Hadronization & Underlying Event

Peter Skands (CERN Theoretical Physics Dept)

Terascale Monte Carlo School
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Lectures 4+5

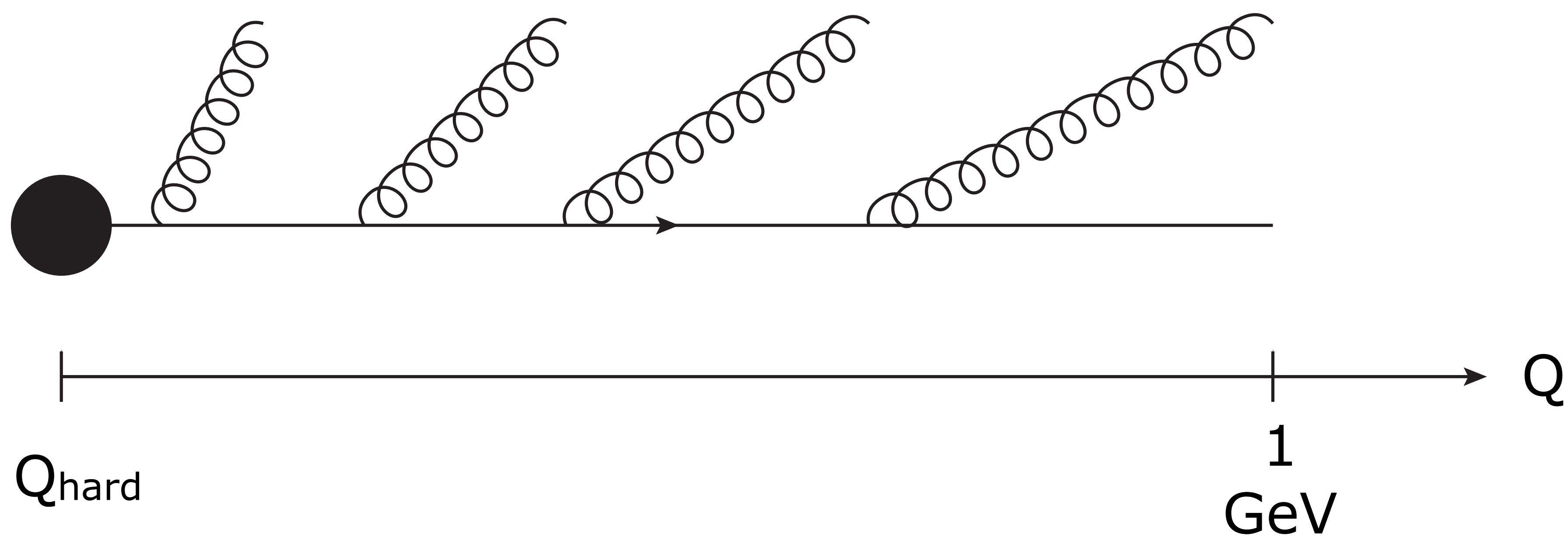
From Partons to Pions

Here's a fast parton

Fast: It starts at a high factorization scale $Q = Q_F = Q_{hard}$

It showers (perturbative bremsstrahlung)

It ends up at a low effective factorization scale $Q \sim m_\rho \sim 1$ GeV



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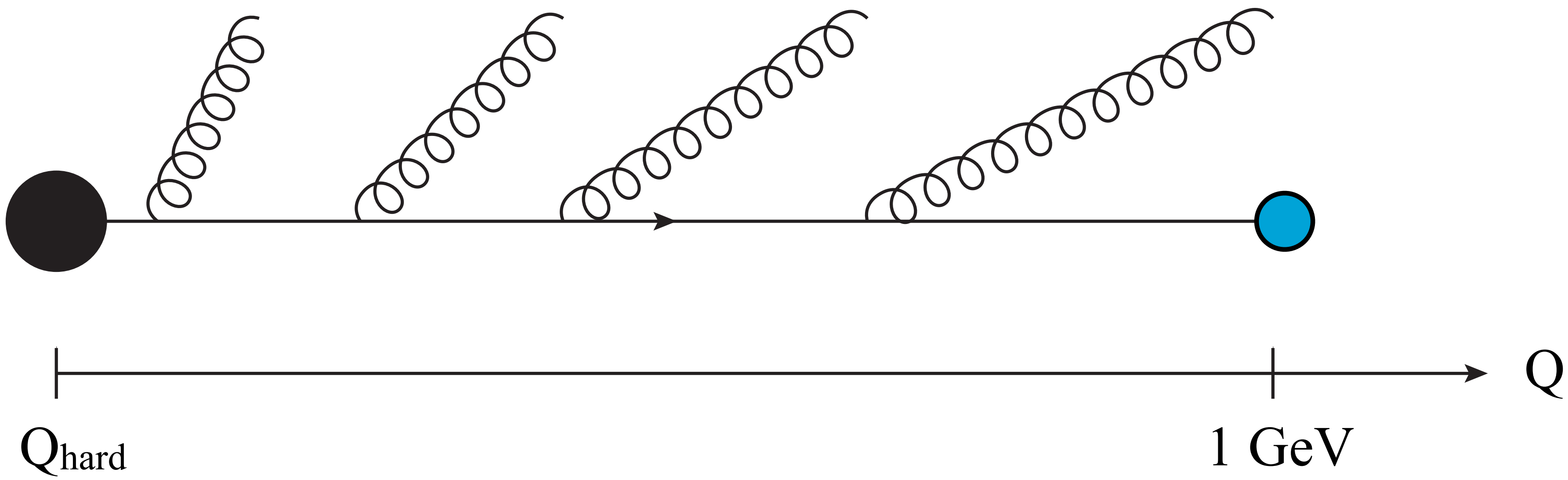

How about I just call it a hadron?

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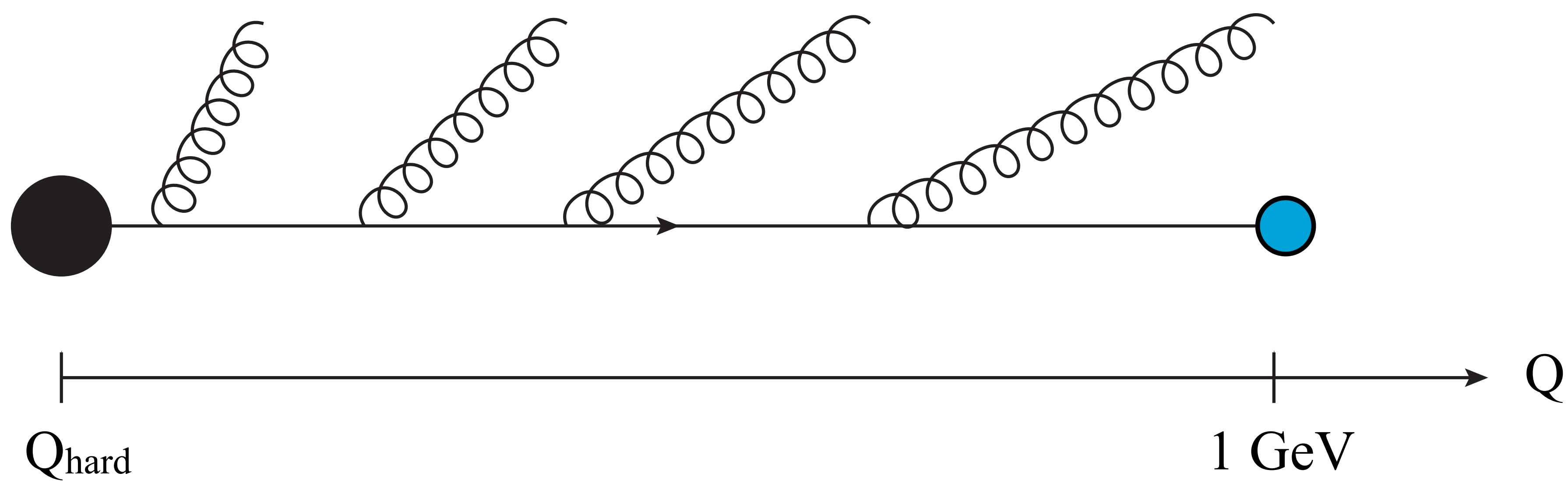


How about I just call it a hadron?

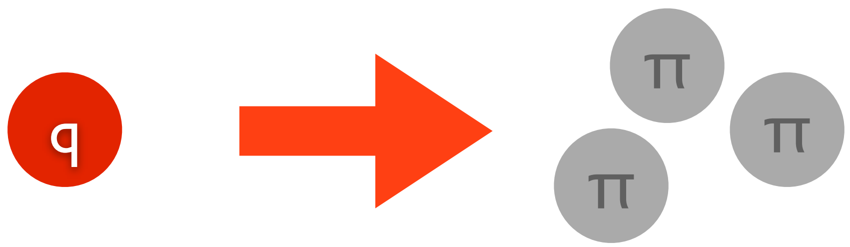
→ "Local Parton-Hadron Duality"

Parton → Hadrons?

Early models: "Independent Fragmentation"

Local Parton Hadron Duality (LPHD) can give useful results for **inclusive** quantities in collinear fragmentation

Motivates a simple model:



But …

The point of confinement is that partons are coloured

Hadronization = the process of colour neutralization

→ Unphysical to think about independent fragmentation of a single parton into hadrons

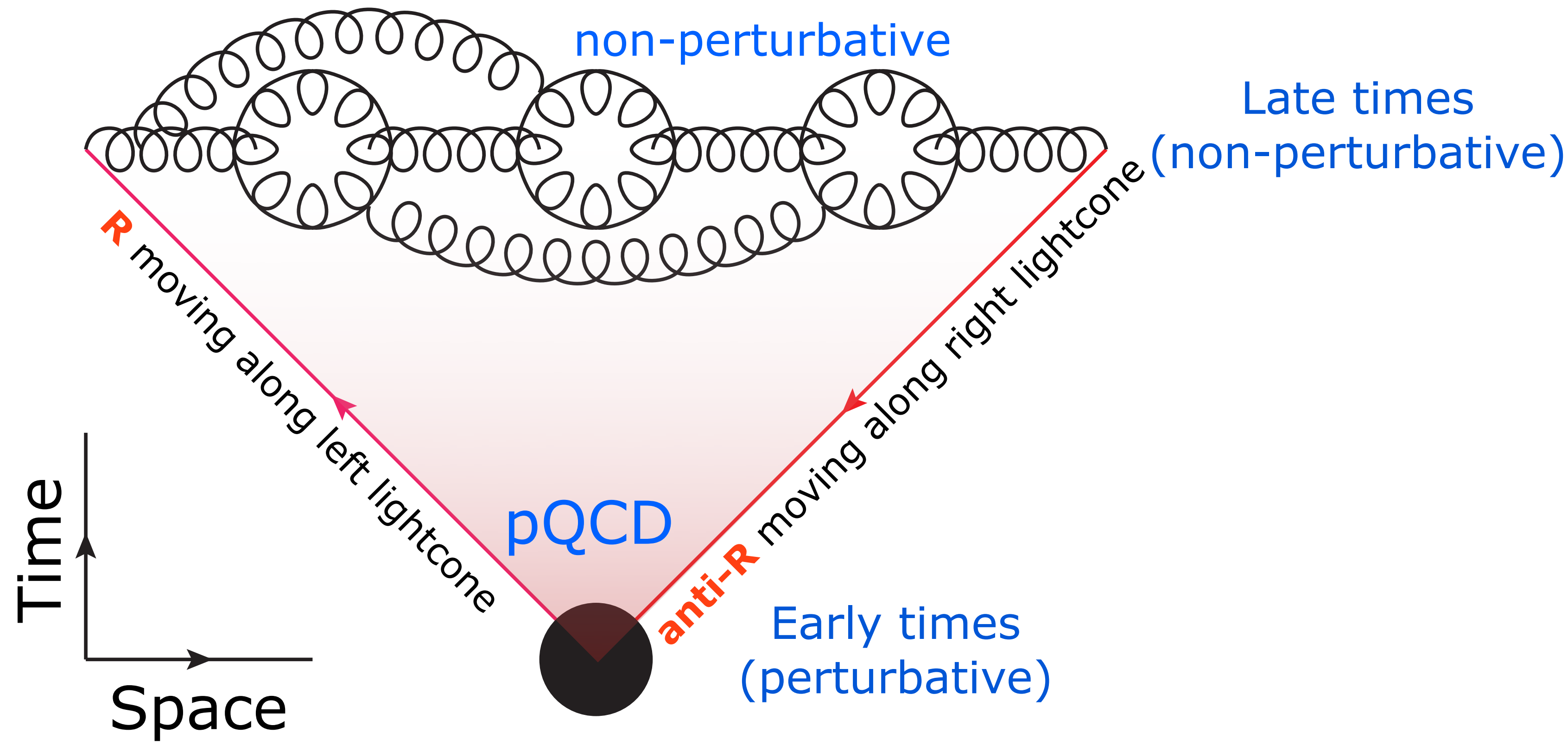
→ Too naive to see LPHD (inclusive) as a justification for Independent Fragmentation (exclusive)

→ More physics needed

Colour Neutralization

A physical hadronization model

Should involve at least TWO partons, with opposite color charges (e.g., **R** and **anti-R**)



Strong "confining" field emerges between the two charges when their separation > ~ 1fm

Color Flow

Between which partons do confining potentials arise?

Set of simple rules for color flow, based on large-N_C limit

(Never Twice Same Color: true up to $O(1/N_C^2)$)

Illustrations from: P.Nason & P.S.,
PDG Review on *MC Event Generators*, 2012

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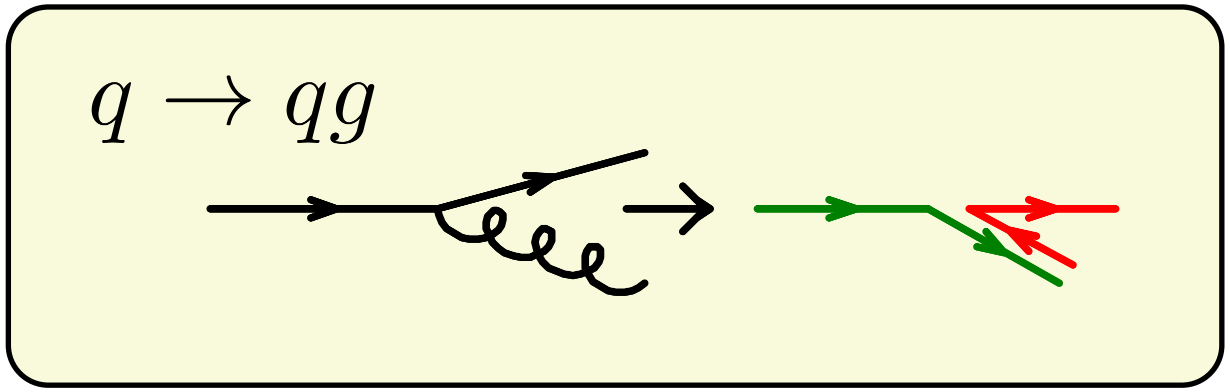

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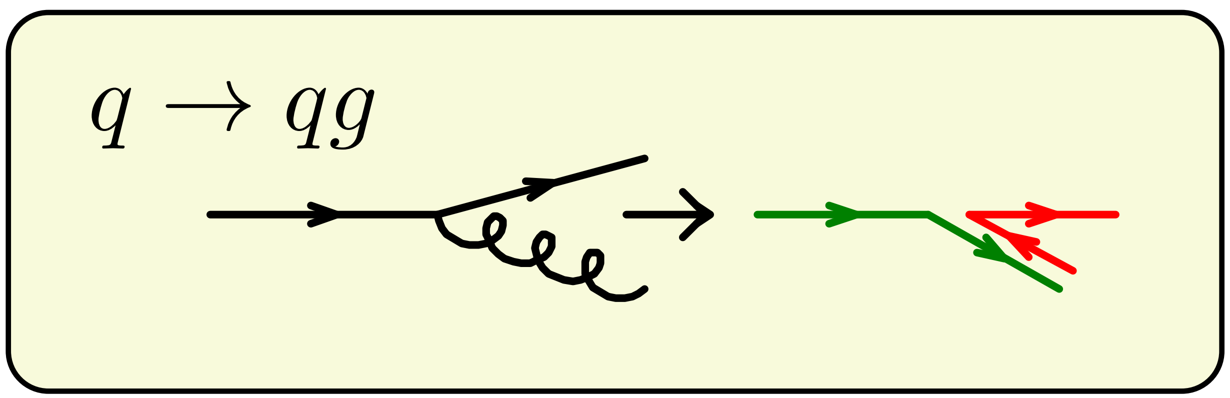

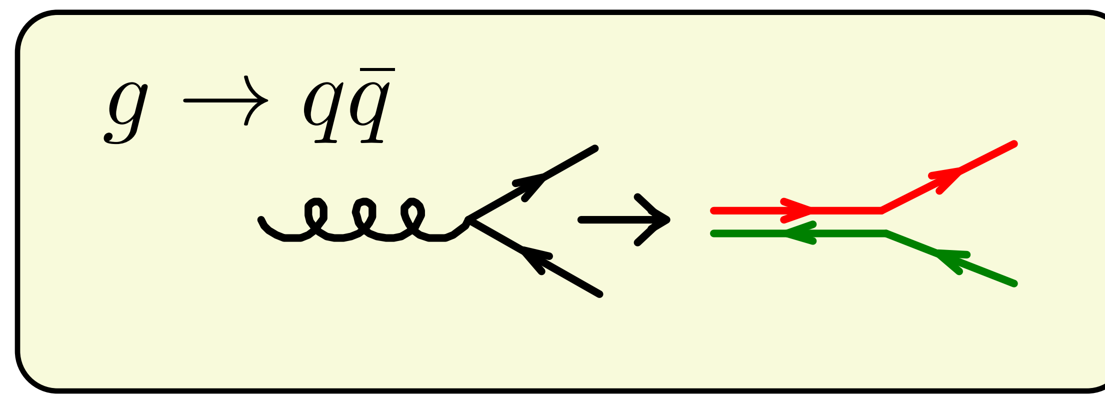

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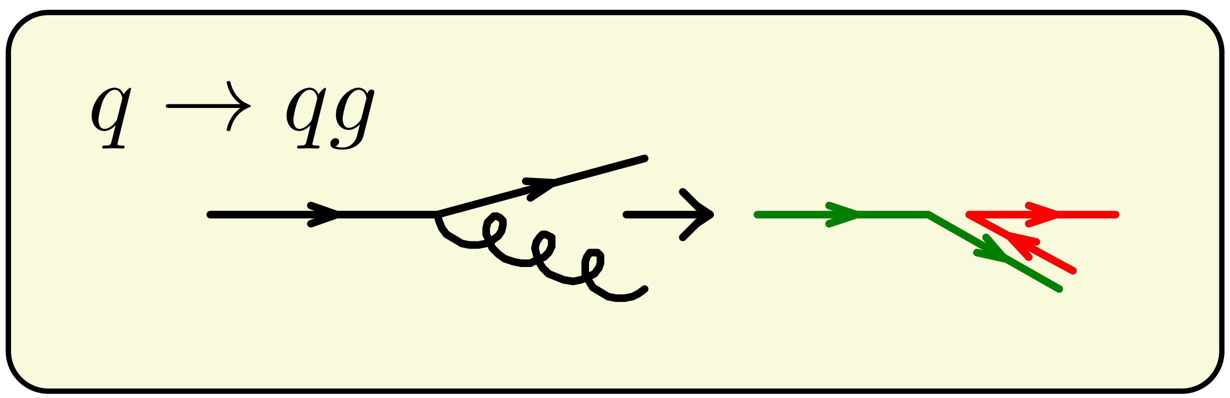

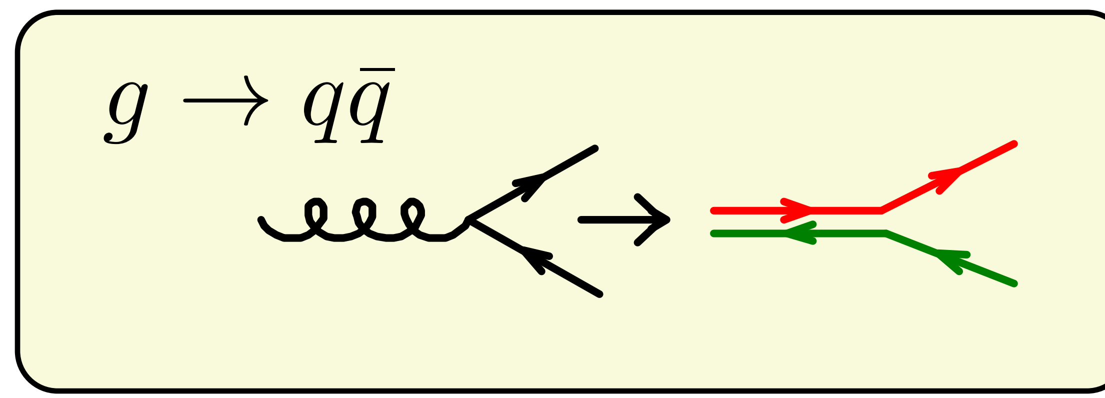

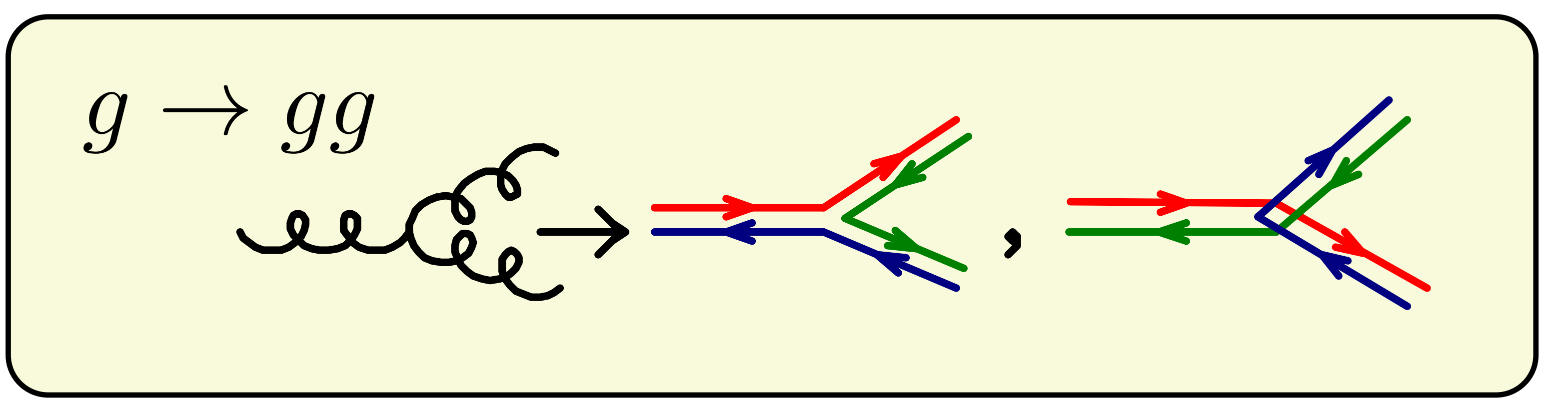

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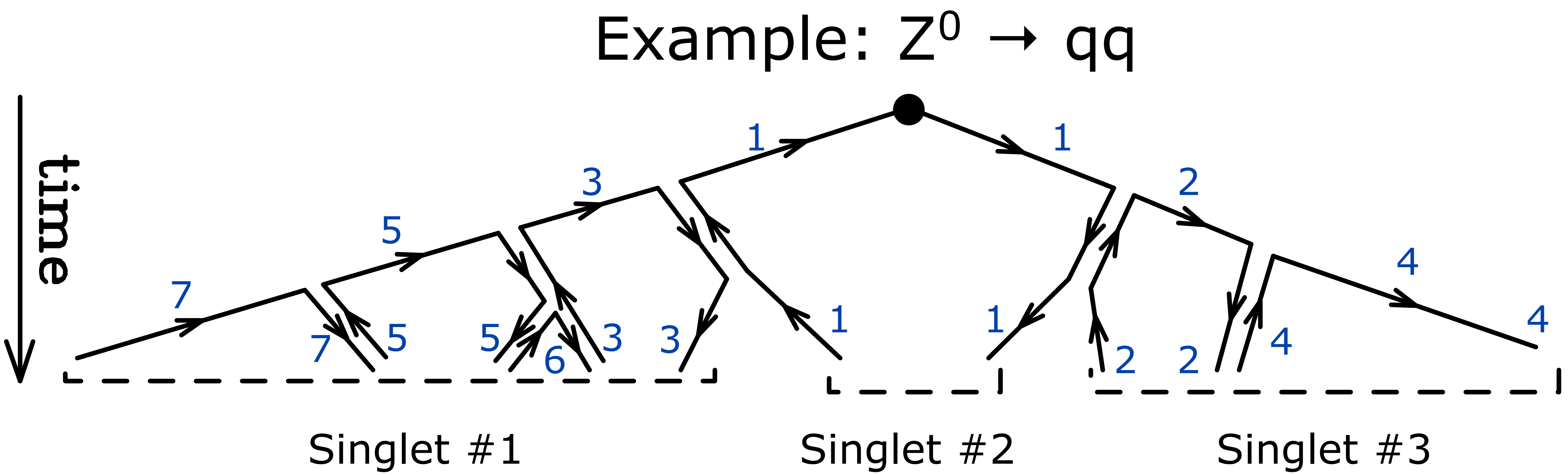
Color Flow

