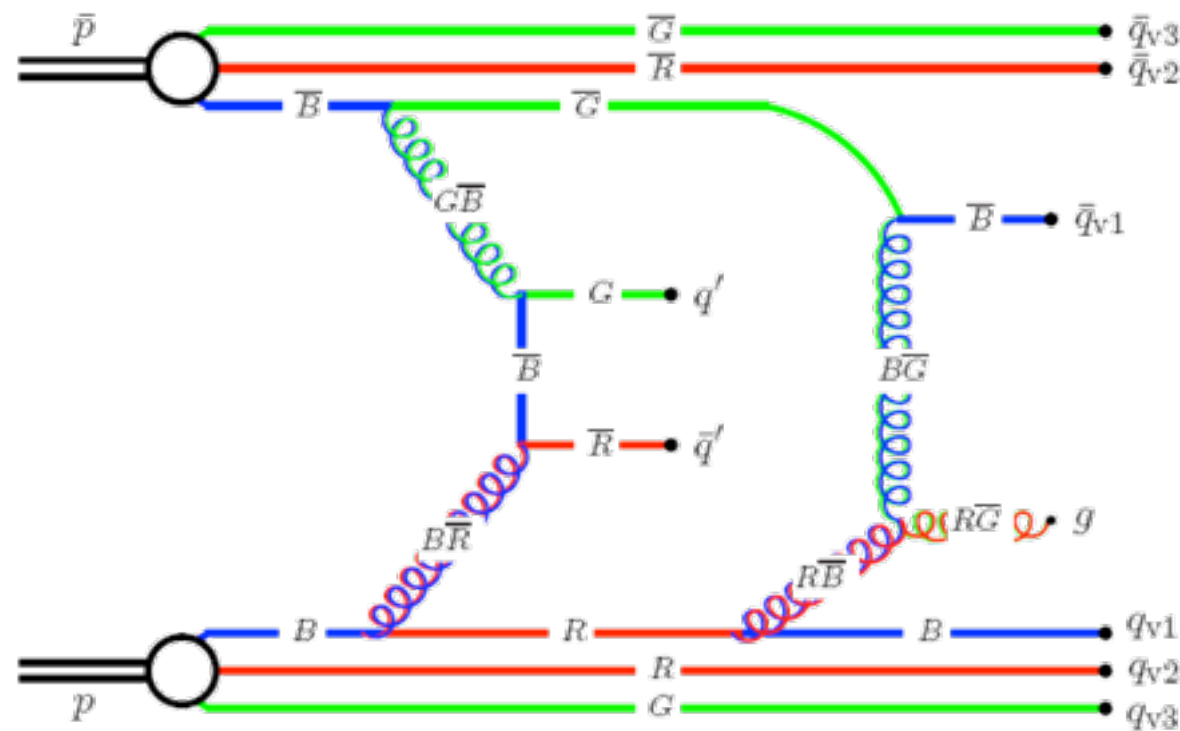
Color Correlations

Each MPI (or cut Pomeron) exchanges color between the beams

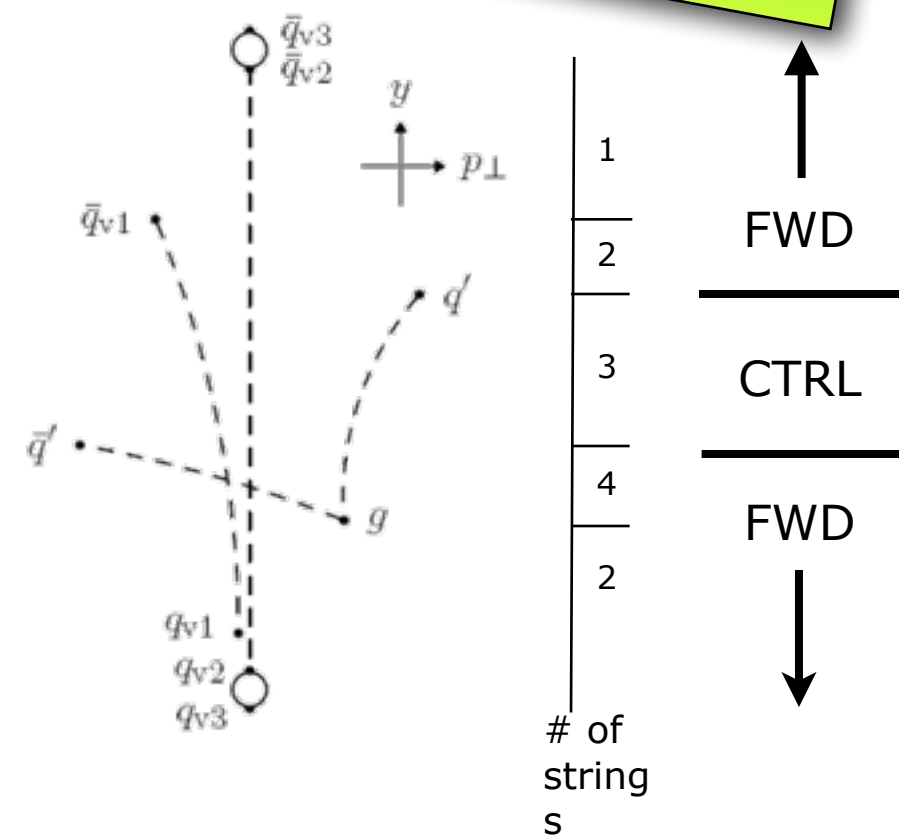
► The colour flow determines the hadronizing string topology

- Each MPI, even when soft, is a color spark
- Final distributions crucially depend on color space