For an entire Cascade



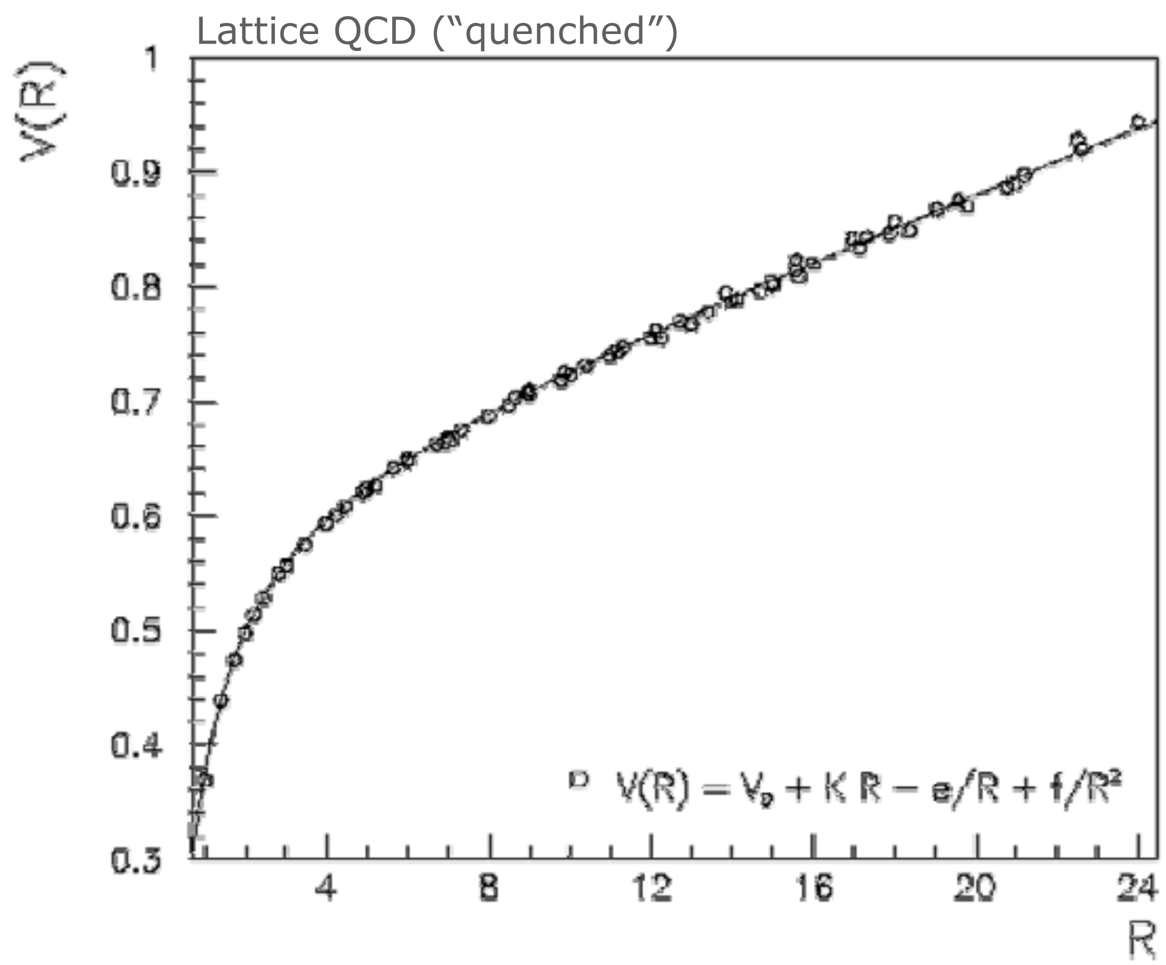
Coherence of pQCD cascades → not much "overlap" between singlet subsystems
→ Leading-colour approximation pretty good

LEP measurements in WW confirm this (at least to order 10% ~ $1/N_c^2$)

Note: (much) more color getting kicked around in hadron collisions → more later

Confinement

Potential between a quark and an antiquark as function of distance, R



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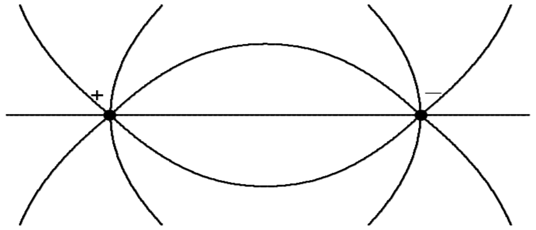
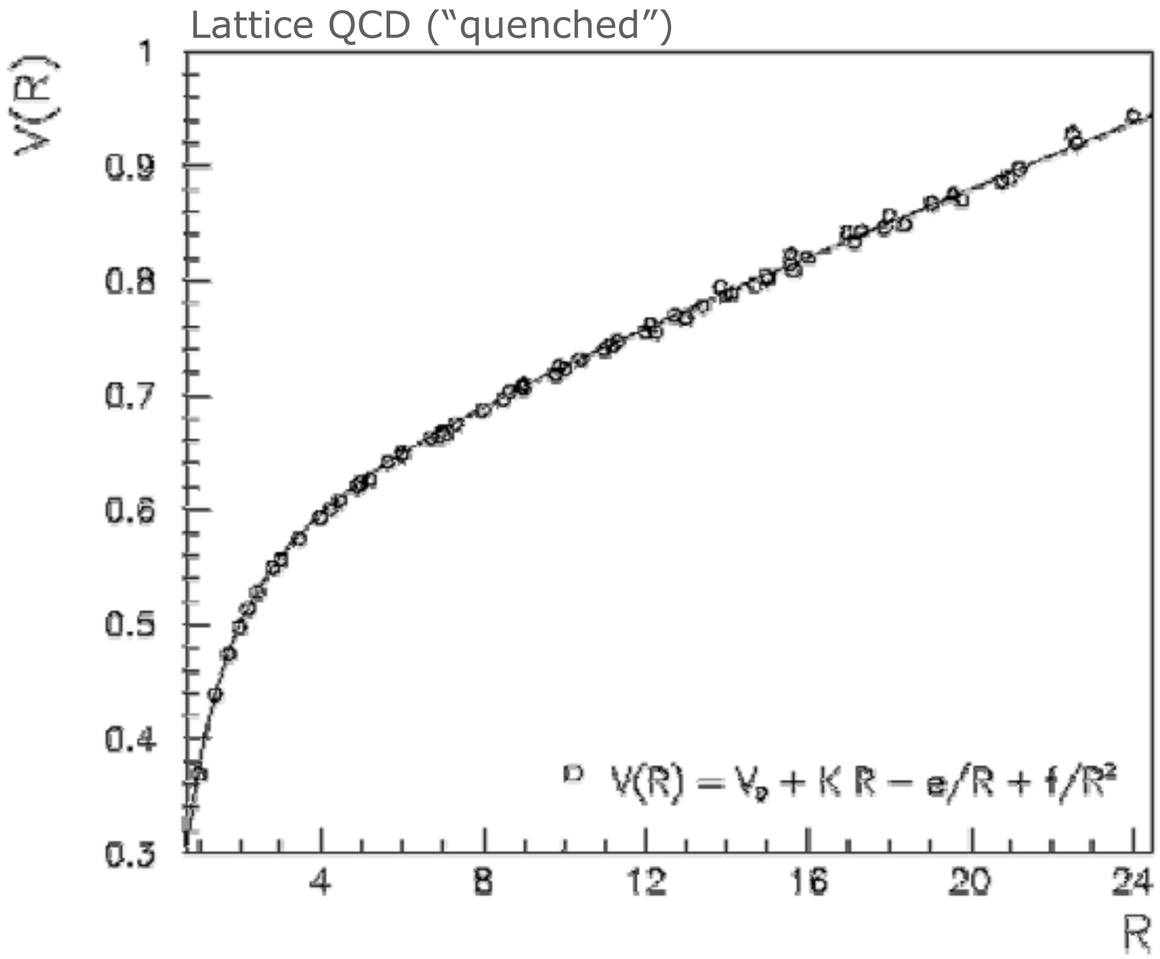
Potential between a quark and an antiquark as function of distance, R

Short Distances ~ "Coulomb"



Partons



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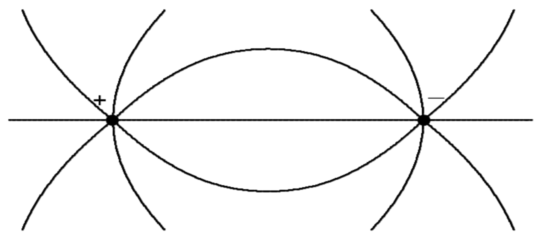
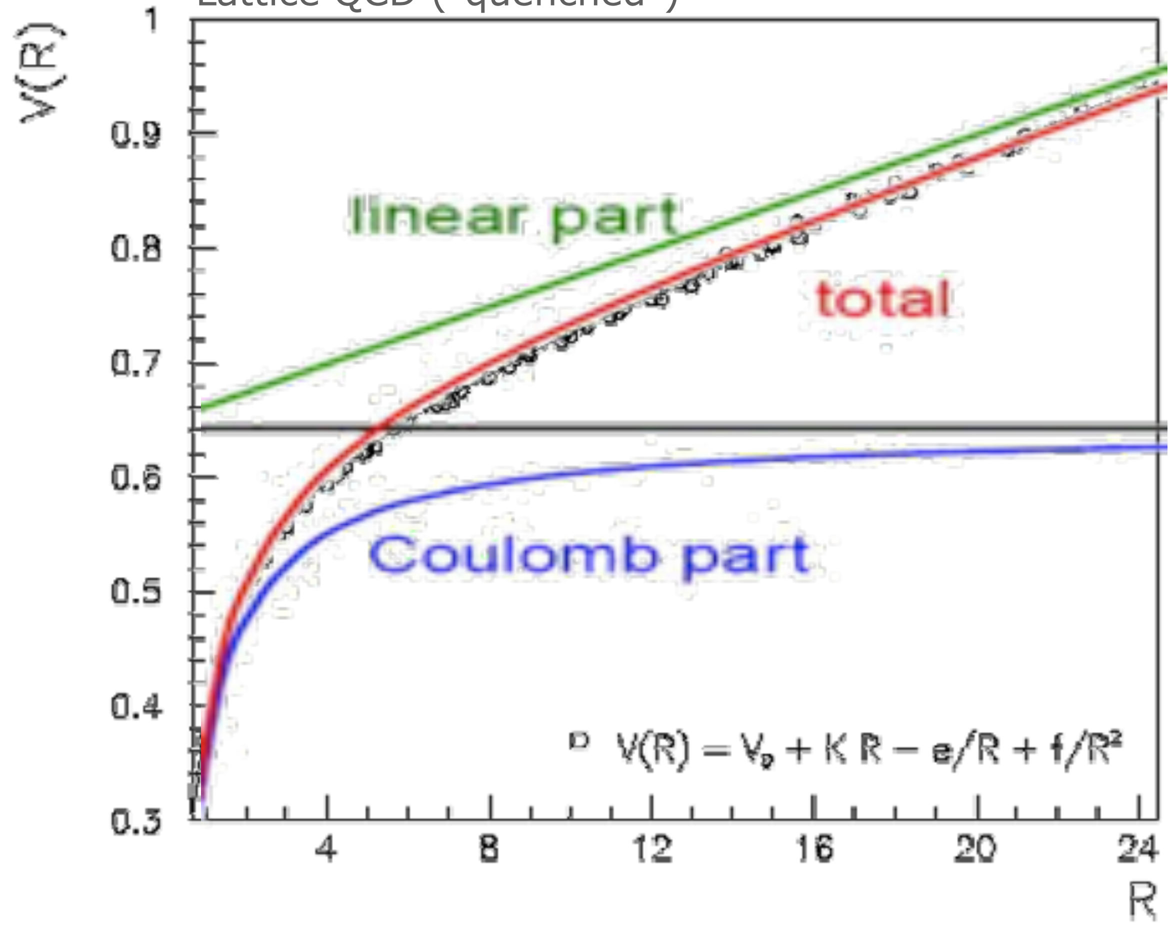


Partons

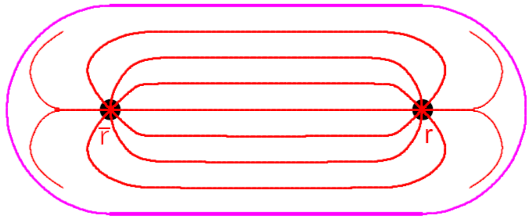
Lattice QCD ("quenched")



Confinement

Potential between a quark and an antiquark as function of distance, R

Long Distances ~ Linear Potential



Quarks (and gluons) confined inside hadrons

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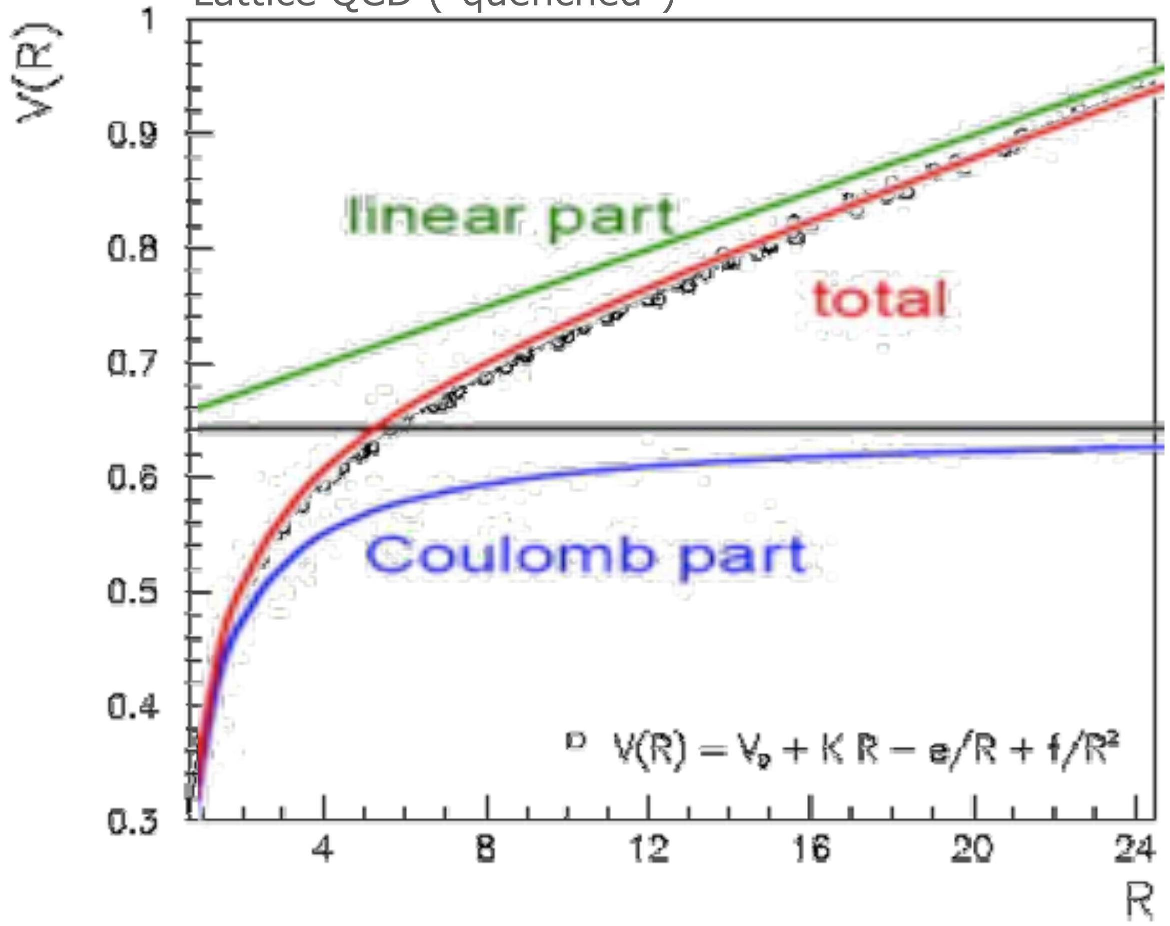

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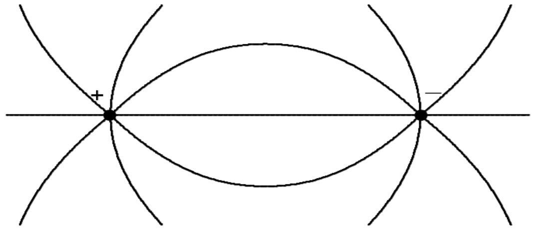

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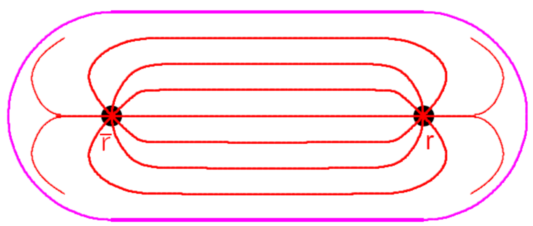

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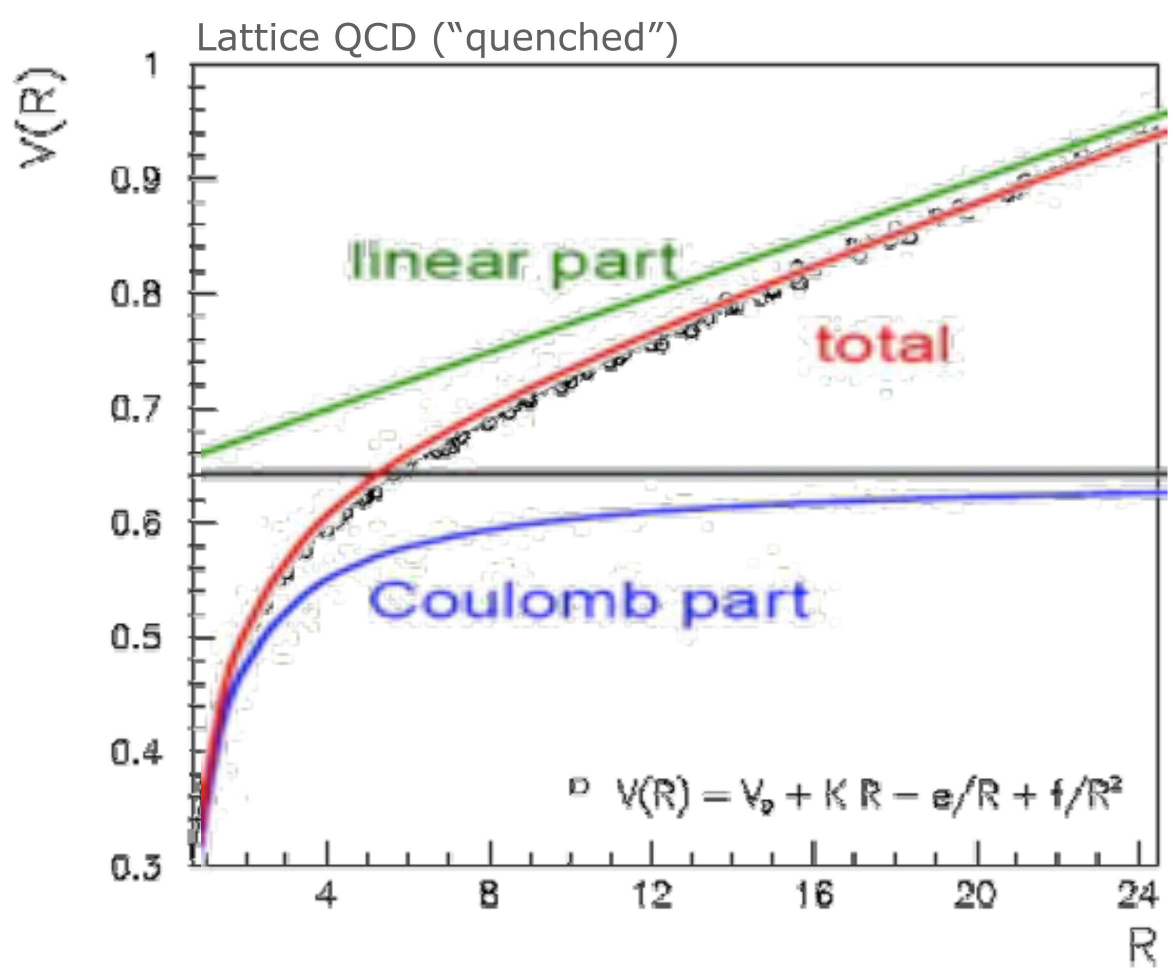

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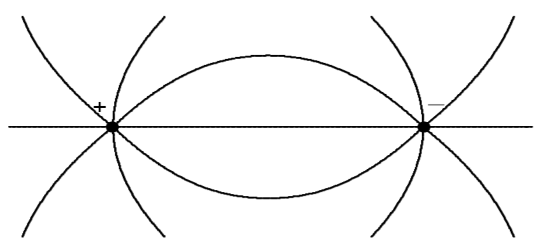

Partons

$$F(r) \approx \text{const} = \kappa \approx 1\ \text{GeV/fm} \quad \Longleftrightarrow \quad V(r) \approx \kappa r$$

~ Force required to lift a 16-ton truck

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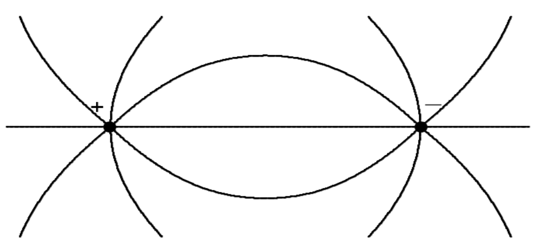

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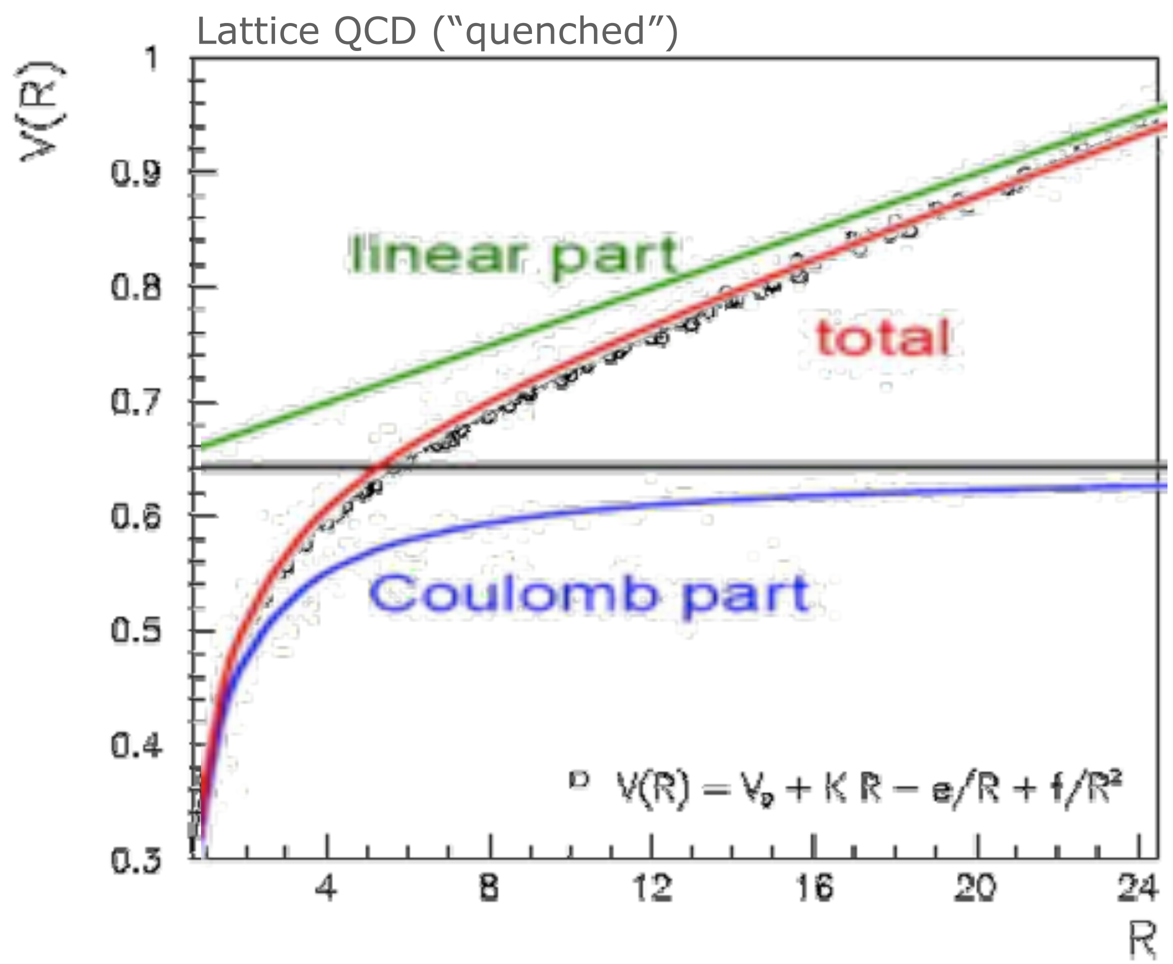