Different models make different ansätze



Sjöstrand & PS, JHEP 03(2004)053



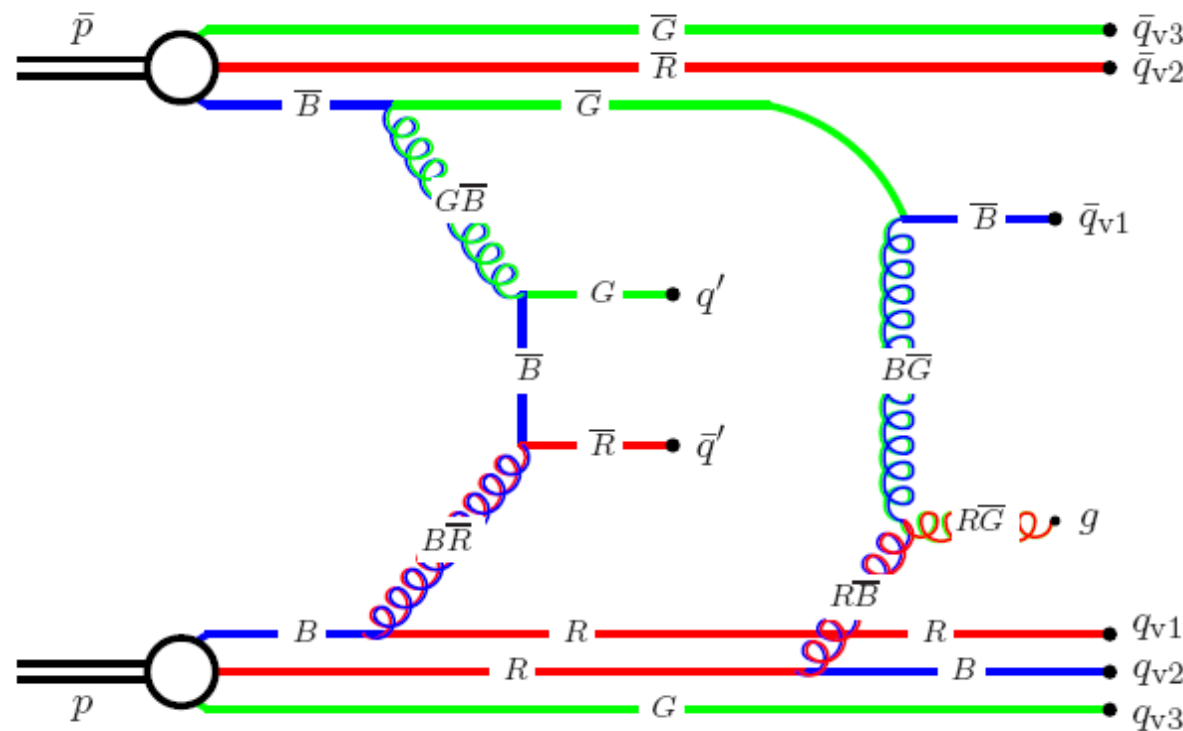
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