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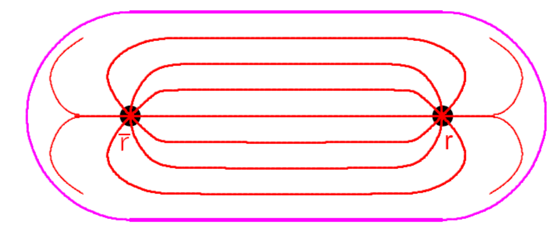


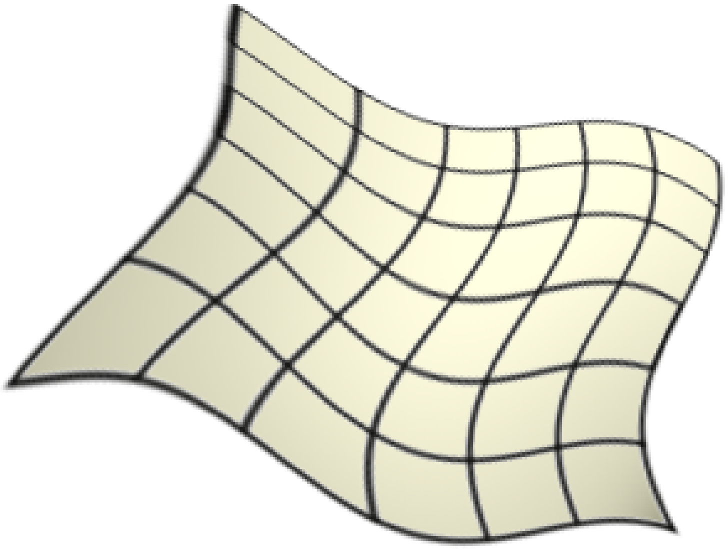
Quarks (and gluons) confined inside hadrons

What physical system has a linear potential?

$$F(r) \approx \text{const} = \kappa \approx 1\ \text{GeV/fm} \quad \Longleftrightarrow \quad V(r) \approx \kappa r$$

~ Force required to lift a 16-ton truck

From Partons to Strings



Motivates a model:

Let color field collapse into a (infinitely) narrow flux tube of uniform energy density $\kappa \sim 1$ GeV / fm

→ Relativistic 1+1 dimensional worldsheet – string

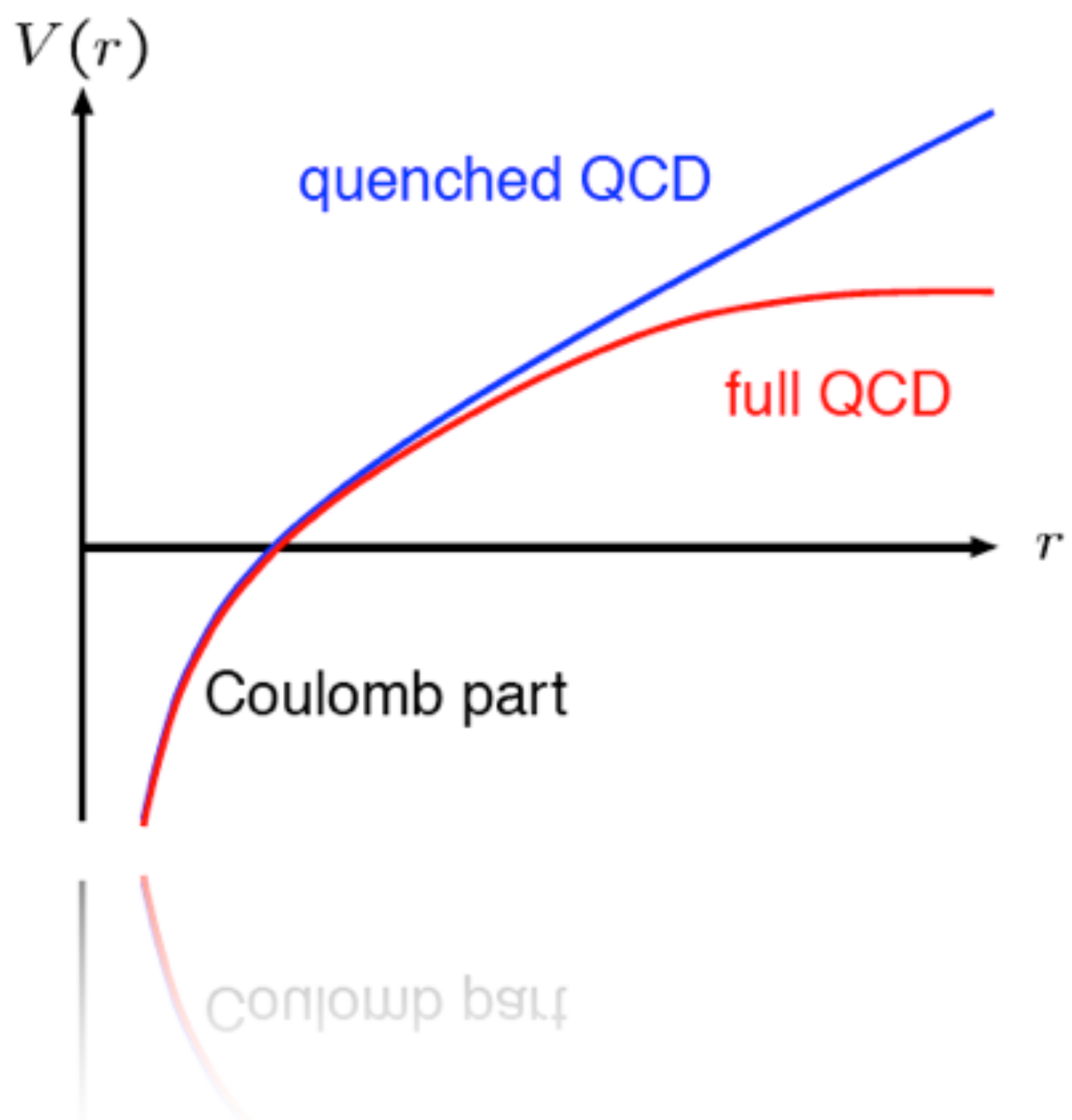
Pedagogical Review: B. Andersson, *The Lund model.* Camb. Monogr. Part. Phys. Nucl. Phys. Cosmol., 1997.

String Breaks

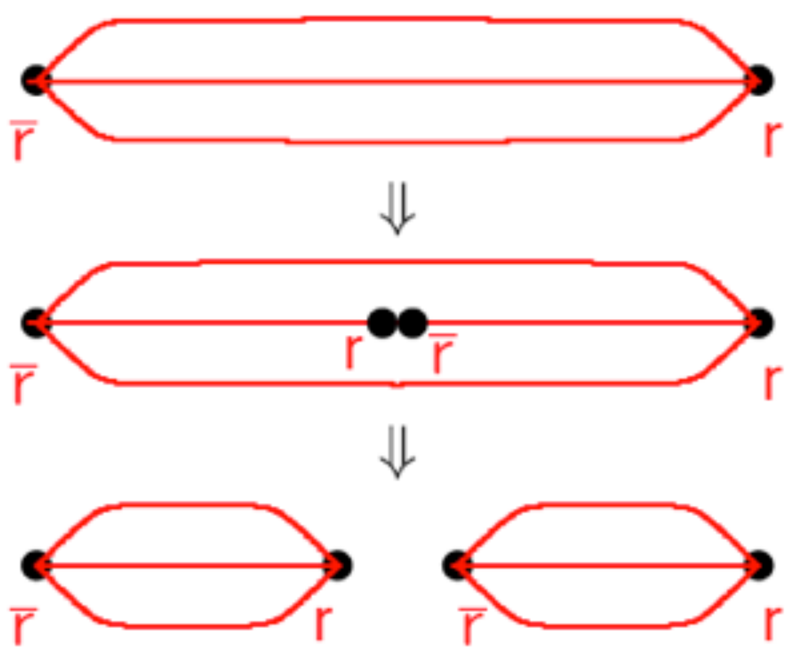
String Breaks

In “unquenched” QCD

g→qq → The strings would break



String Breaks:
via Quantum Tunneling



(simplified colour representation)

$$\mathcal{P} \propto \exp\left(\frac{-m_q^2 - p_\perp^2}{\kappa/\pi}\right)$$

Illustrations by T. Sjöstrand

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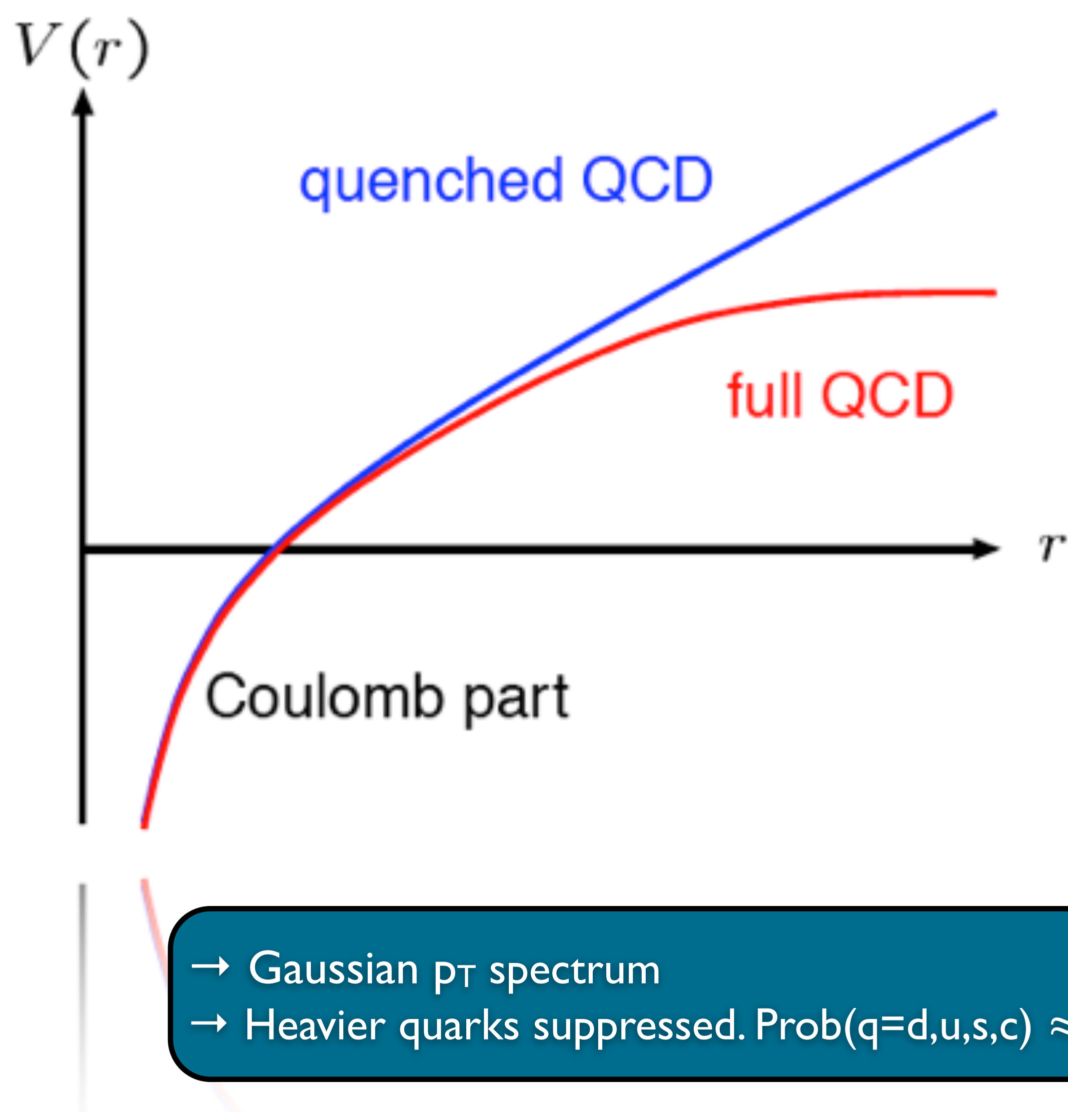

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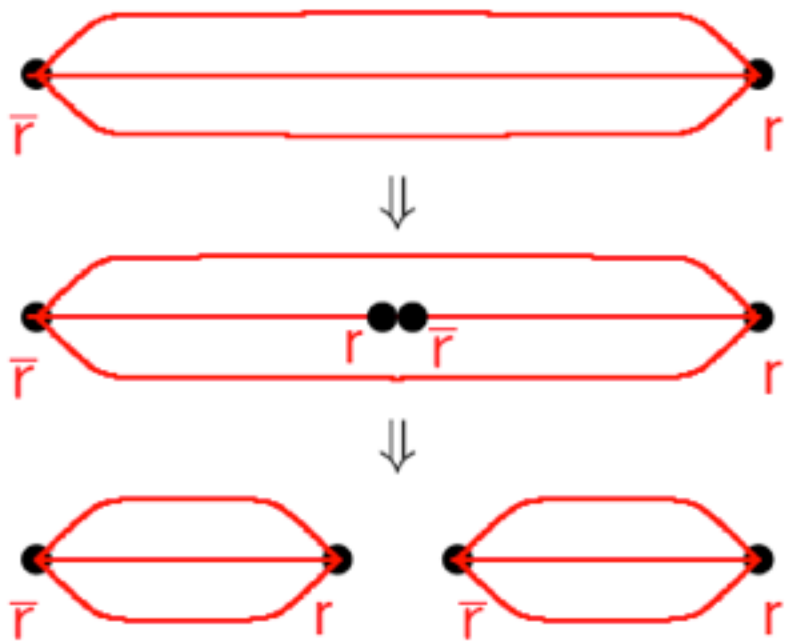

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→ Gaussian p_T spectrum
→ Heavier quarks suppressed. Prob(q=d,u,s,c) ≈ 1 : 1 : 0.2 : 10^{-11}

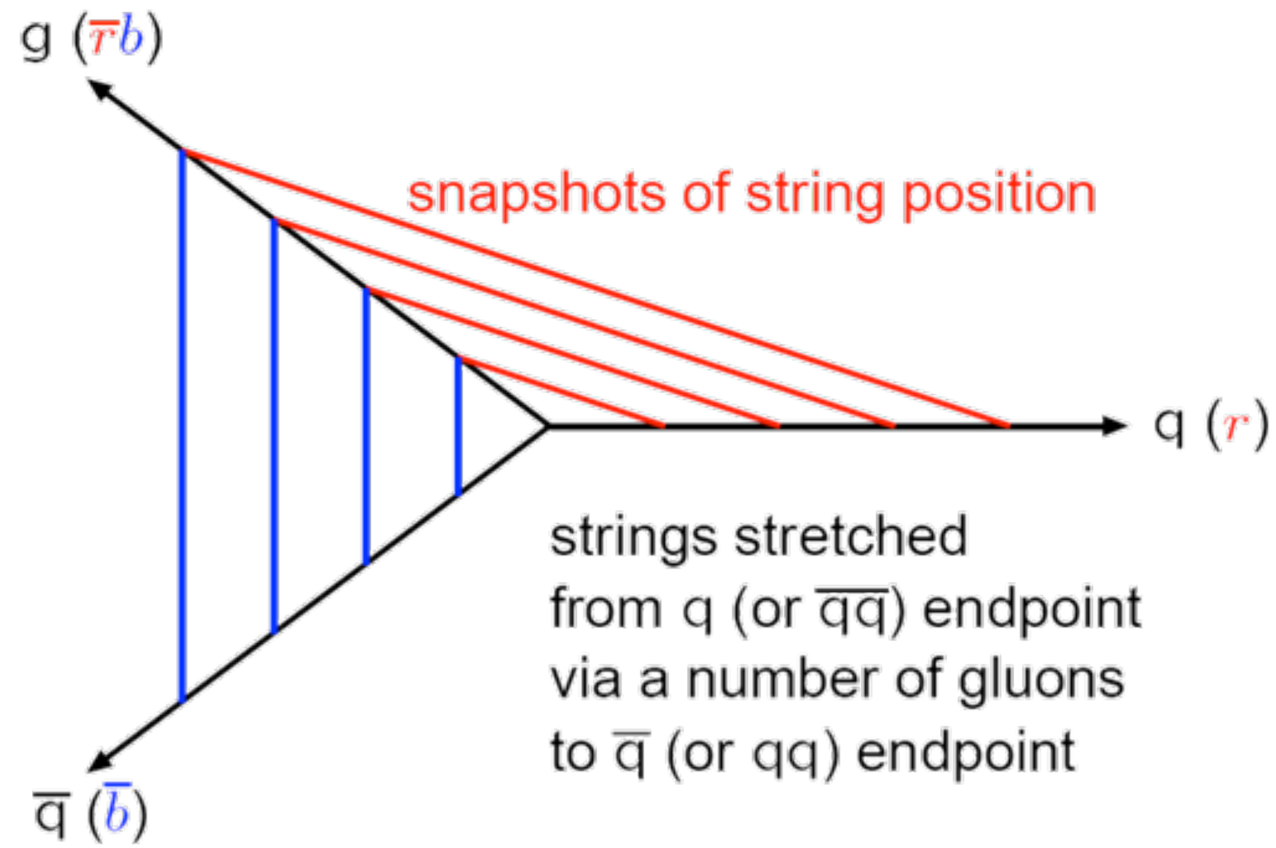
Illustrations by T. Sjöstrand

The (Lund) String Model

Map:

- **Quarks** → String Endpoints
- **Gluons** → Transverse Excitations (kinks)
- Physics then in terms of string worldsheet evolving in spacetime
- Probability of string break (by quantum tunneling) constant per unit area → **AREA LAW**

See also Yuri's 2nd lecture



Gluon = kink on string, carrying energy and momentum

→ **STRING EFFECT**

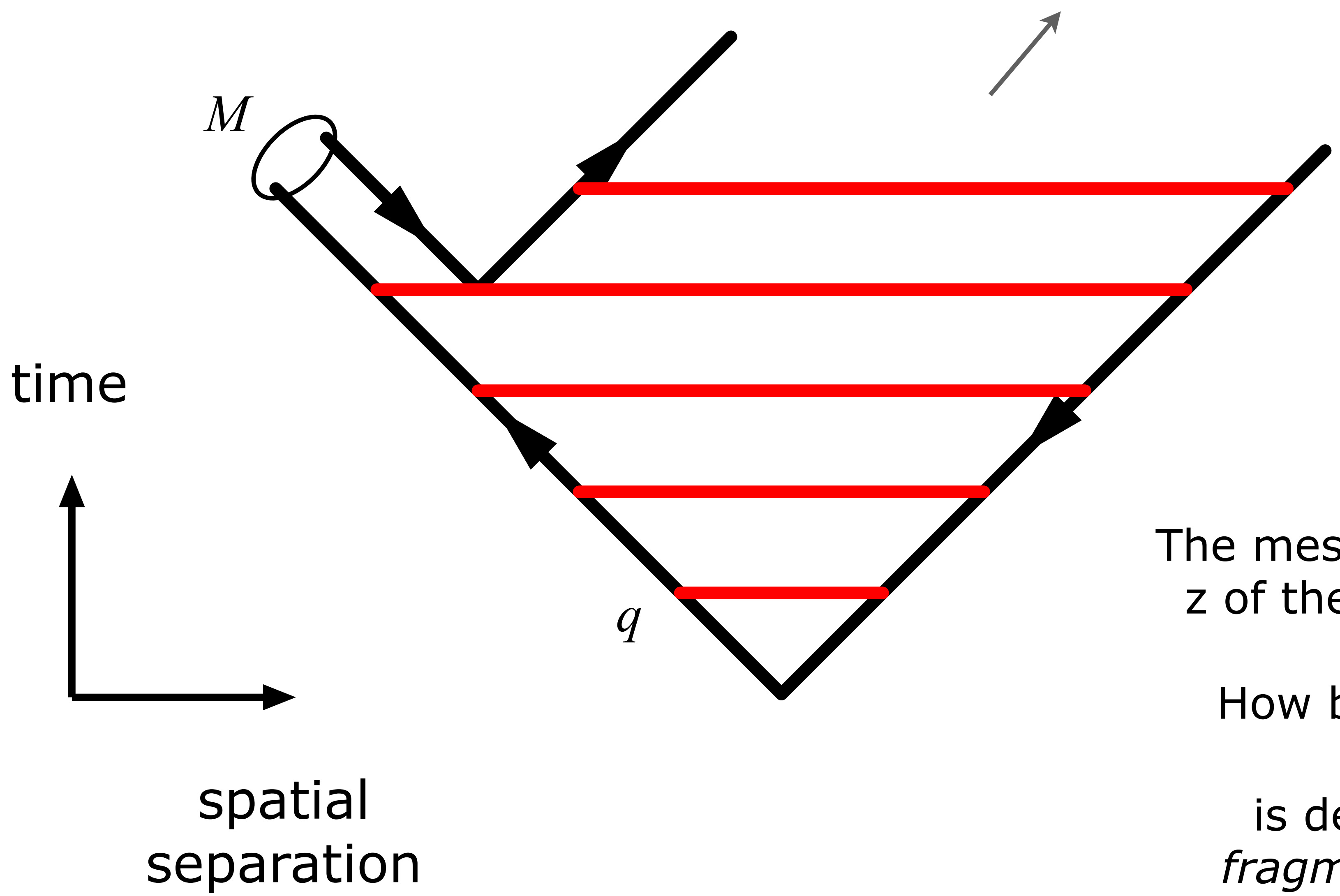
Simple space-time picture

Details of string breaks more complicated (e.g., baryons, spin multiplets)