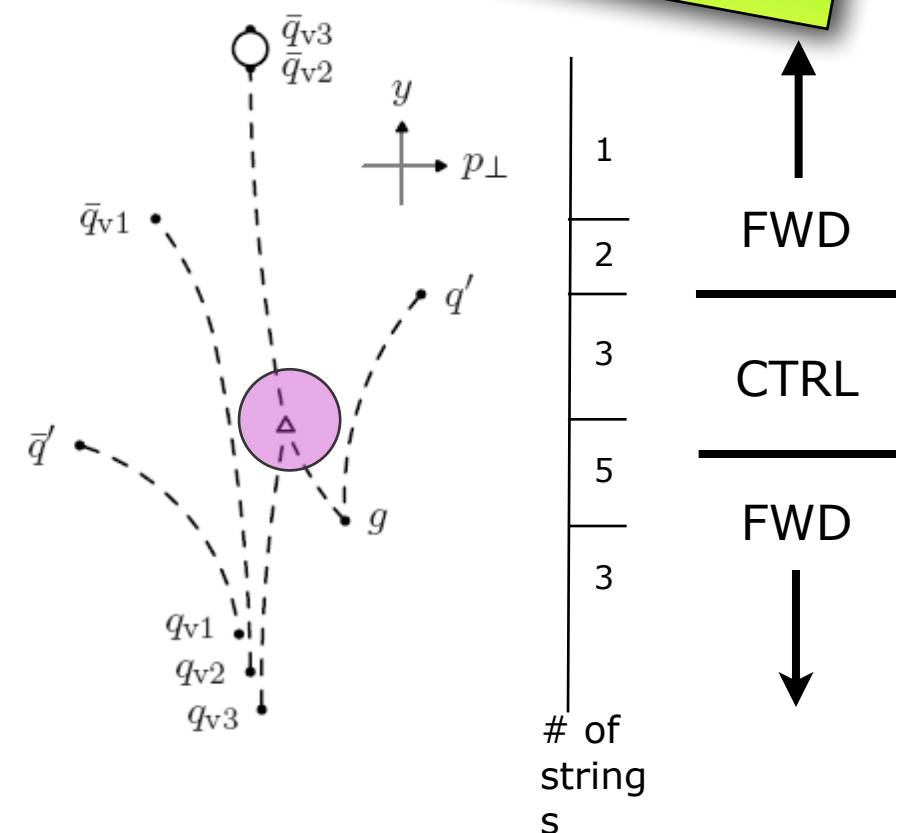
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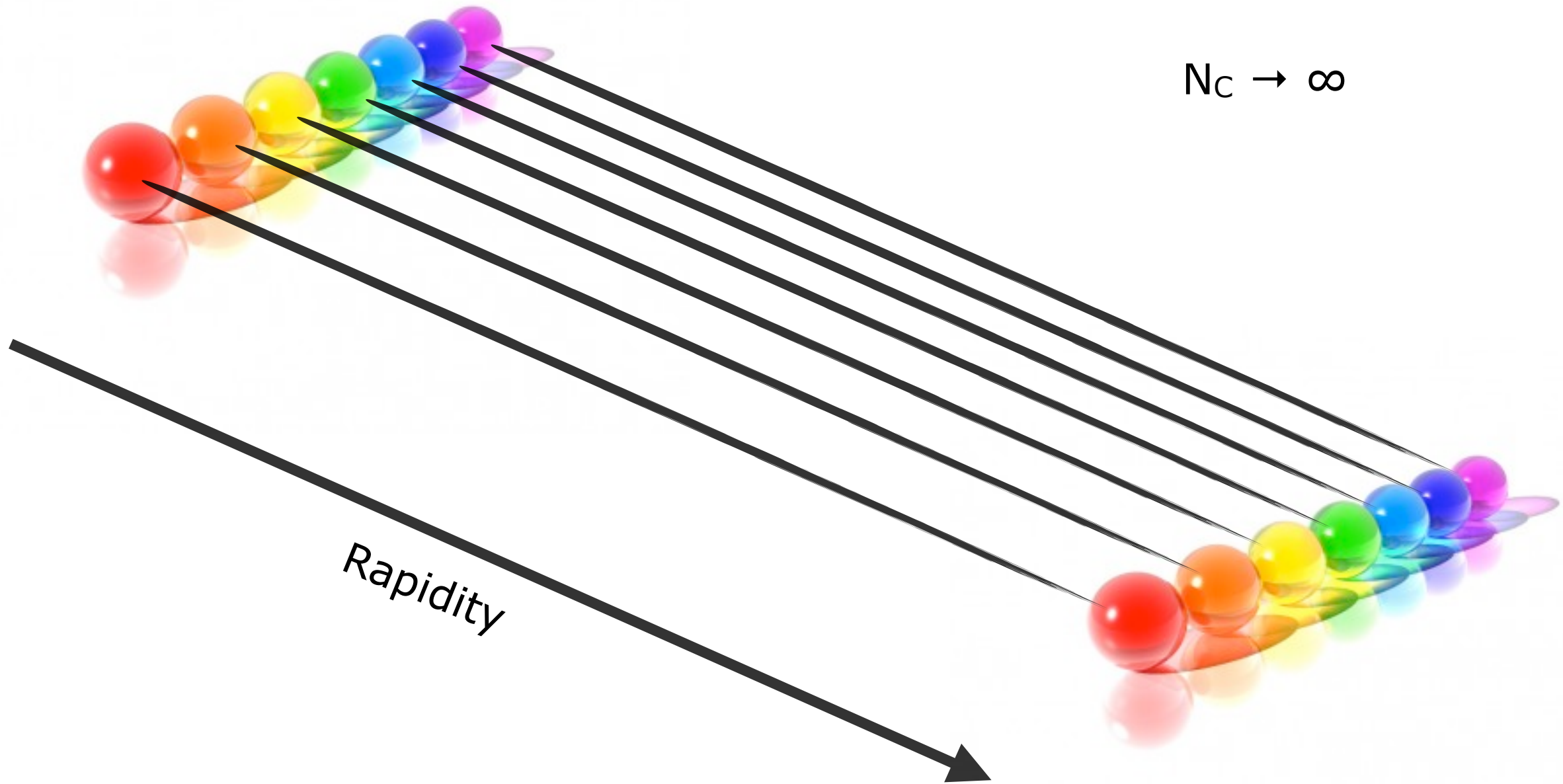
Sjöstrand & PS, JHEP 03(2004)053



Color Connections

Better theory models needed

$$N_c \rightarrow \infty$$



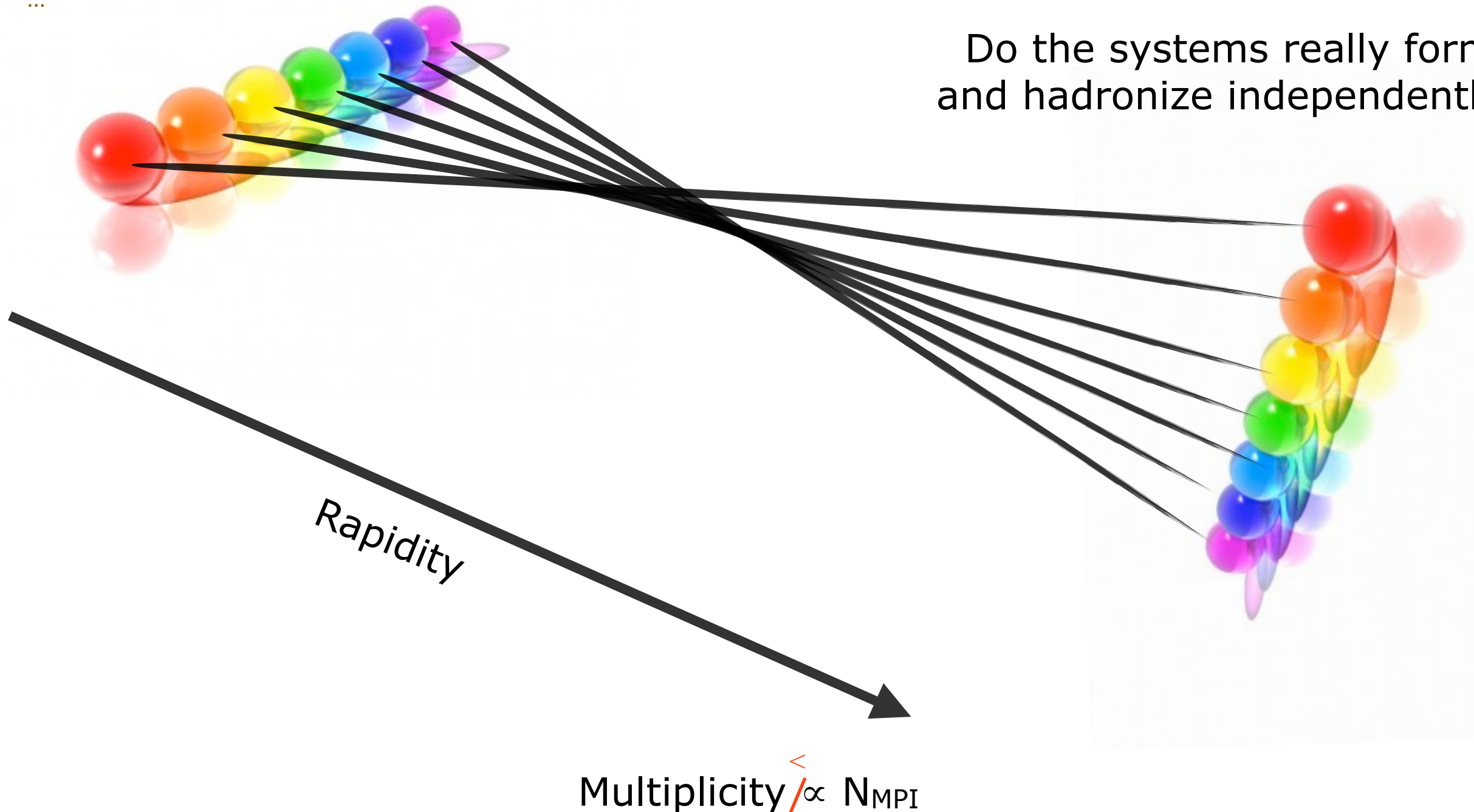
$$\text{Multiplicity} \propto N_{\text{MPI}}$$