Fragmentation Function

Spacetime Picture

leftover string,
further string breaks

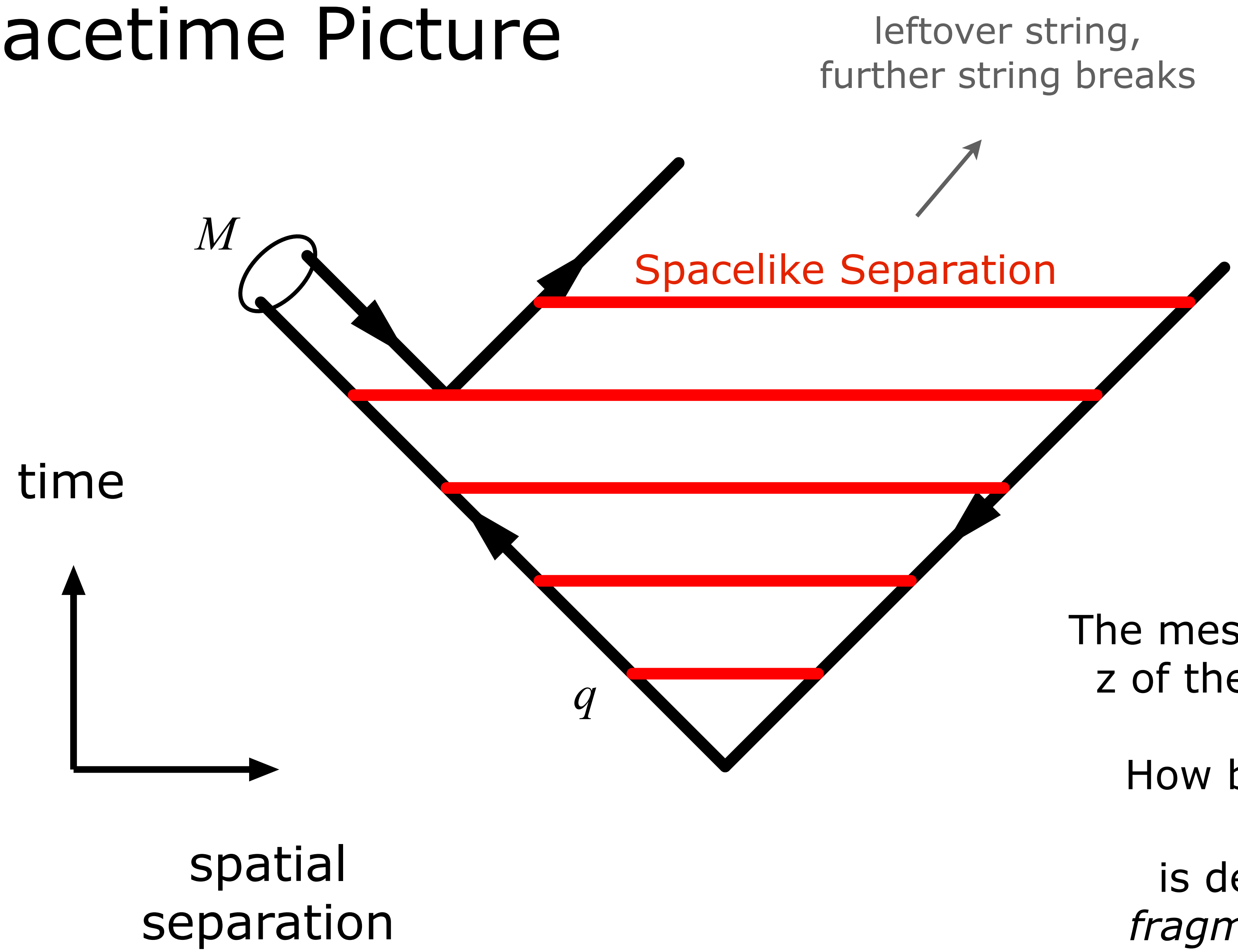


The meson M takes a fraction
z of the quark momentum,

How big that fraction is,
$z \in [0,1]$,
is determined by the
fragmentation function,
$f(z,Q_0^2)$

Fragmentation Function

Spacetime Picture



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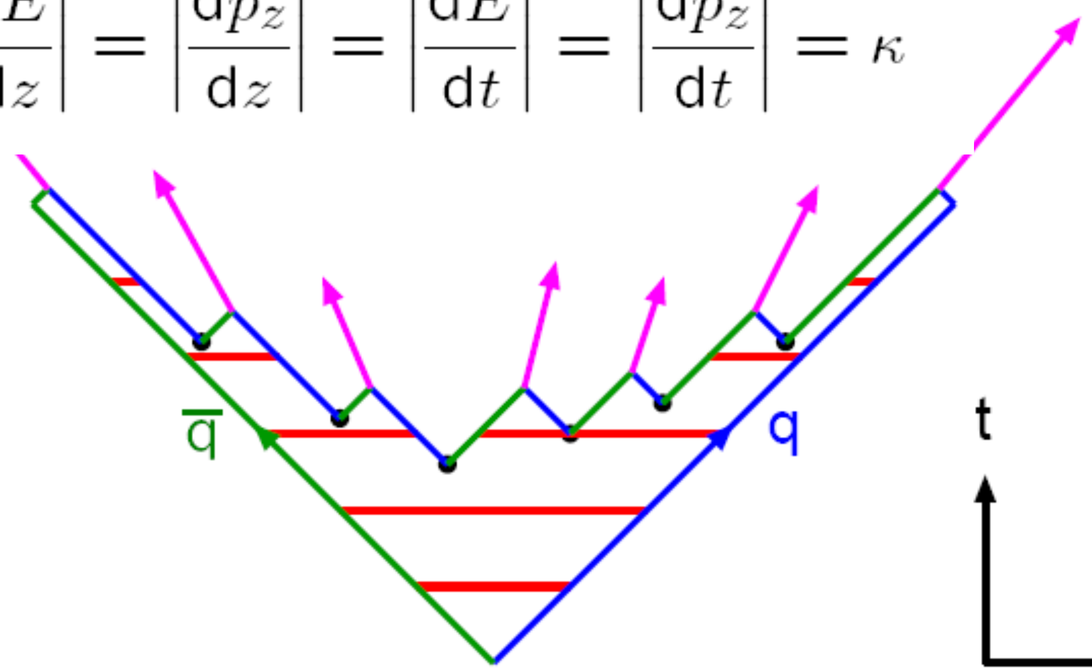
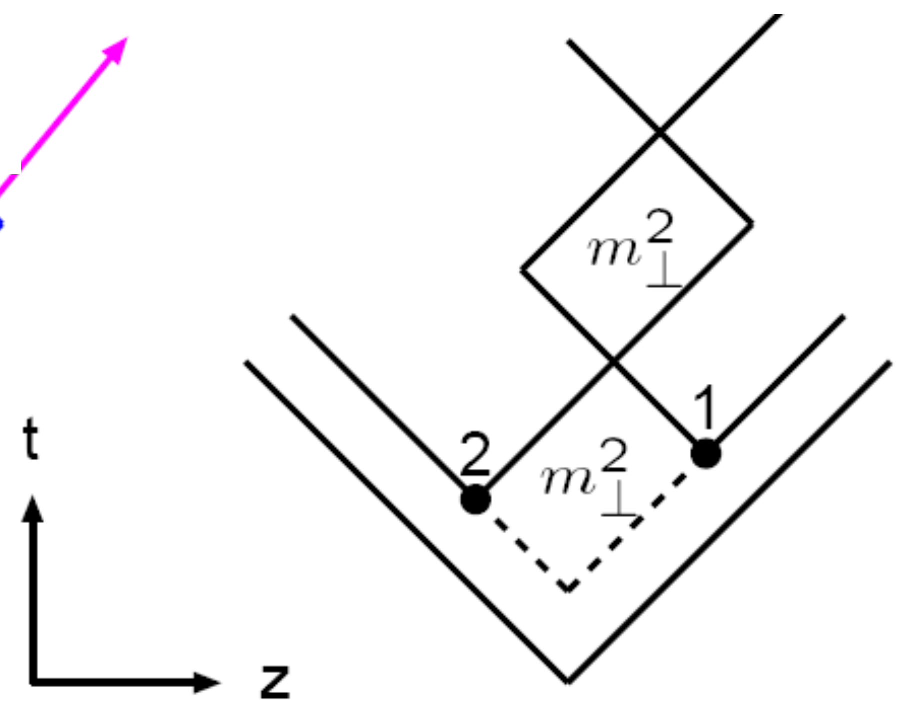
Large System

Illustrations by T. Sjöstrand

$$\left|\frac{\mathrm{d}E}{\mathrm{d}z}\right| = \left|\frac{\mathrm{d}p_z}{\mathrm{d}z}\right| = \left|\frac{\mathrm{d}E}{\mathrm{d}t}\right| = \left|\frac{\mathrm{d}p_z}{\mathrm{d}t}\right| = \kappa$$





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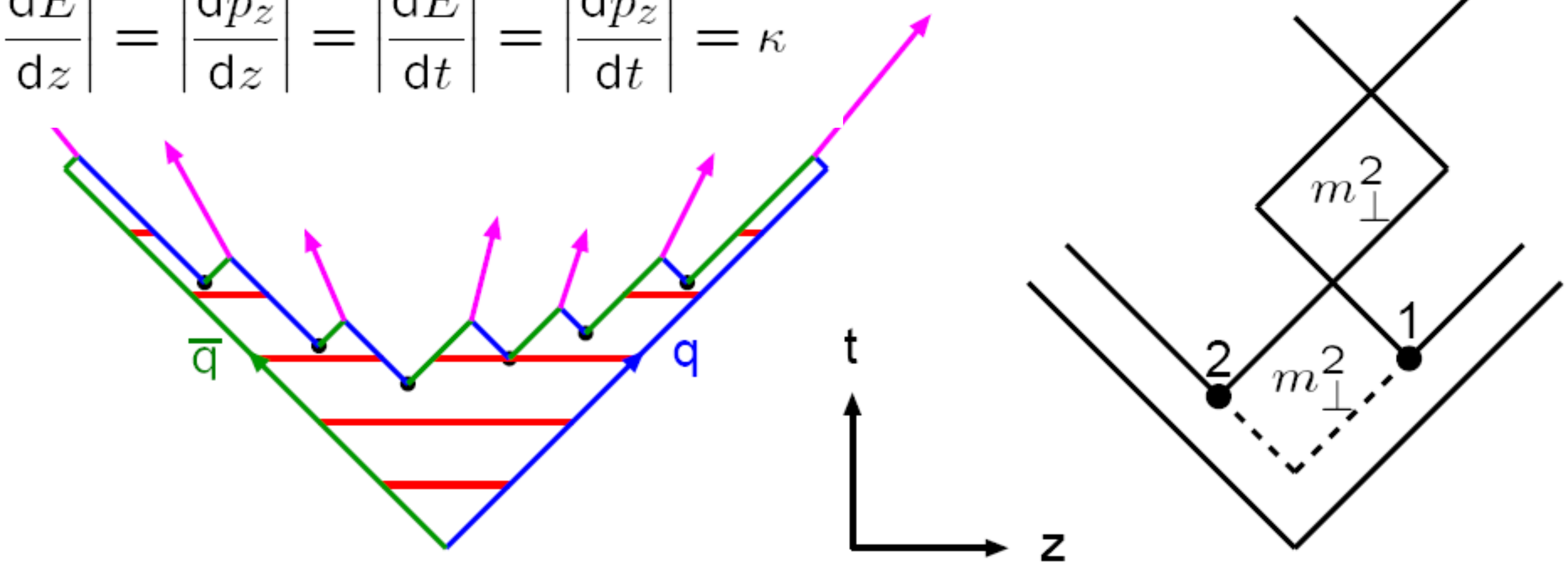


String breaks causally disconnected

→ can proceed in arbitrary order (left-right, right-left, in-out, …)

→ constrains possible form of fragmentation function

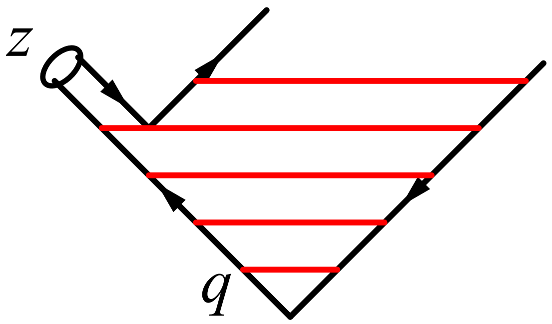
→ Justifies iterative ansatz (useful for MC implementation)

Left-Right Symmetry

Causality → Left-Right Symmetry

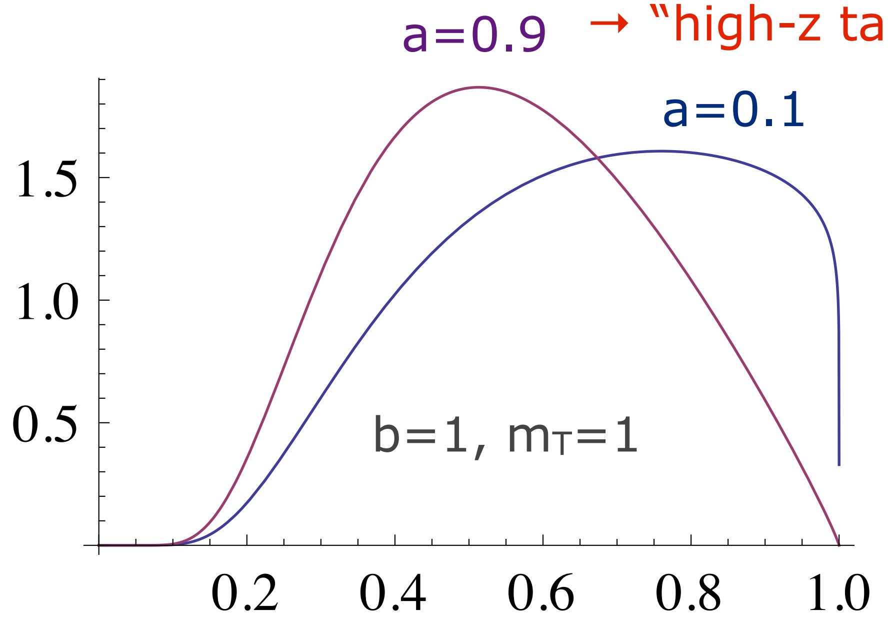
→ Constrains form of fragmentation function!

→ Lund Symmetric Fragmentation Function

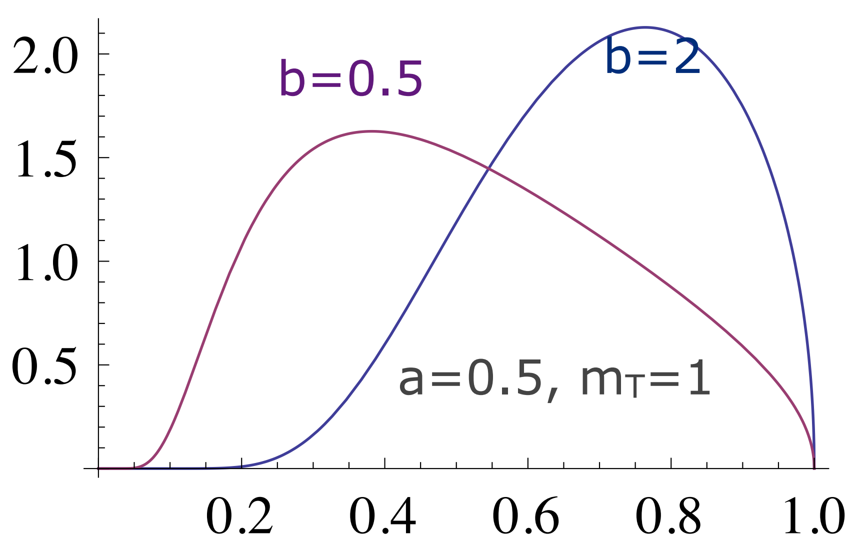


$$f(z) \propto \frac{1}{z}(1-z)^a \exp\left(-\frac{b\,(m_h^2 + p_{\perp h}^2)}{z}\right)$$

Small a
→ "high-z tail"



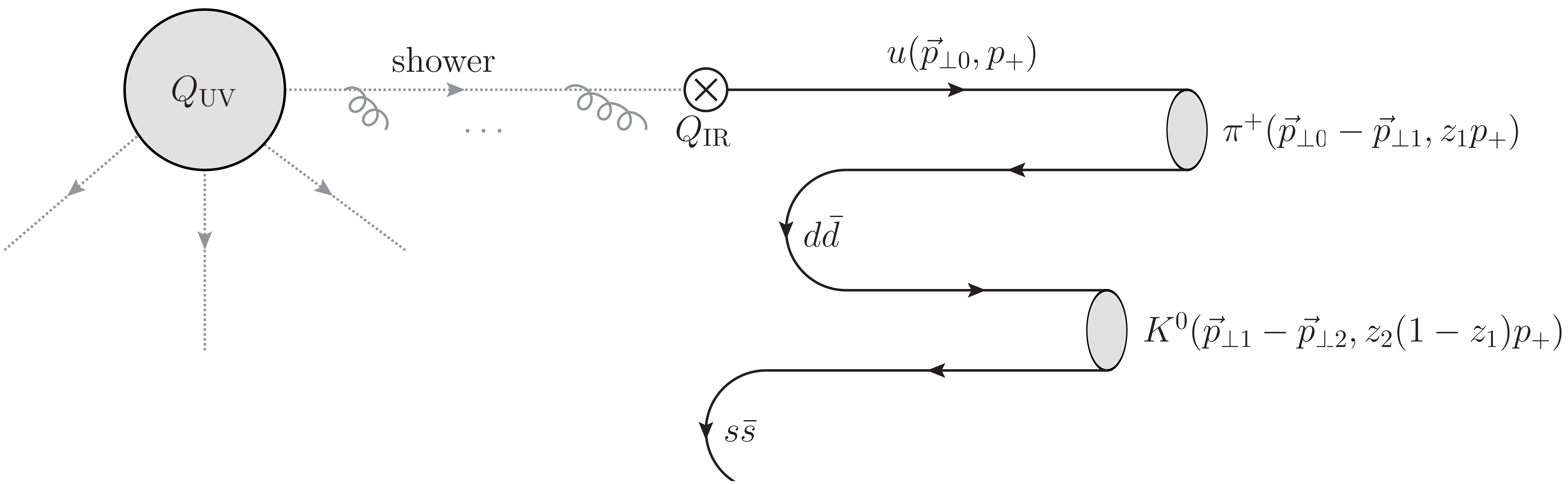
Small b
→ "low-z enhancement"



Note: In principle, a can be flavour-dependent. In practice, we only distinguish between baryons and mesons

Iterative String Breaks

Causality → May iterate from outside-in



The Length of Strings

In Space:

String tension ≈ 1 GeV/fm → a 5-GeV quark can travel 5 fm before all its kinetic energy is transformed to potential energy in the string.

Then it must start moving the other way. String breaks will have happened behind it → yo-yo model of mesons

In Rapidity :

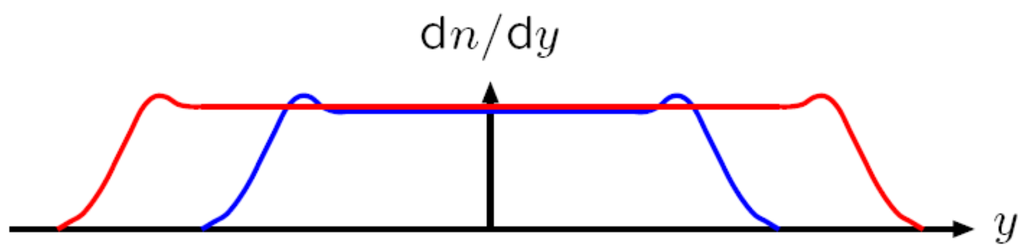
$$y = \frac{1}{2}\ln\left(\frac{E+p_z}{E-p_z}\right) = \frac{1}{2}\ln\left(\frac{(E+p_z)^2}{E^2-p_z^2}\right)$$

For a pion with z=1 along string direction (For beam remnants, use a proton mass):

$$y_{\mathrm{max}} \sim \ln\left(\frac{2E_q}{m_\pi}\right)$$

Note: Constant average hadron multiplicity per unit y → logarithmic growth of total multiplicity

Scaling in lightcone $p_\pm = E \pm p_z$ (for $q\bar{q}$ system along z axis) implies flat central rapidity plateau + some endpoint effects:



$\langle n_{\mathrm{ch}}\rangle \approx c_0 + c_1 \ln E_{\mathrm{cm}}$, ∼ Poissonian multiplicity distribution

Alternative: The Cluster Model

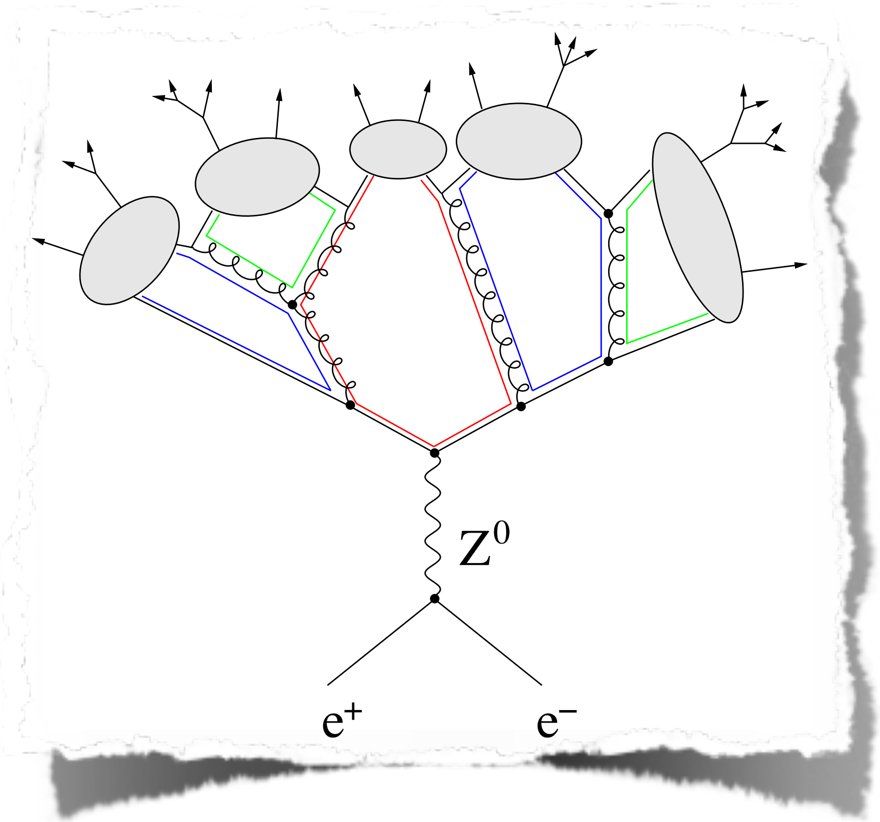
"Preconfinement"

+ Force g→qq splittings at Q_0

→ high-mass q-qbar "clusters"

Isotropic 2-body decays to hadrons

according to PS ≈ $(2s_1+1)(2s_2+1)(p^*/m)$



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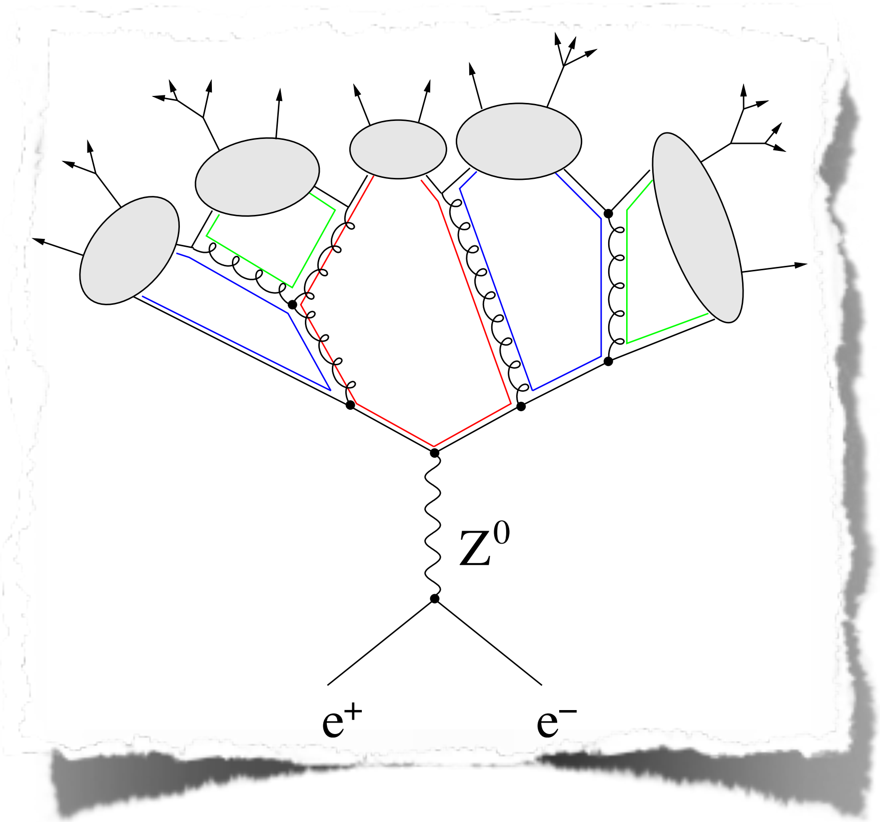

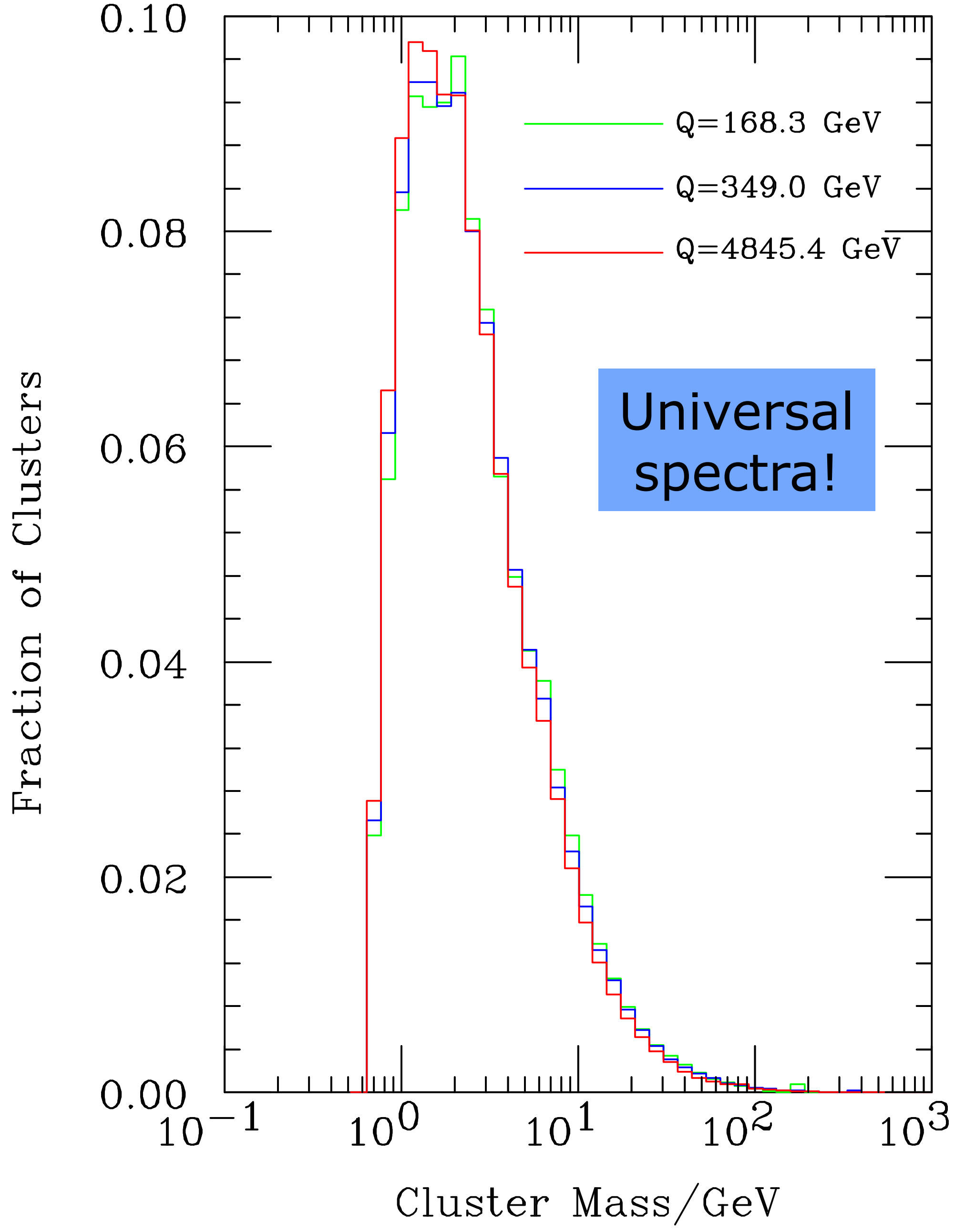

Universal spectra!

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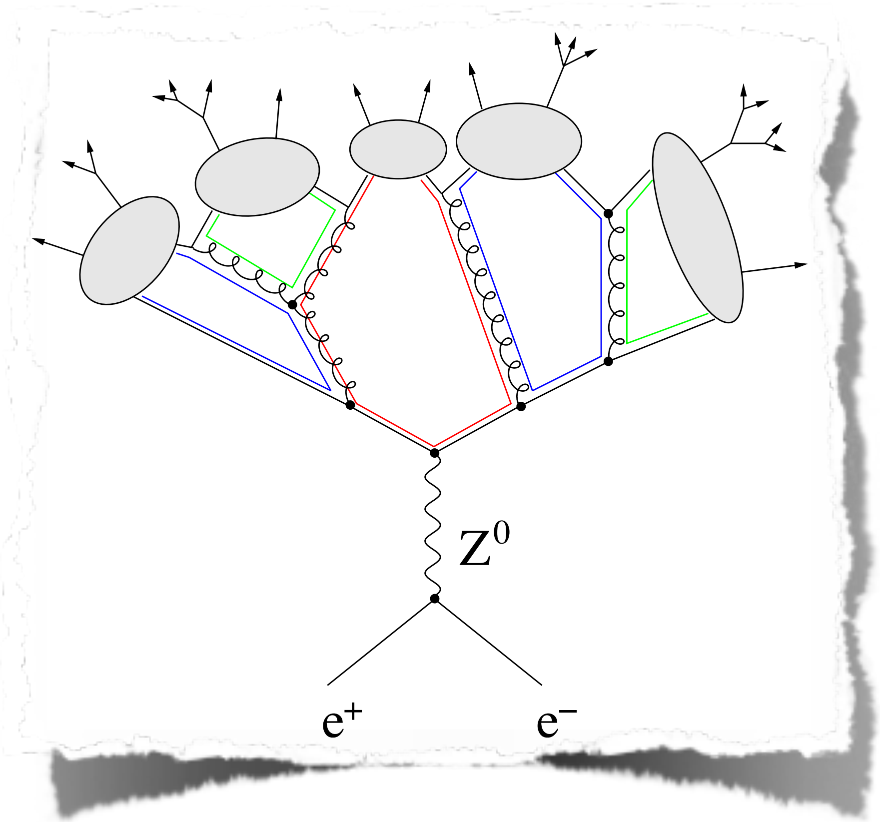
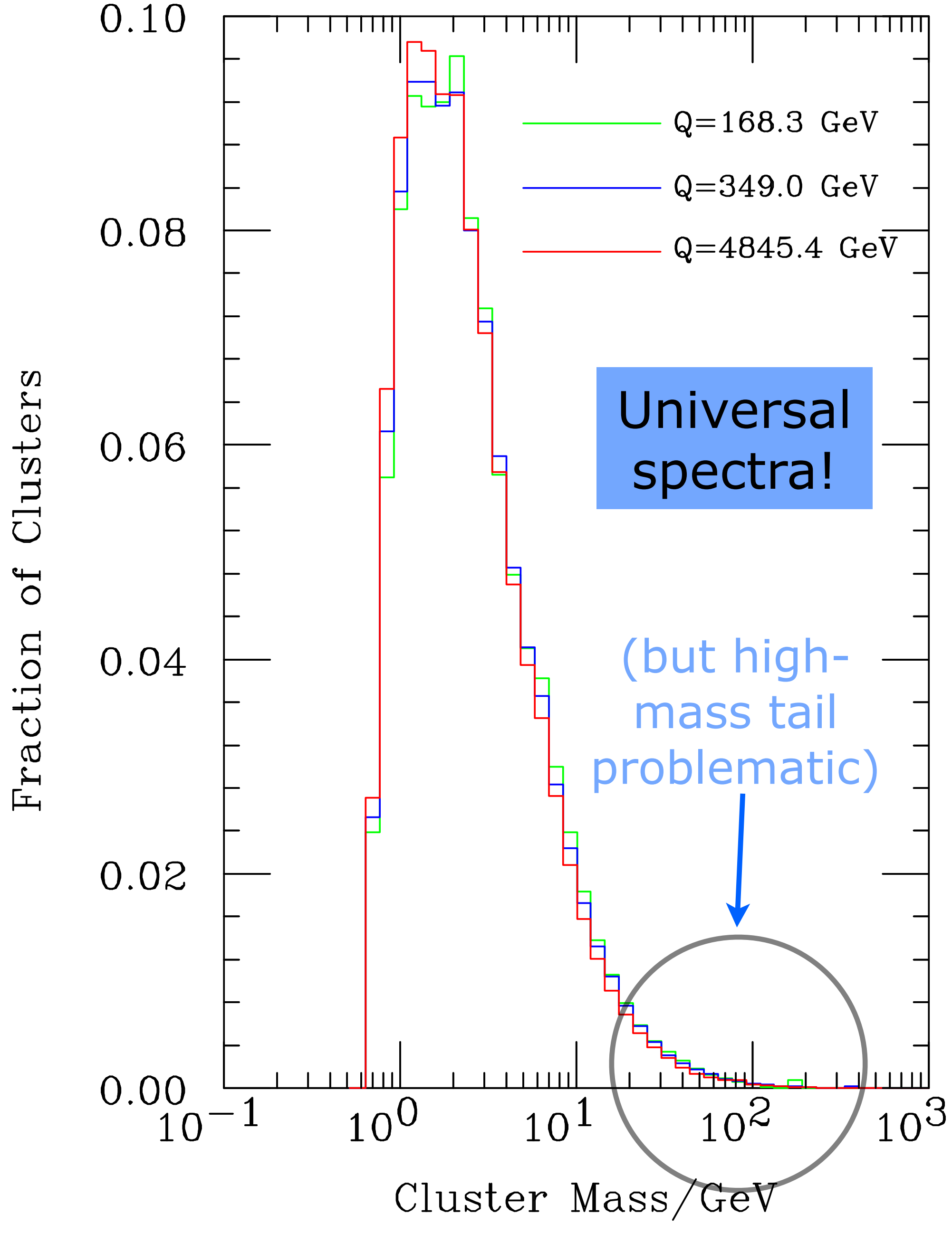
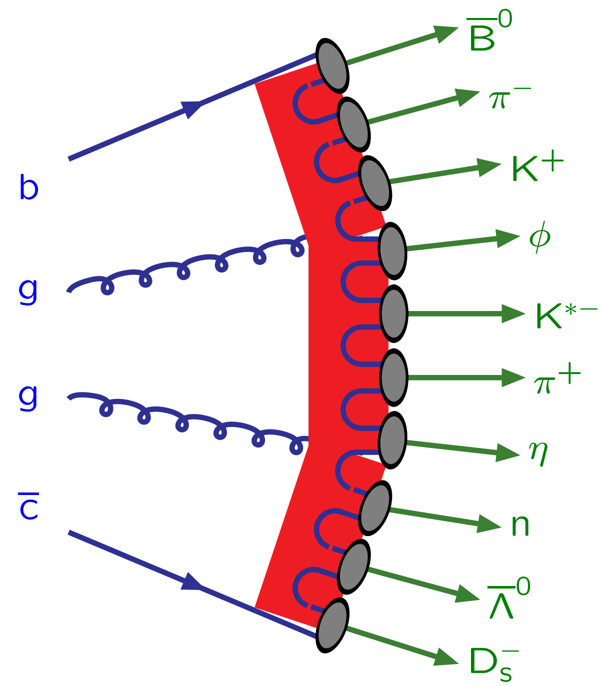


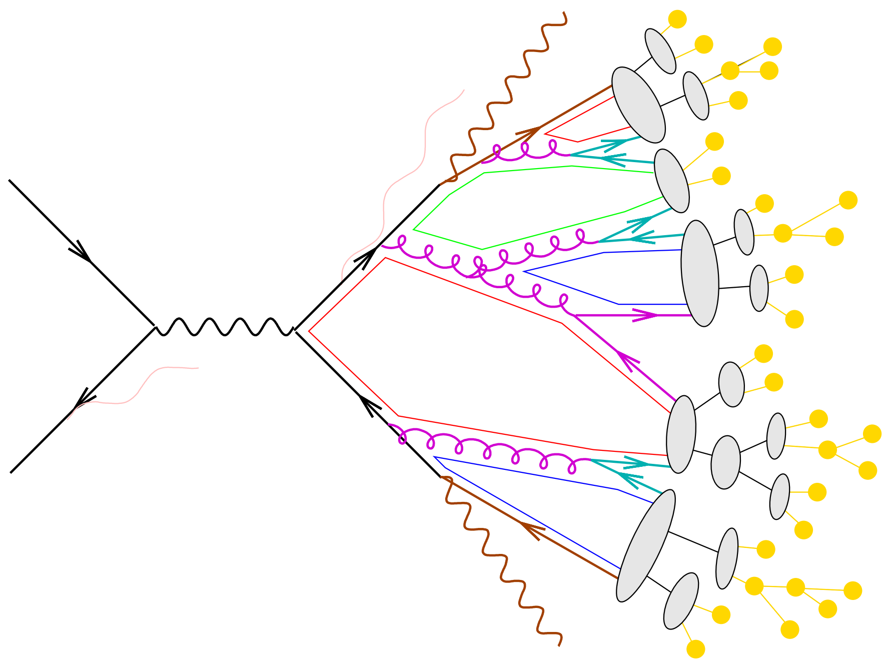


Universal spectra!

(but high-mass tail problematic)

Strings and Clusters





program	PYTHIA	HERWIG (&SHERPA)
model	string	cluster
energy–momentum picture	powerful	simple
	predictive	unpredictive
parameters	few	many
flavour composition	messy	simple
	unpredictive	in-between
parameters	many	few

Small strings → clusters. Large clusters → strings

Hadron Collisions

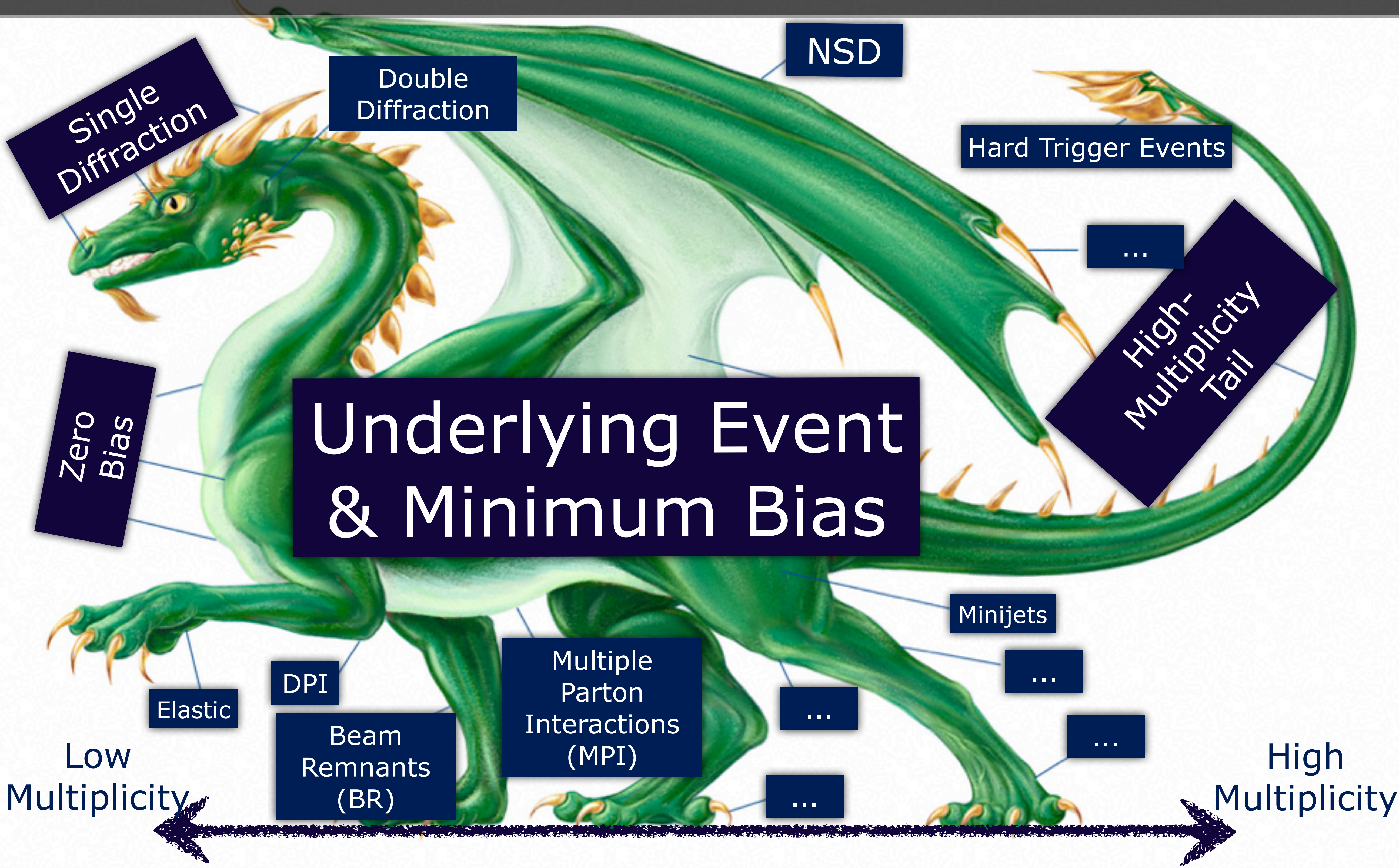


Image credits: E. Arenhaus & J. Walker

Hadron Collisions

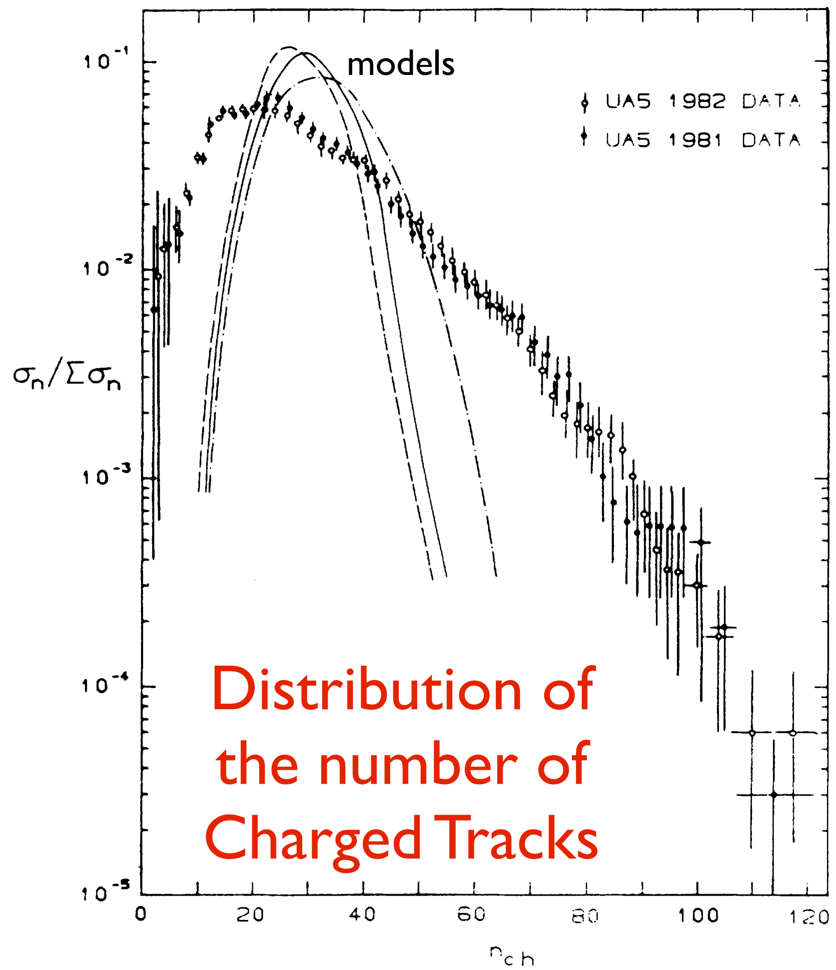
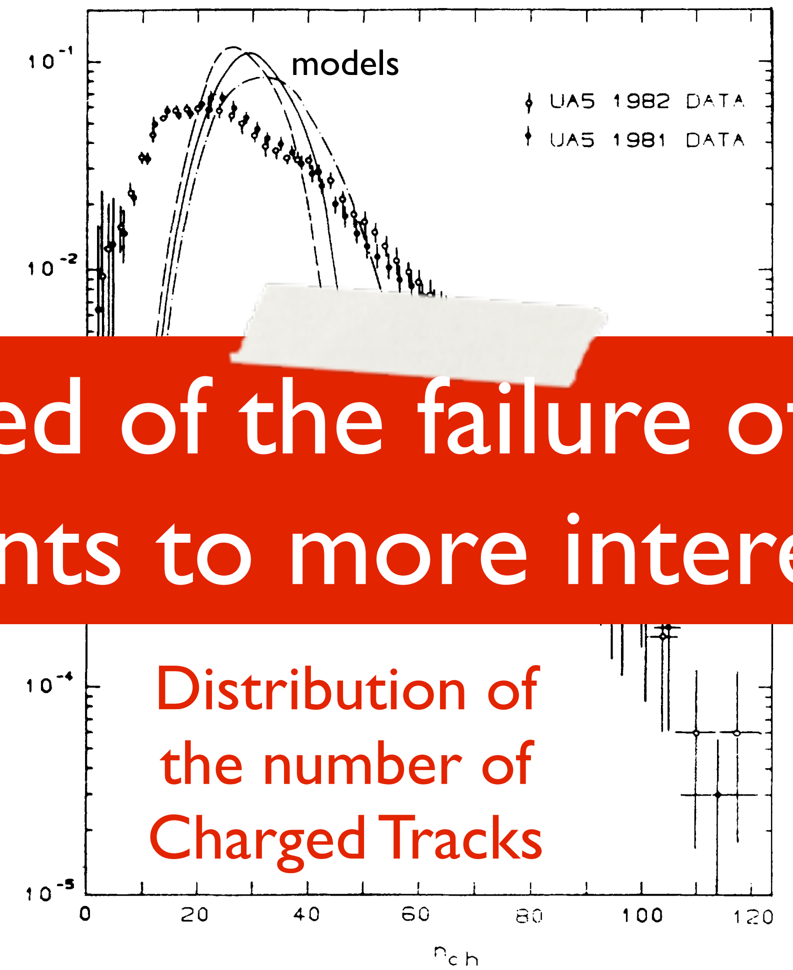


FIG. 3. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs simple models: dashed low p_T only, full including hard scatterings, dash-dotted also including initial- and final-state radiation.