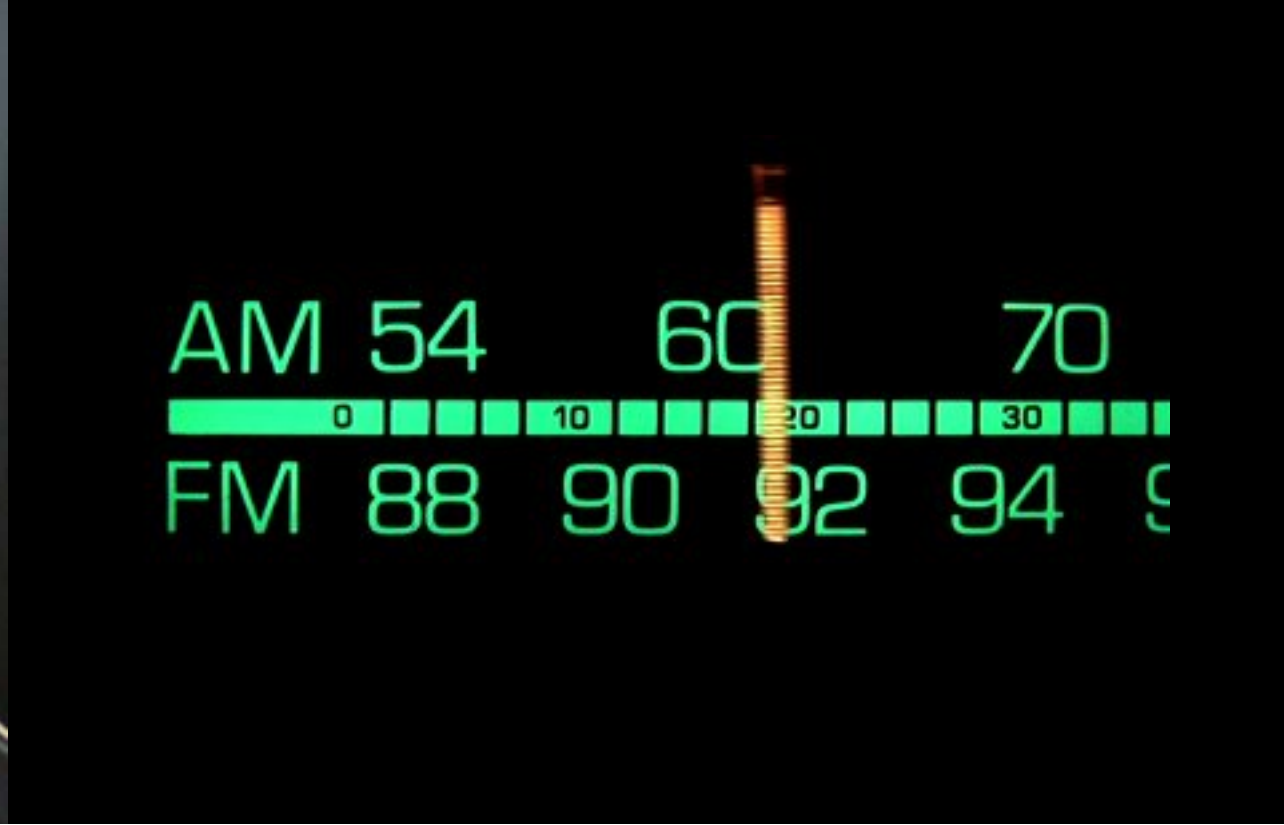
Color Reconnections?

E.g.,
Generalized Area Law (Rathsman: Phys. Lett. B452 (1999) 364)
Color Annealing (P.S., Wicke: Eur. Phys. J. C52 (2007) 133)
...

Better theory models needed

Do the systems really form
and hadronize independently?





Tuning
means different things to different people



MCnet Studentships

MCnet projects:

- PYTHIA (+ VINCIA)
- HERWIG
- SHERPA
- MadGraph
- Ariadne (+ DIPSY)
- Cedar (Rivet/Professor)

Activities include

- summer schools
(2014: Manchester?)
- short-term studentships
- graduate students
- postdocs
- meetings (open/closed)

Monte Carlo training studentships



3-6 month fully funded studentships for current PhD students at one of the MCnet nodes. An excellent opportunity to really understand and improve the Monte Carlos you use!