Sjöstrand & v. Zijl, Phys.Rev.D36(1987)2019

Hadron Collisions

Do not be scared of the failure of physical models (typically points to more interesting physics)



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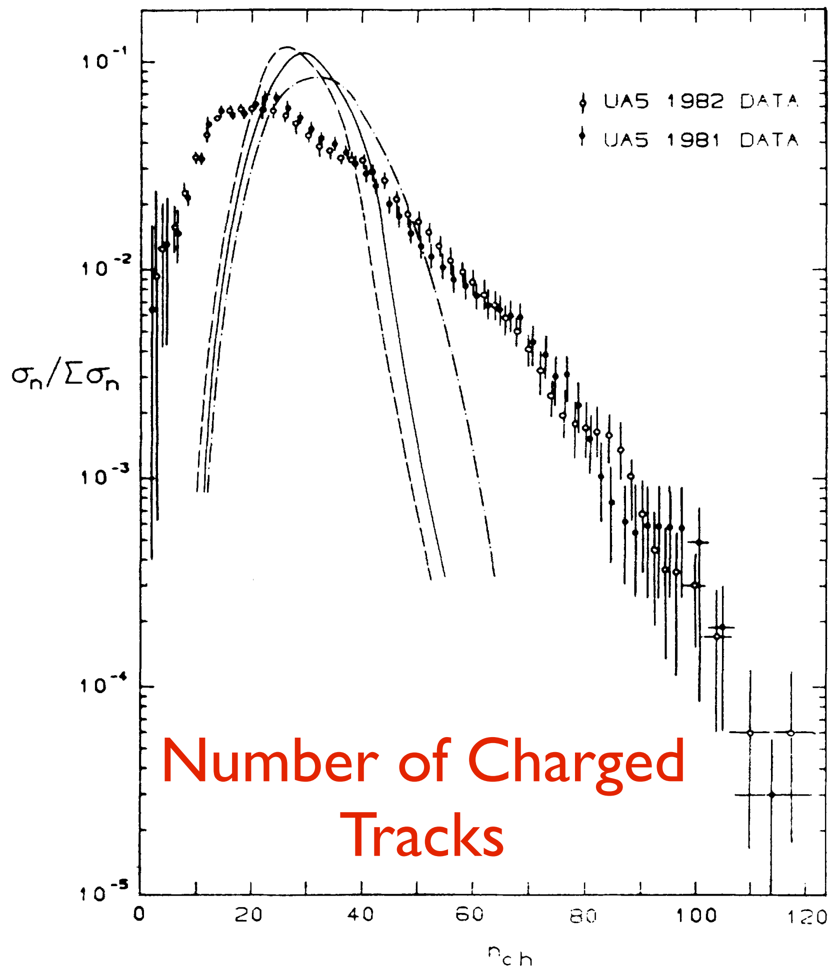
Sjöstrand & v. Zijl, Phys.Rev.D36(1987)2019

Hadron Collisions



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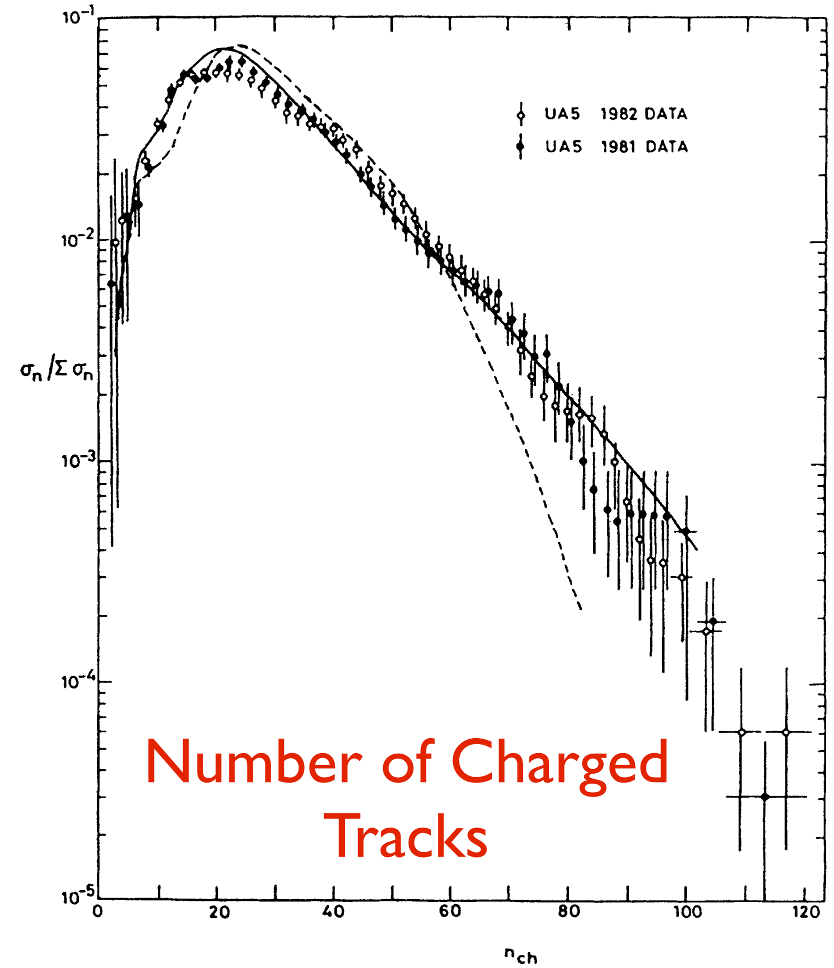


FIG. 12. Charged-multiplicity distribution at 540 GeV, UA5 results (Ref. 32) vs multiple-interaction model with variable impact parameter: solid line, double-Gaussian matter distribution; dashed line, with fix impact parameter [i.e., $\tilde{O}_0(b)$].

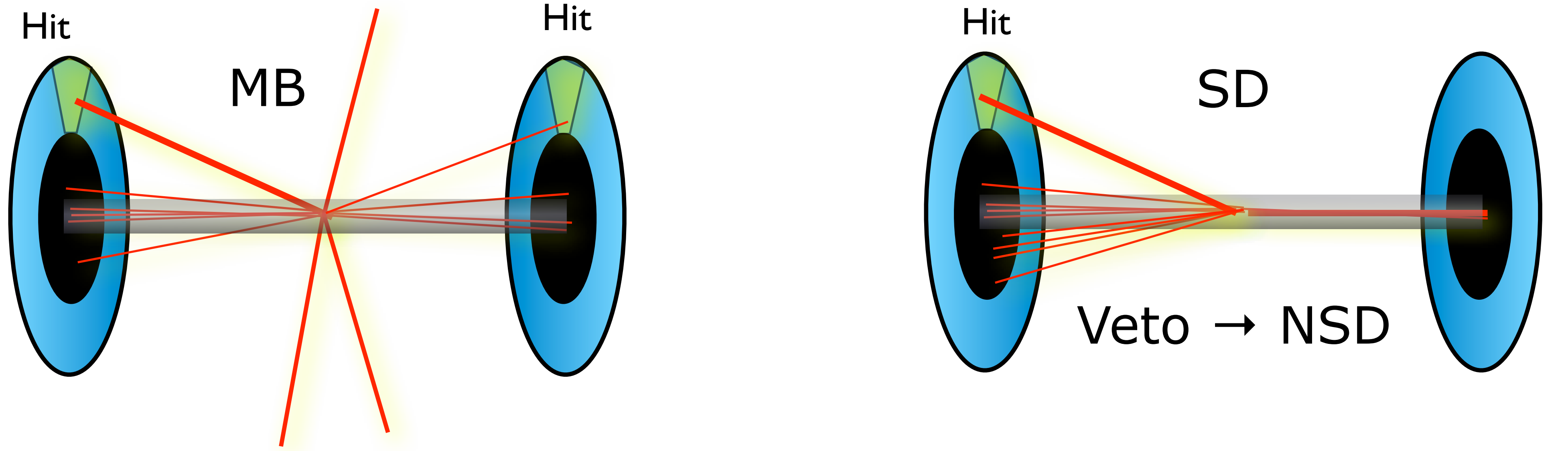
Sjöstrand & v. Zijl, Phys.Rev.D36(1987)2019

What is Pileup / Min-Bias?

We use Minimum-Bias (MB) data to test soft-QCD models

Pileup = “Zero-bias”

“Minimum-Bias” typically suppresses diffraction by requiring two-armed coincidence, and/or ≥ n particle(s) in central region



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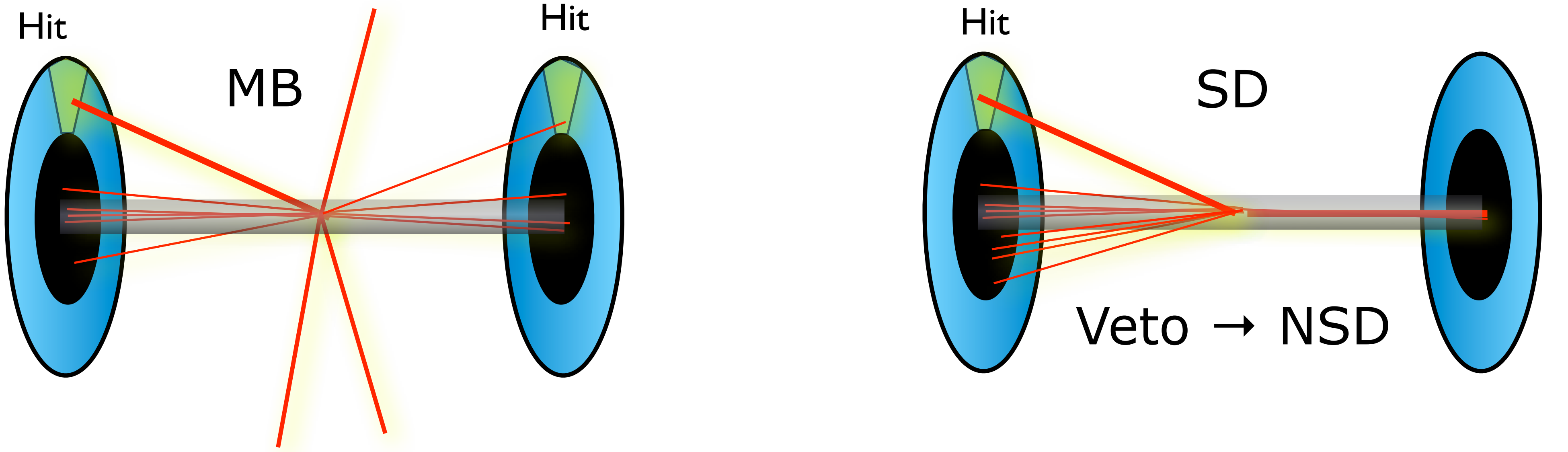


→ Pileup contains more diffraction than Min-Bias

Total diffractive cross section ~ 1/3 σ_{inel}

Most diffraction is low-mass → no contribution in central regions

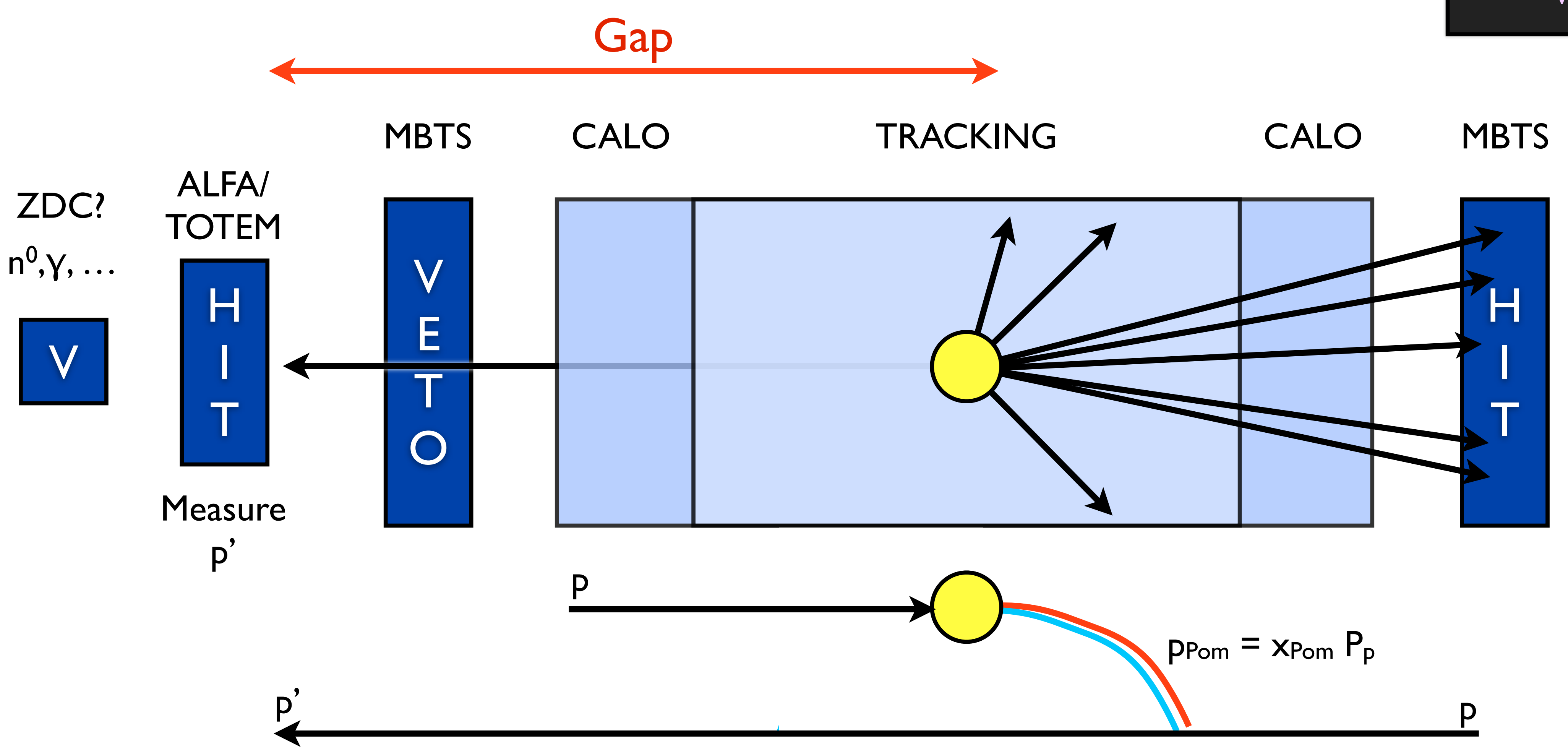
High-mass tails could be relevant in FWD region

→ direct constraints on diffractive components (→ later)

What is diffraction?

Single Diffraction

Glueball-Proton Collider with variable E_{CM}



What is diffraction?

Single Diffraction

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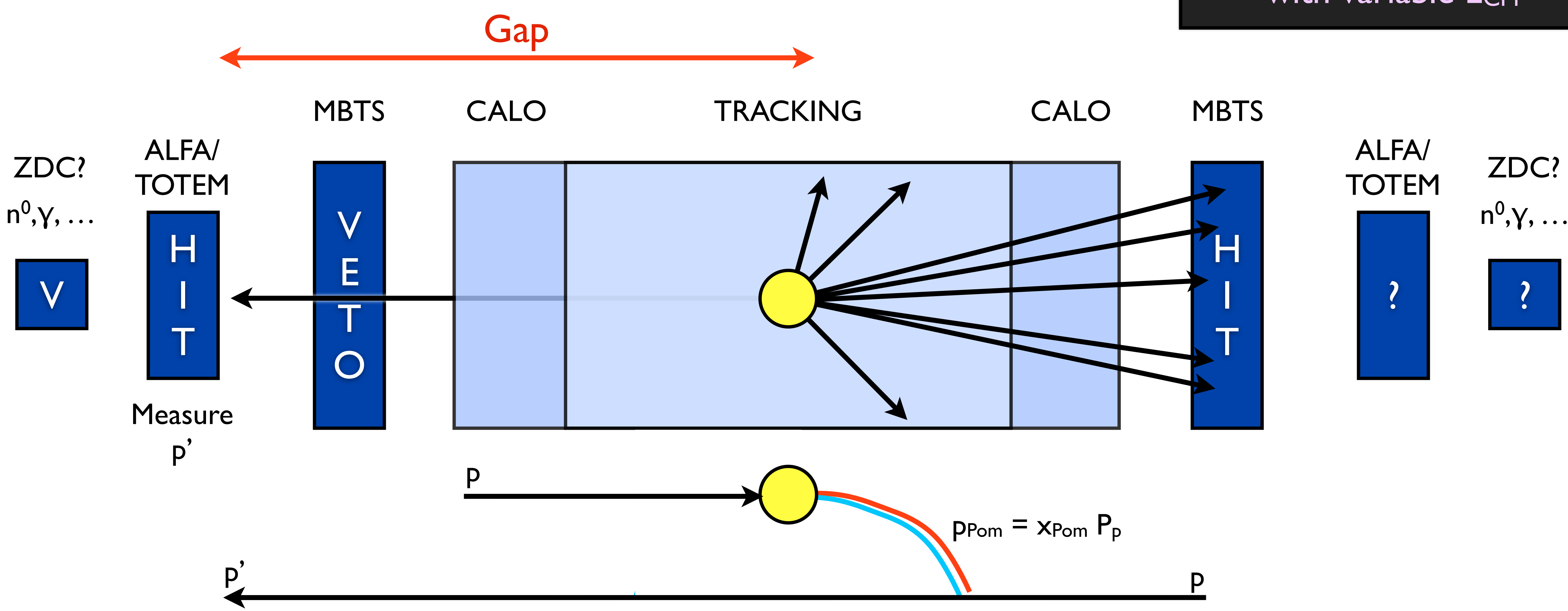
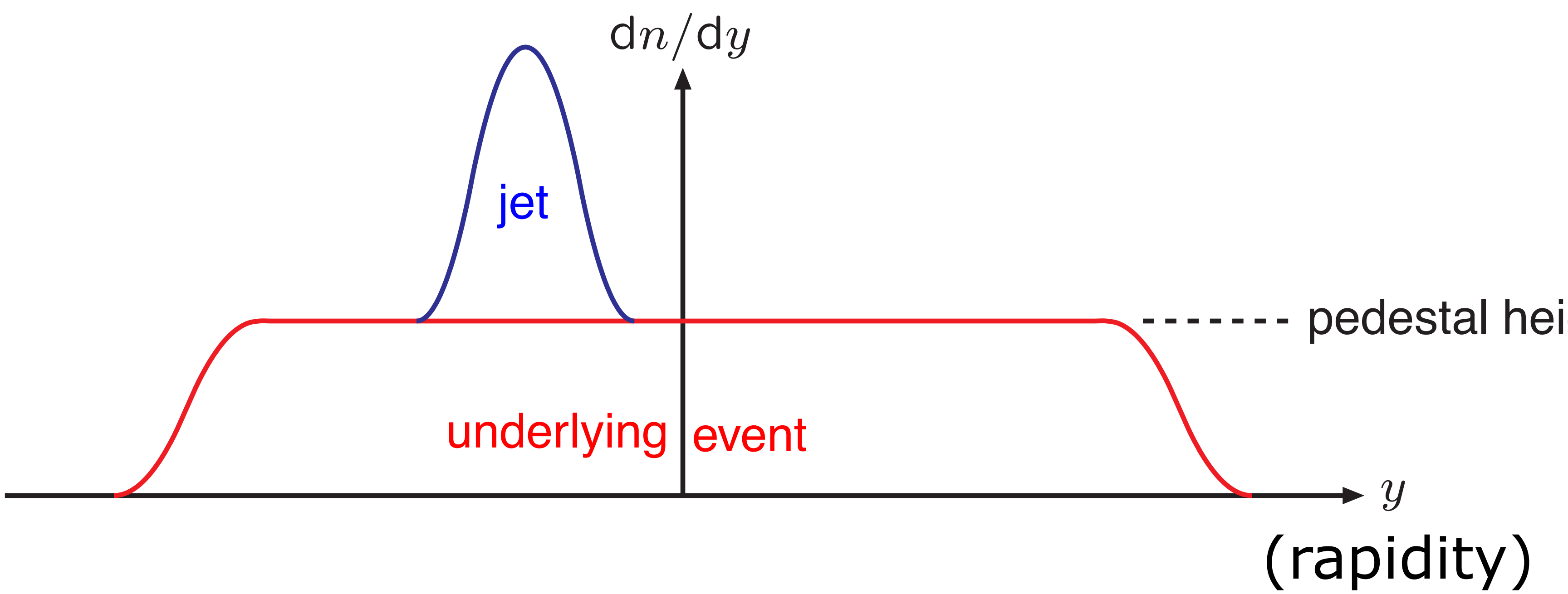


Double Diffraction: both protons explode; gap inbetween
Central Diffraction: two protons + a central (exclusive) system

What is Underlying Event ?

"Pedestal Effect"



Useful variable in hadron collisions: **Rapidity** (now along beam axis)

Designed to be additive under Lorentz Boosts along beam (z) direction

$$y = \frac{1}{2}\ln\left(\frac{E+p_z}{E-p_z}\right)$$

$$y \to -\infty \quad \text{for} \quad p_z \to -E \qquad y \to 0 \quad \text{for} \quad p_z \to 0 \qquad y \to \infty \quad \text{for} \quad p_z \to E$$

Illustrations by T. Sjöstrand

Questions

Pileup

- How much? In central & fwd acceptance?
- Structure: averages + fluctuations, particle composition, lumpiness, …
- Scaling to 13 TeV and beyond

Underlying Event ~ "A handful of pileup" ?

- Hadronizes with Main Event → "Color reconnections"
- Additional "minijets" from multiple parton interactions

Hadronization

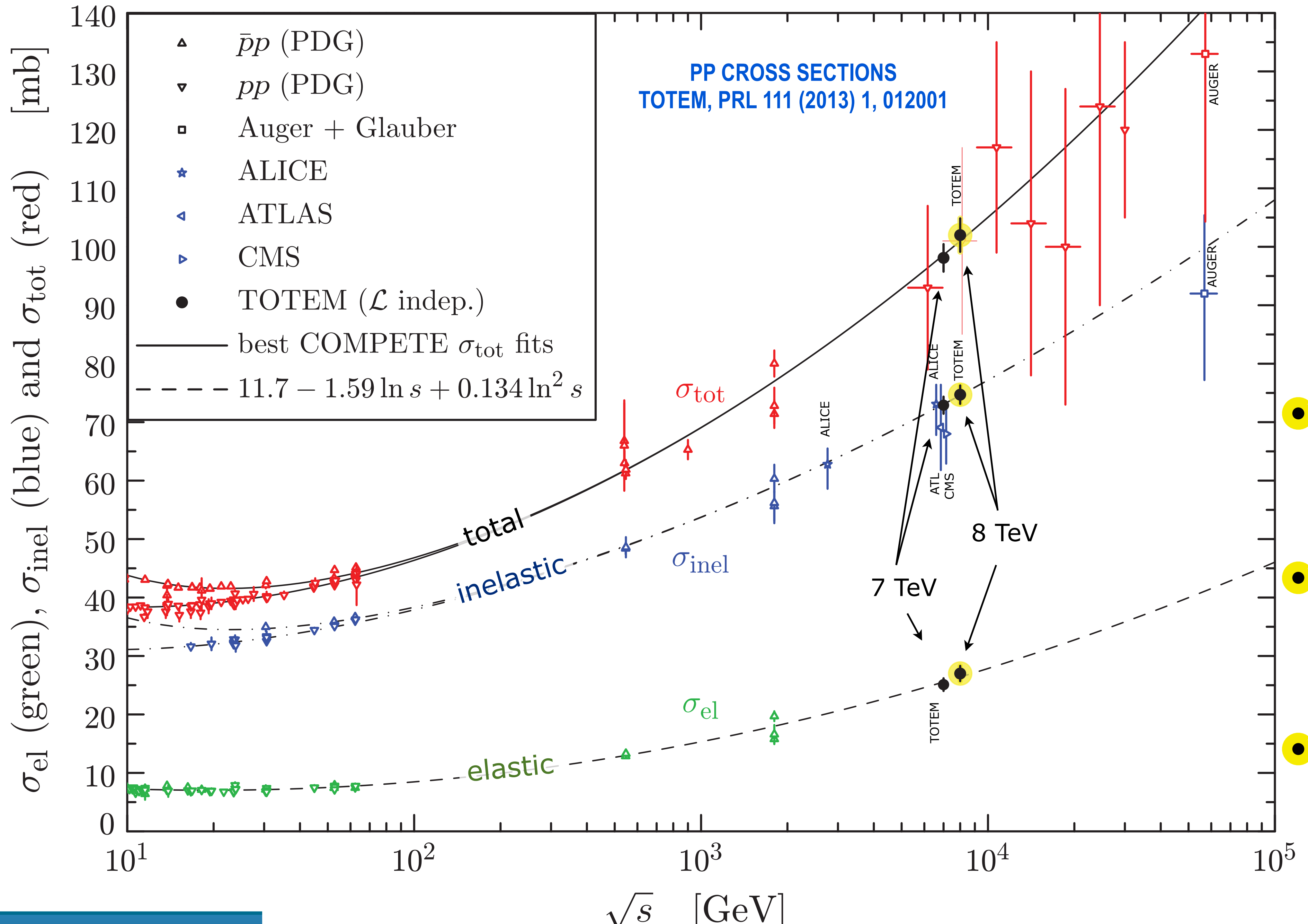
- Models from the 80ies, mainly constrained in 90ies
- Meanwhile, perturbative models have evolved
  - Dipole/Antenna showers, ME matching, NLO corrections, …
  - Precision → re-examine non-perturbative models and constraints
  - New clean constraints from LHC (& future colliders)?
- Hadronization models ⇄ analytical NP corrections?

Uses and Limits of "Tuning"

The Total Cross Section

$$\text{Pileup rate} \propto \sigma_{\text{tot}}(s) = \sigma_{\text{el}}(s) + \sigma_{\text{inel}}(s) \propto s^{0.08} \text{ or } \ln^2(s)\ ?$$

Donnachie-Landshoff Froissart-Martin Bound



(2.9%)
$\sigma_{\text{tot}}(8\text{ TeV}) = 101 \pm 2.9$ mb

(2.3%)
$\sigma_{\text{inel}}(8\text{ TeV}) = 74.7 \pm 1.7$ mb

(5.1%)
$\sigma_{\text{el}}(8\text{ TeV}) = 27.1 \pm 1.4$ mb

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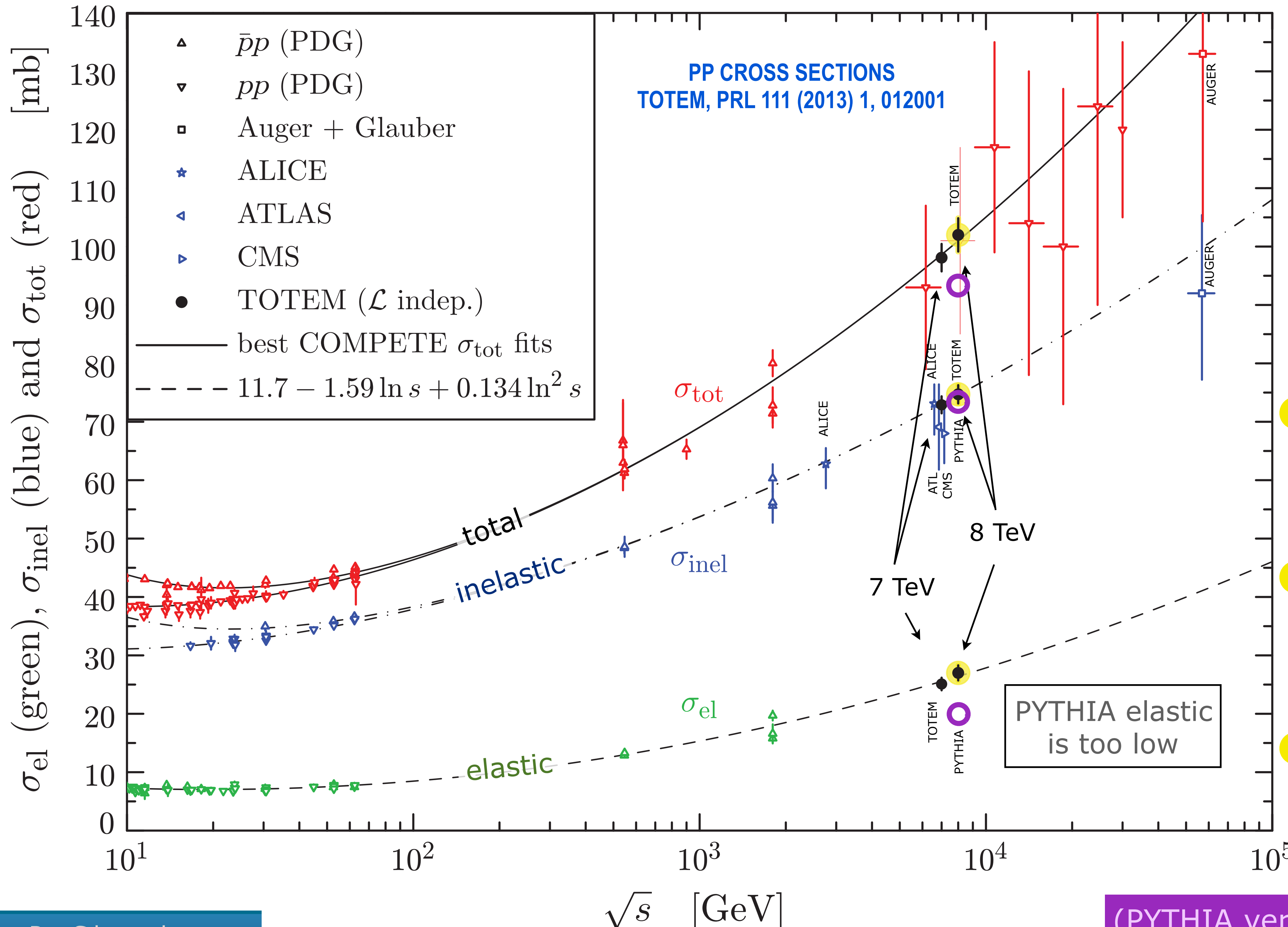

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(PYTHIA versions: 6.4.28 & 8.1.80)

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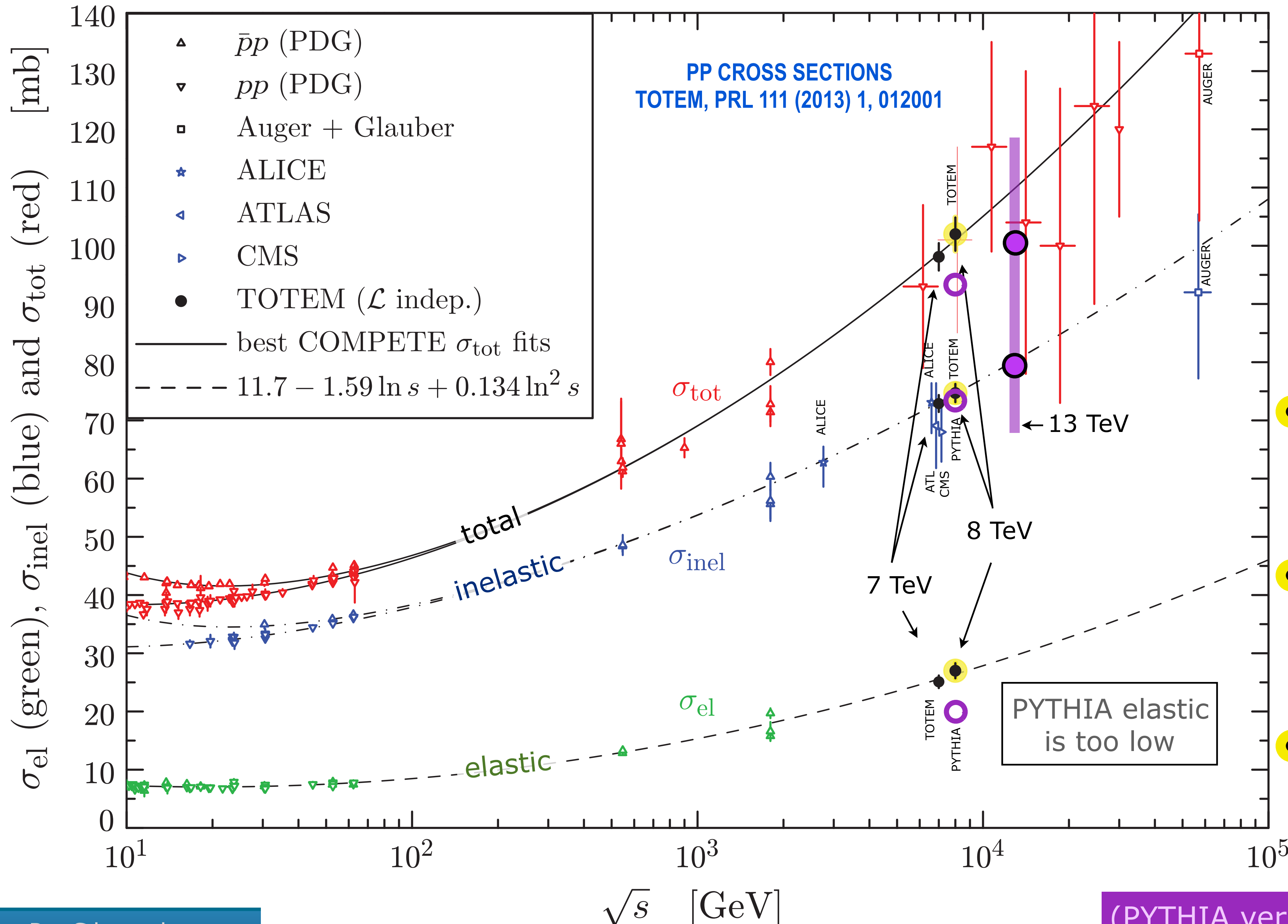

$\sigma_{\rm tot}(13\ \text{TeV}) \sim 110 \pm 6$ mb

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$\sigma_{\rm inel}(13\ \text{TeV}) \sim 80 \pm 3.5$ mb

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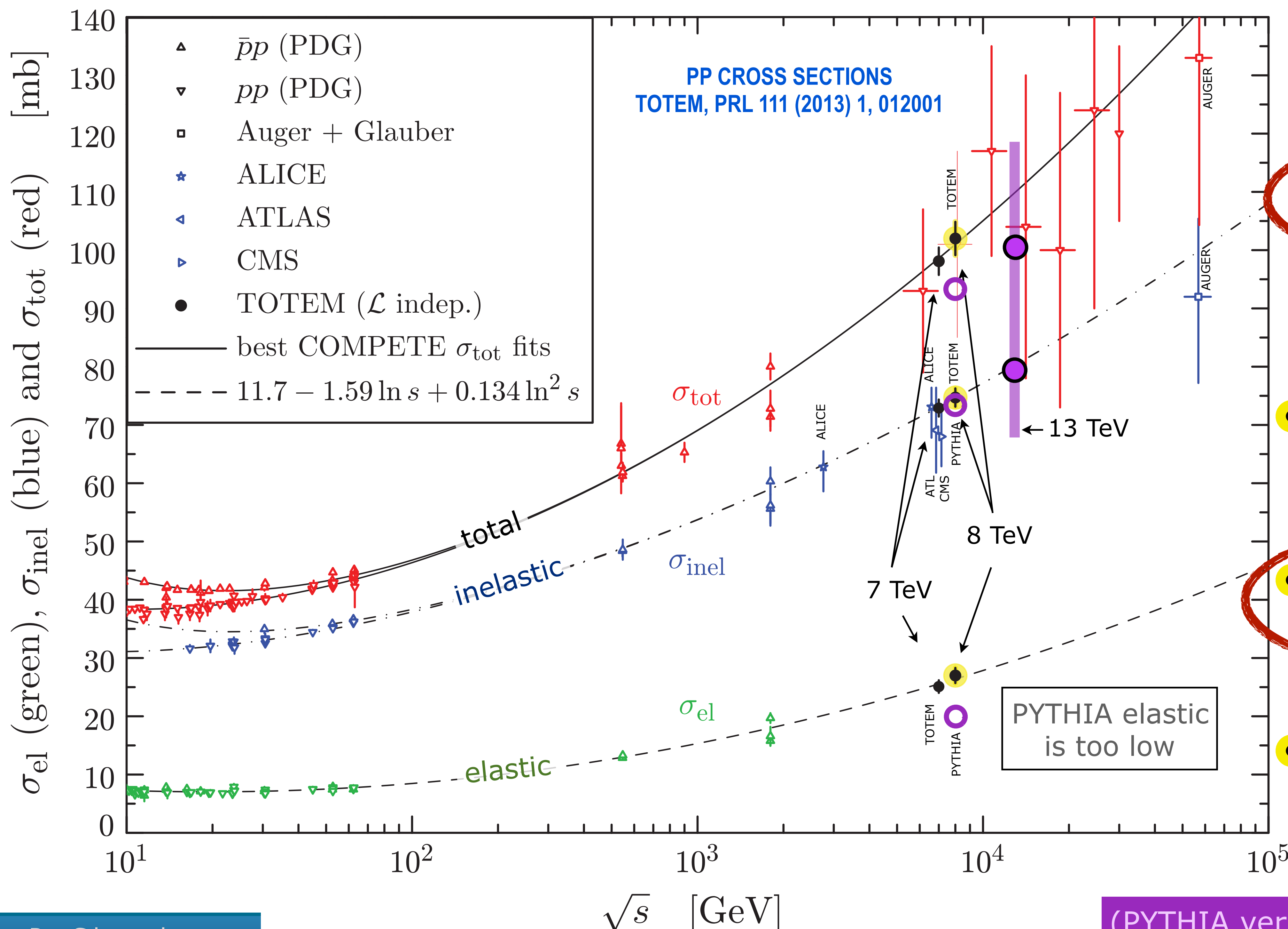

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The Inelastic Cross Section

First try: decompose $\sigma_{\rm inel} = \sigma_{\rm sd} + \sigma_{\rm dd} + \sigma_{\rm cd} + \sigma_{\rm nd}$

\+ Parametrizations of diffractive components: dM^2/M^2

PYTHIA:

$$\frac{\mathrm{d}\sigma_{\mathrm{sd}(AX)}(s)}{\mathrm{d}t\,\mathrm{d}M^2} = \frac{g_{3\mathbb{P}}}{16\pi}\,\beta_{A\mathbb{P}}^2\,\beta_{B\mathbb{P}}\,\frac{1}{M^2}\,\exp(B_{\mathrm{sd}(AX)}t)\,F_{\mathrm{sd}}\,,$$

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\+ Integrate and solve for $\sigma_{\rm nd}$

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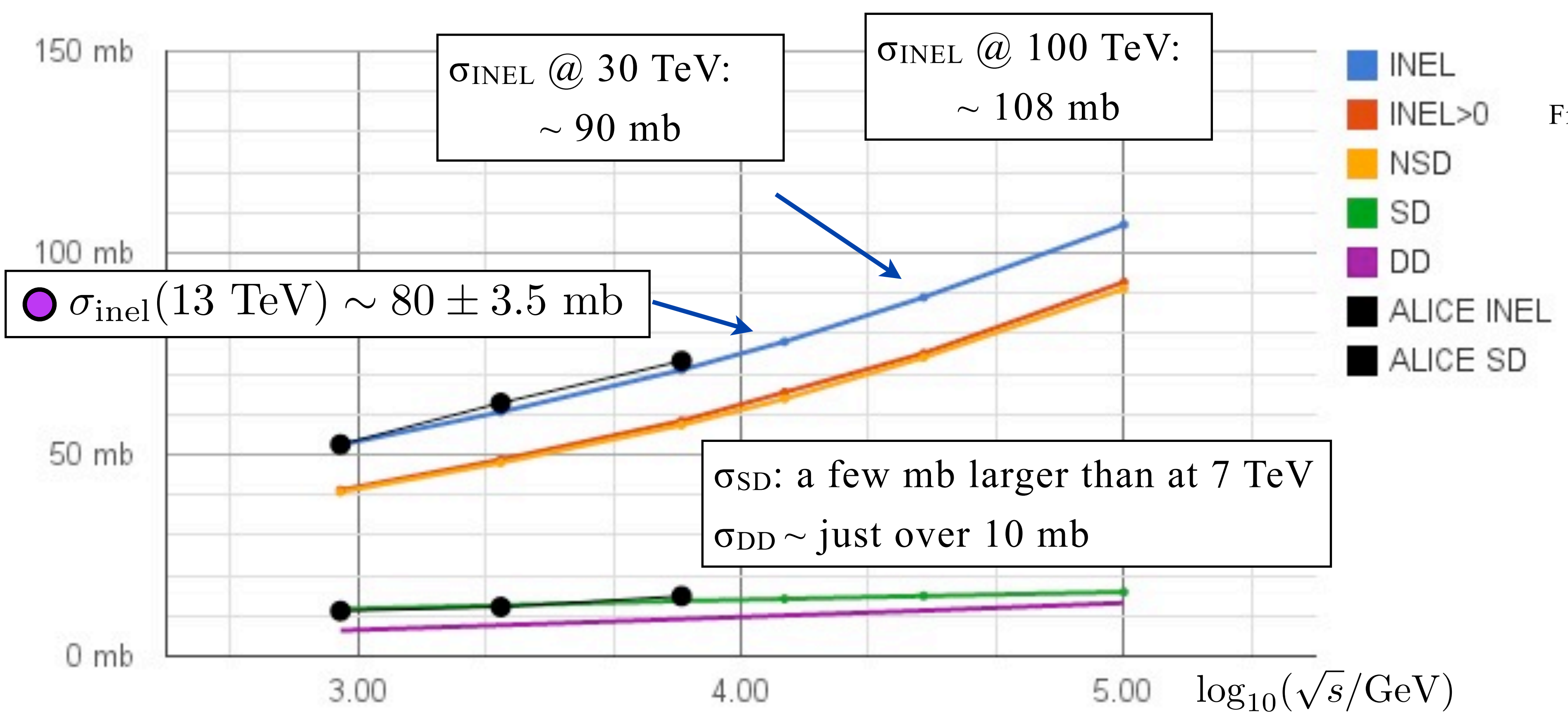

What Cross Section?

INEL	Total Inelastic
INEL>0	Fraction with one charged particle in $|\eta|<1$
NSD	Ambiguous Theory Definition
SD	Ambiguous Theory Definition
DD	Ambiguous Theory Definition
ALICE INEL	Observed fraction corrected to total
ALICE SD	ALICE def : SD has MX<200

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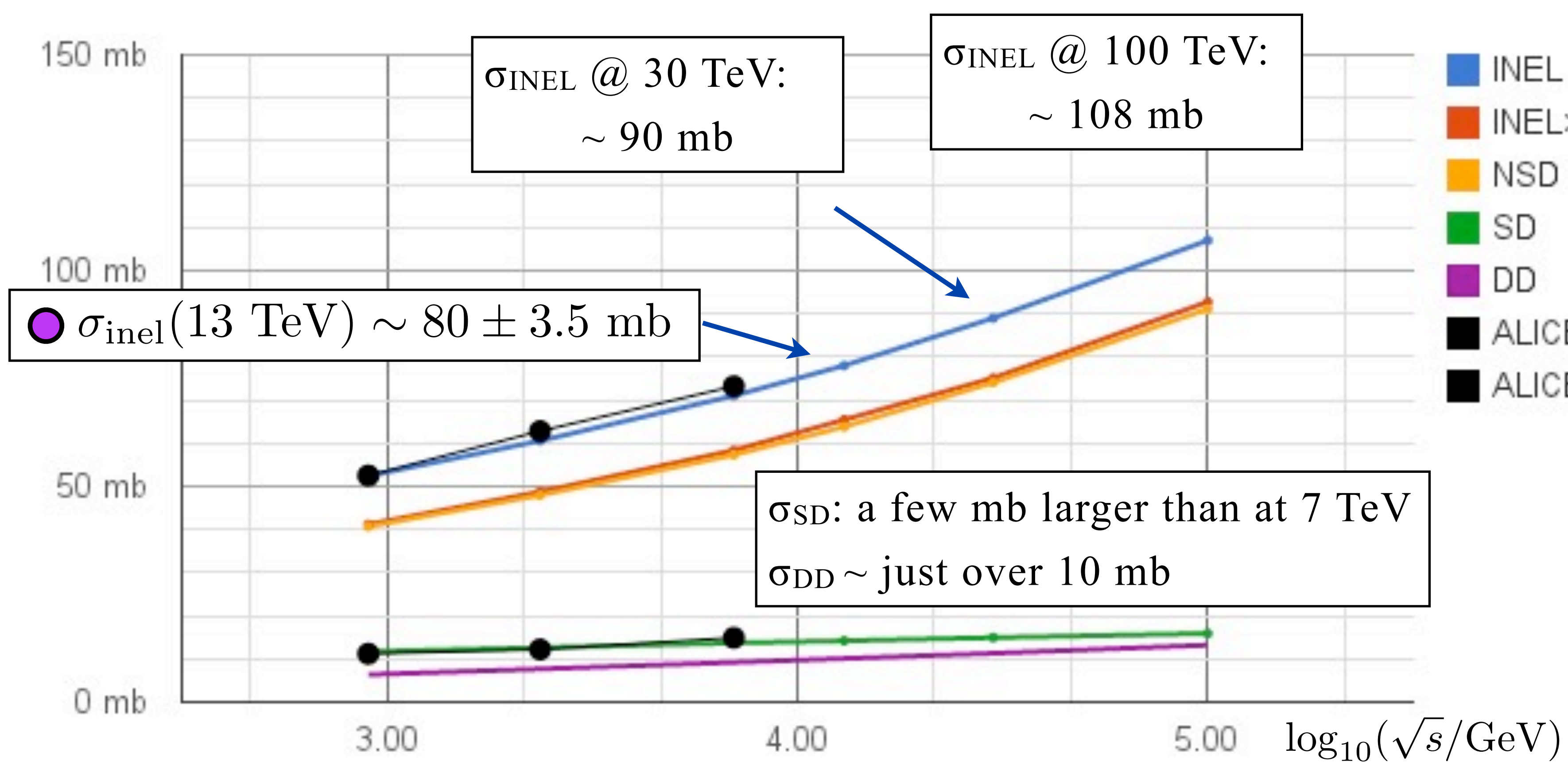

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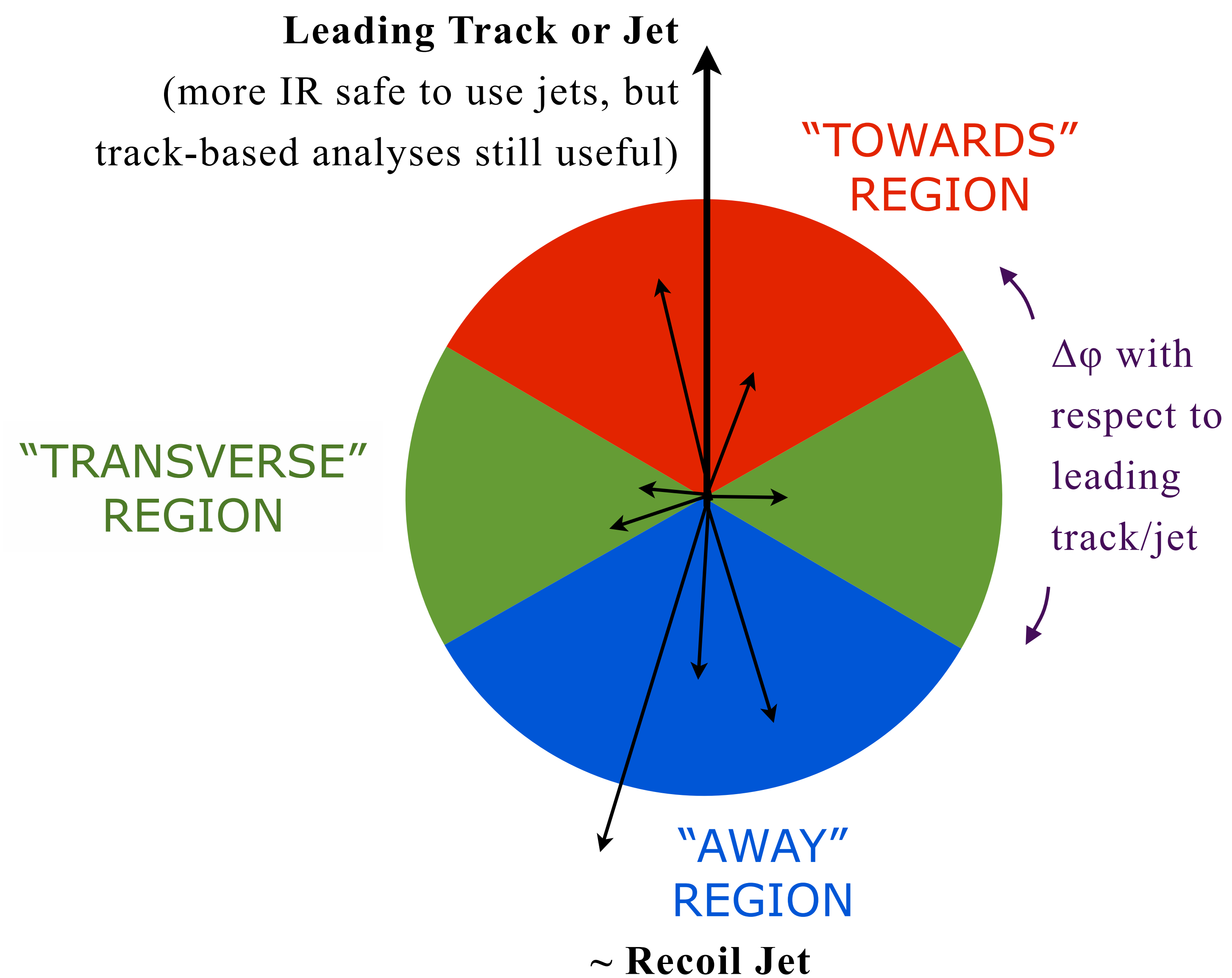
Note problem of principle: Q.M. requires **distinguishable final states**

The "Rick Field" UE Plots

(the same Field as in Field-Feynman)

There are many UE variables.
The most important is $\langle \Sigma p_T \rangle$ in the "Transverse Region"



Transverse Region
(TRNS)