Application rounds every 3 months.



for details go to:
www.montecarlonet.org

Come to Australia



P



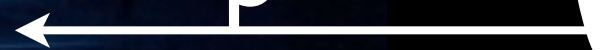
Establishing a new group in **Melbourne**
Working on Precision LHC **phenomenology & soft physics**
PYTHIA & VINCIA

NLO Event Generators

Support LHC **experiments**, **astro-particle** community, and **future** accelerators

Outreach and Citizen Science

P



Soon Advertising:
1 post doc in theoretical physics
2 PhD scholarships in QCD pheno
(1 joint with Warwick ATLAS group, UK)
+ you can apply for Monash scholarships

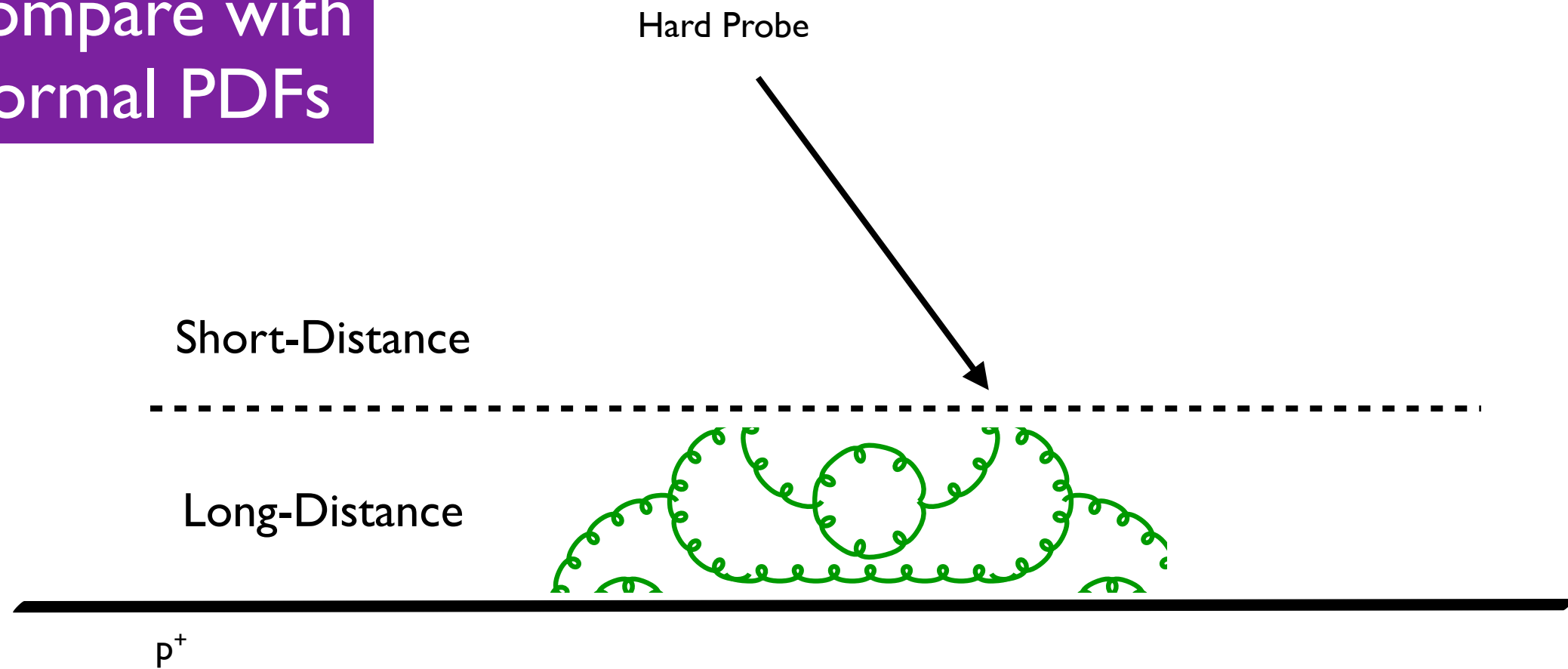
Oct 2014

→ Monash University
Melbourne, Australia

(+ Diffraction)

“Intuitive picture”

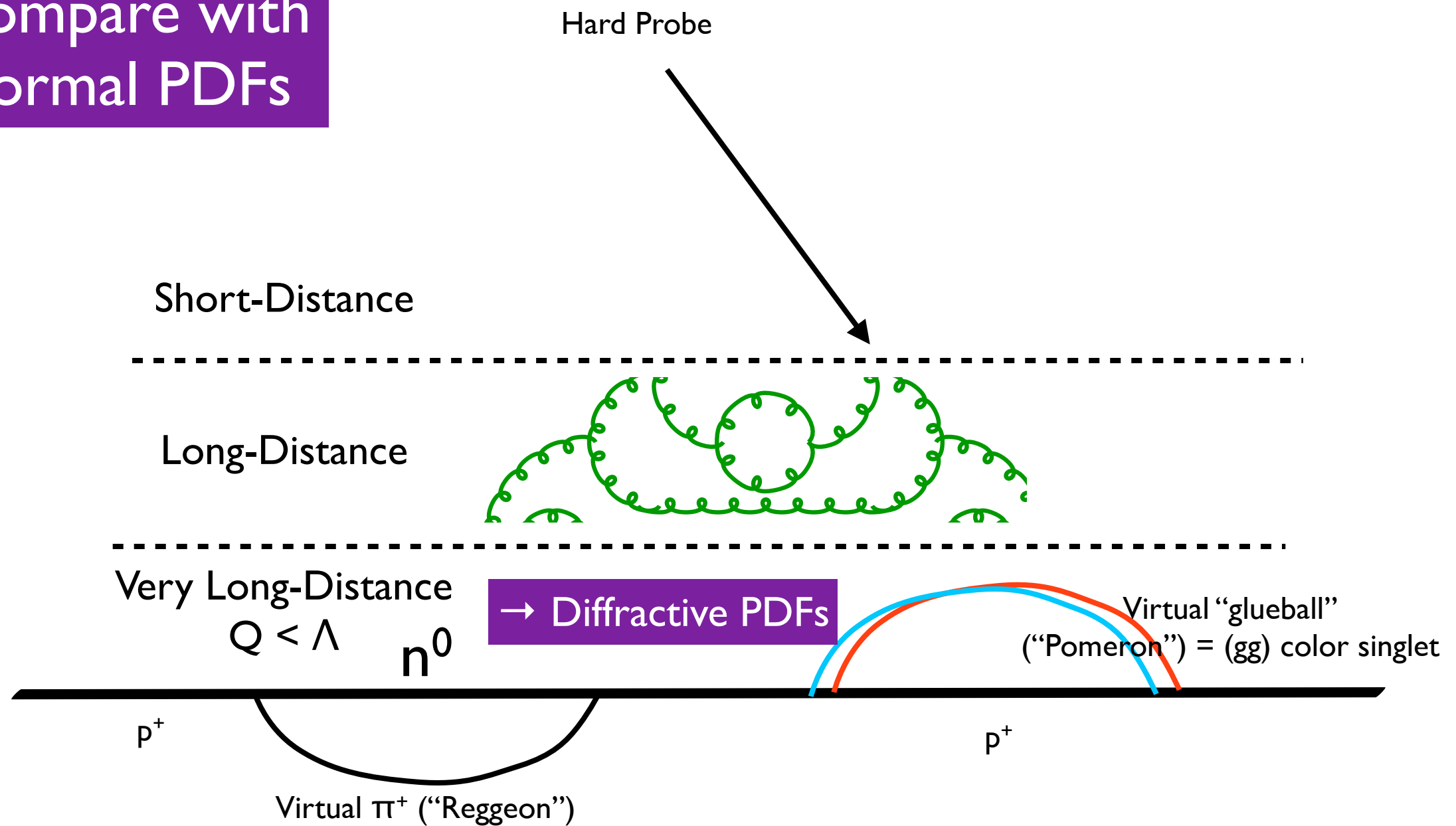
Compare with
normal PDFs



(+ Diffraction)

“Intuitive picture”

Compare with normal PDFs



(+ Diffraction)

“Intuitive picture”

Compare with normal PDFs