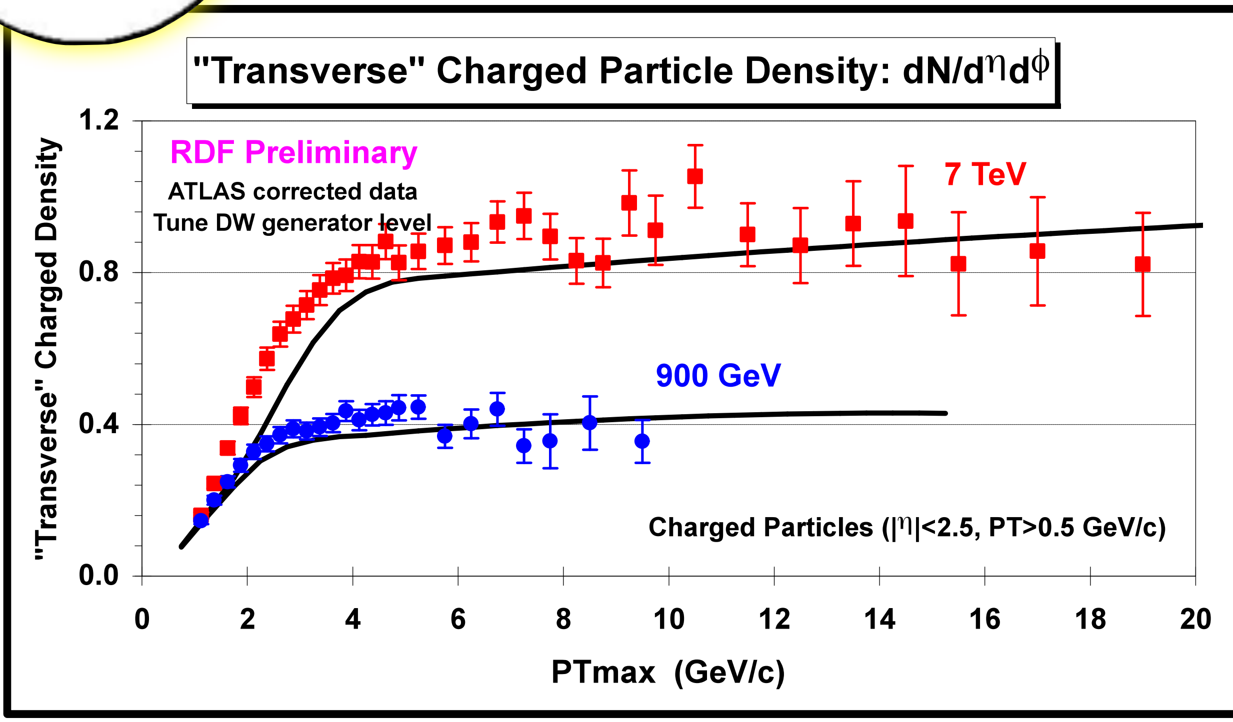
Sensitive to activity at right angles to the hardest jets

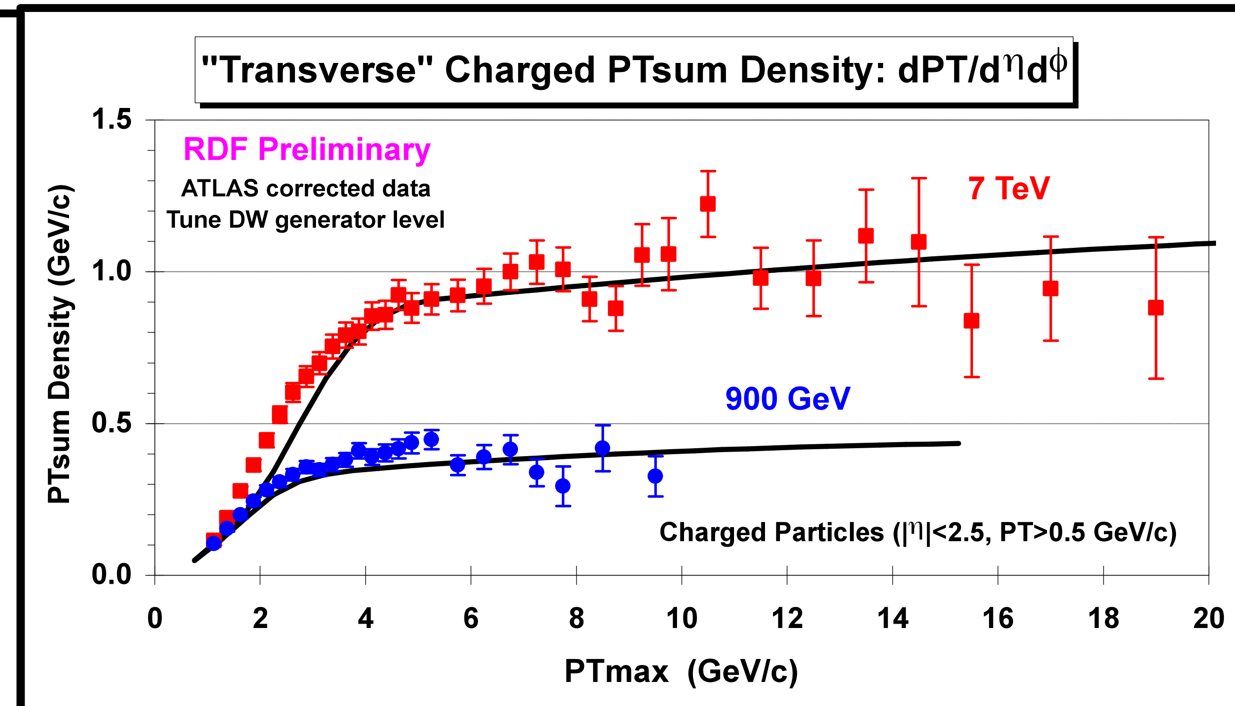
Useful definition of Underlying Event

The Pedestal

(now called the Underlying Event)

LHC from 900 to 7000 GeV - ATLAS





Track Density (TRANS)

Sum(pT) Density (TRANS)

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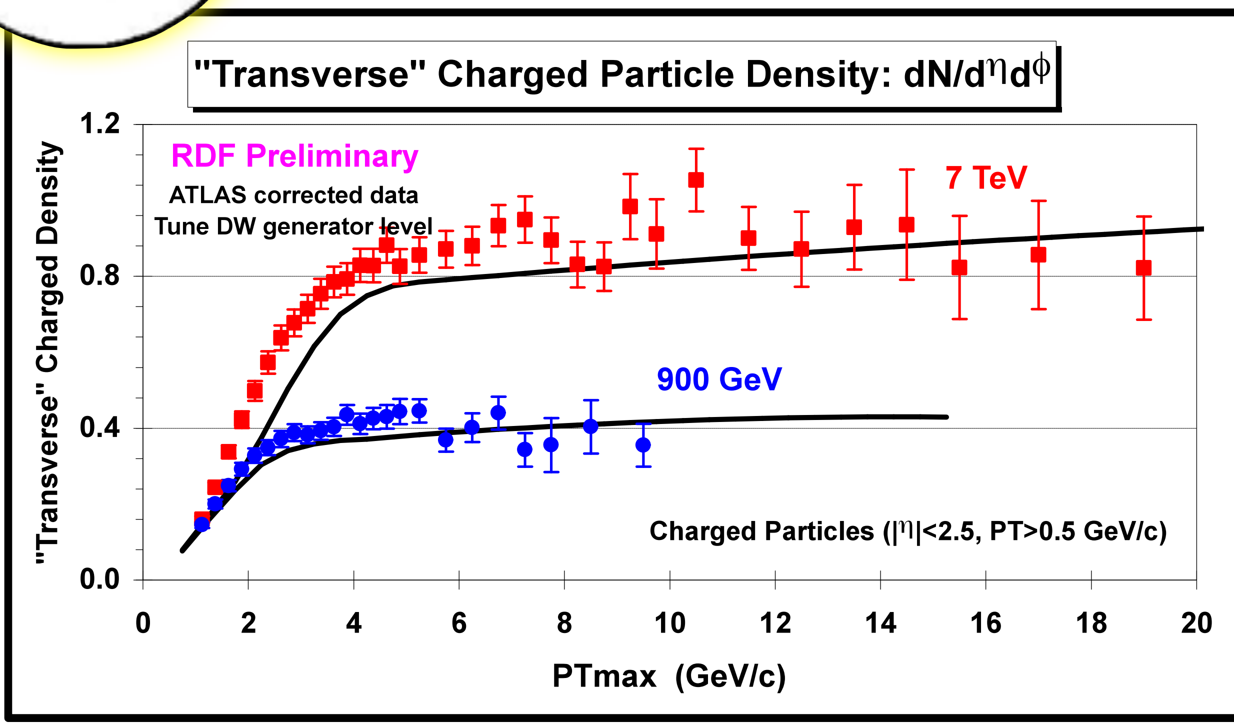

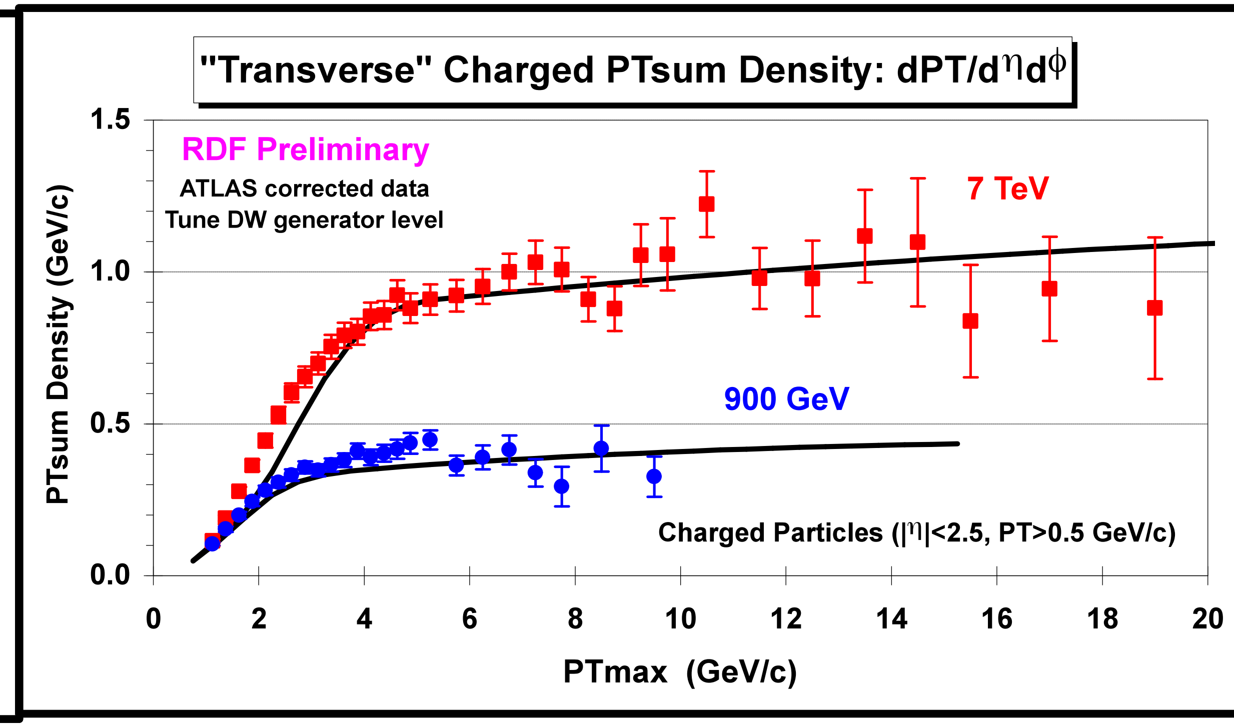

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Sum(pT) Density (TRANS)

Not Infrared Safe

Large Non-factorizable Corrections

Prediction off by ≈ 10%

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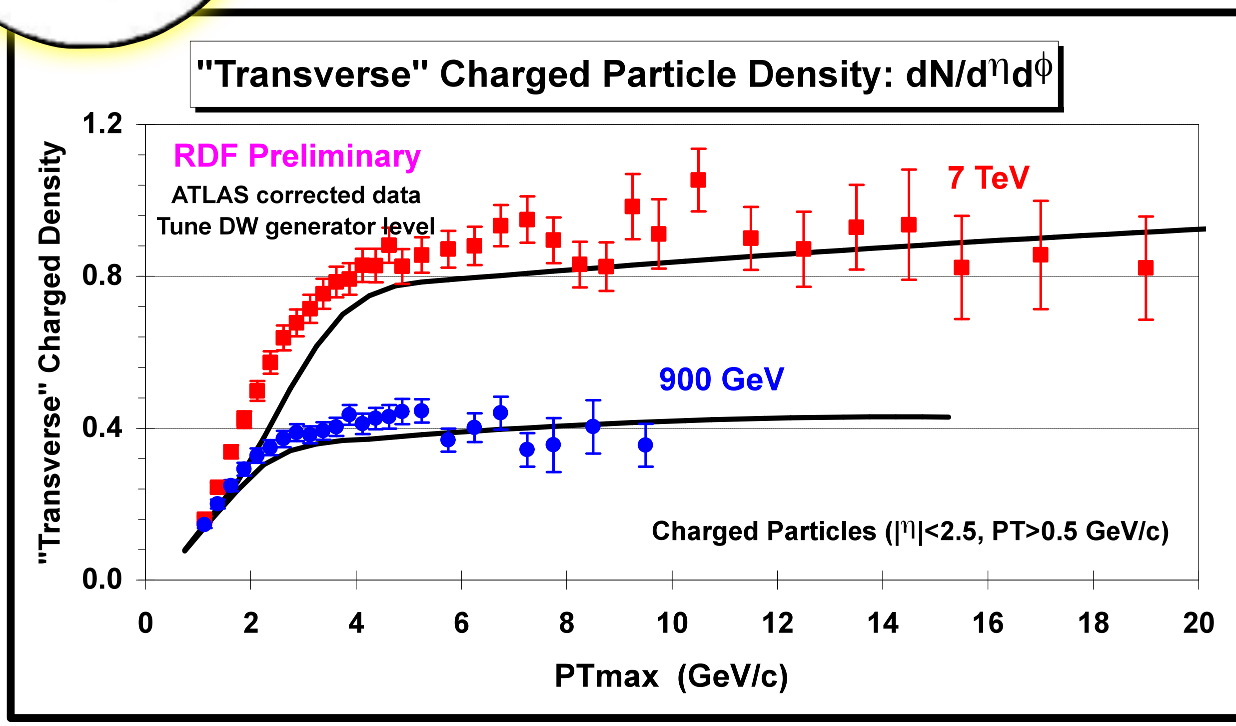

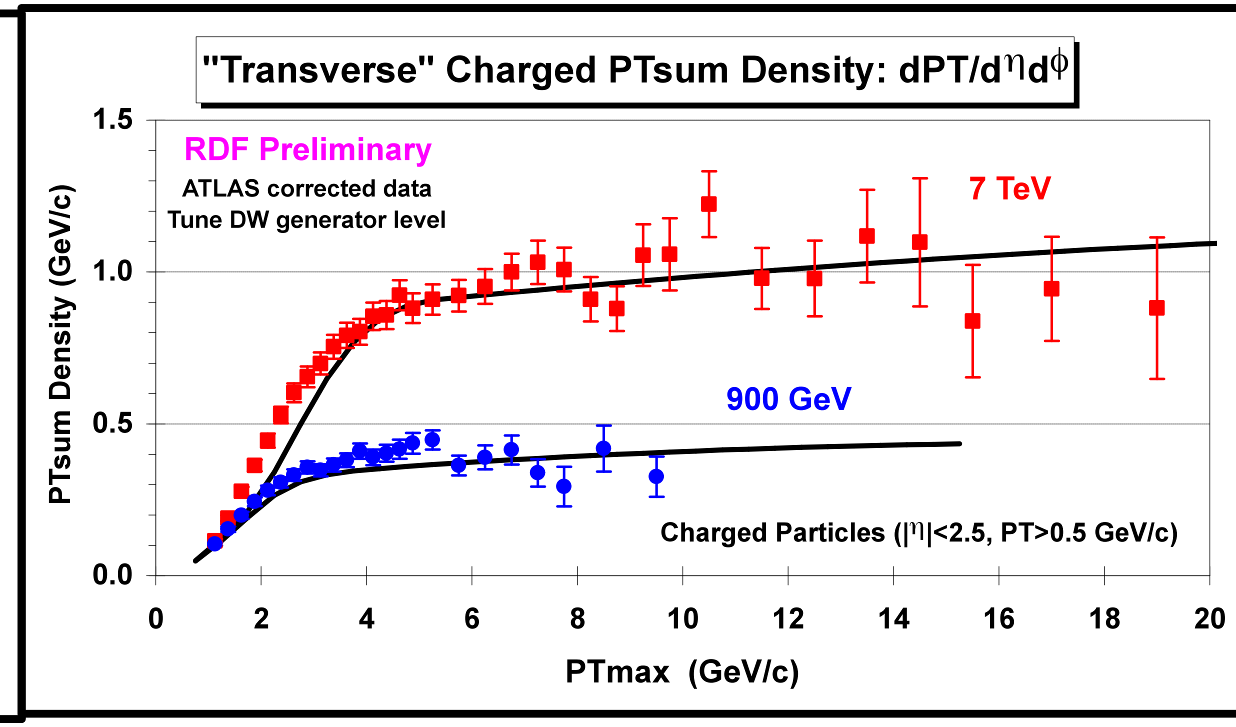

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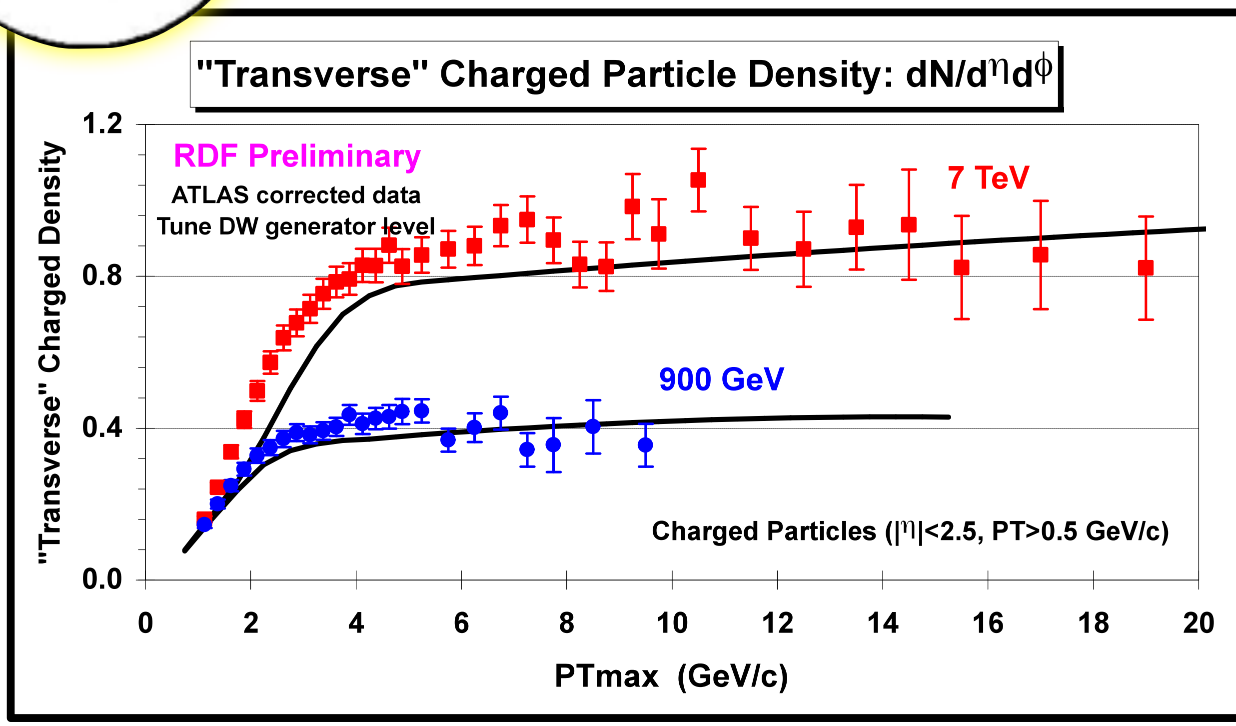

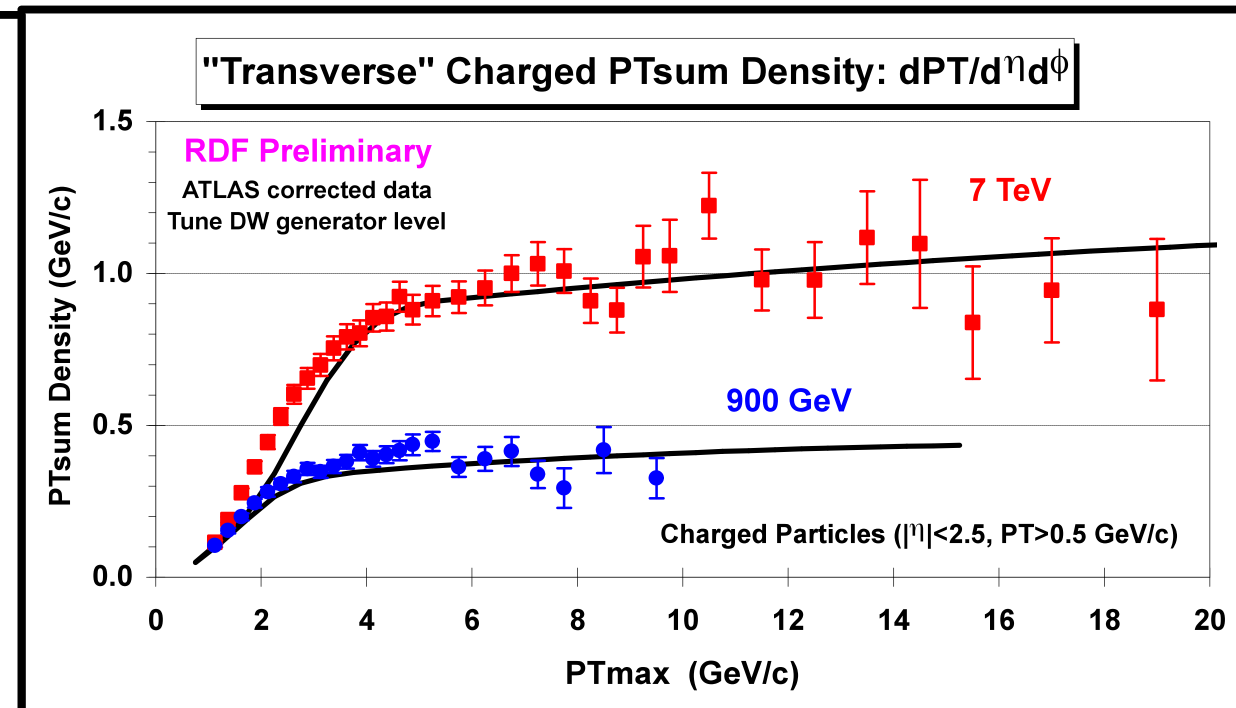

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R. Field: "See, I told you!"

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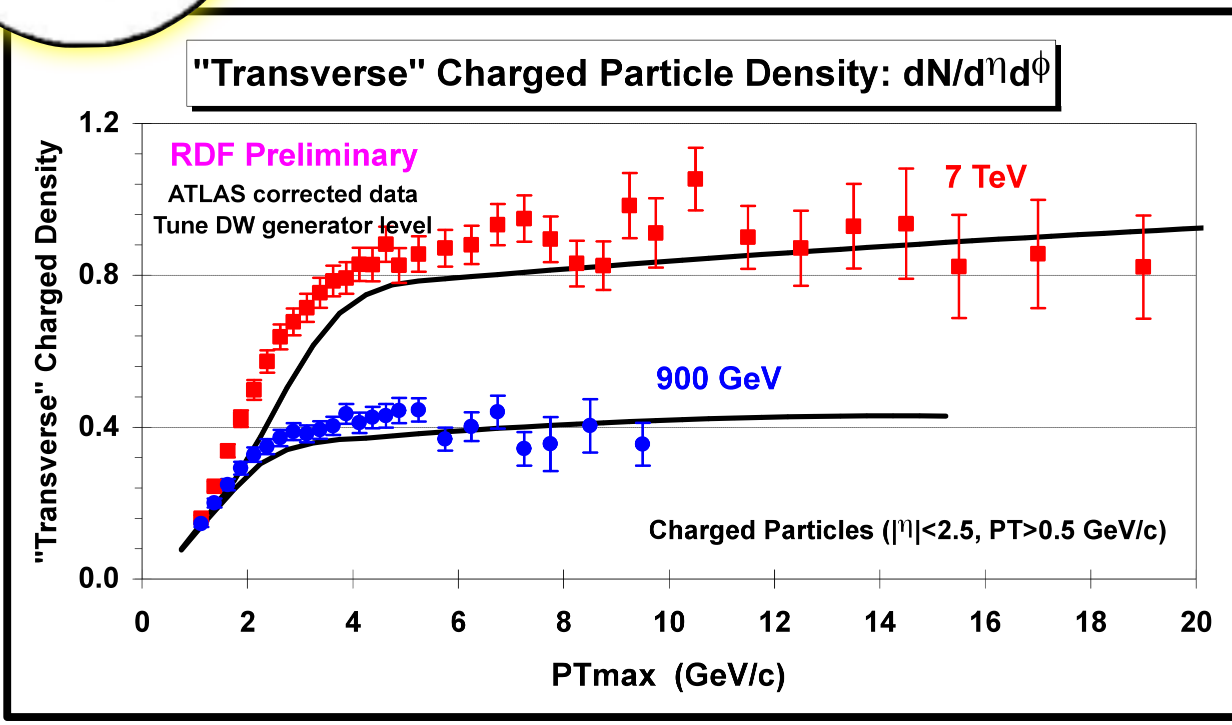

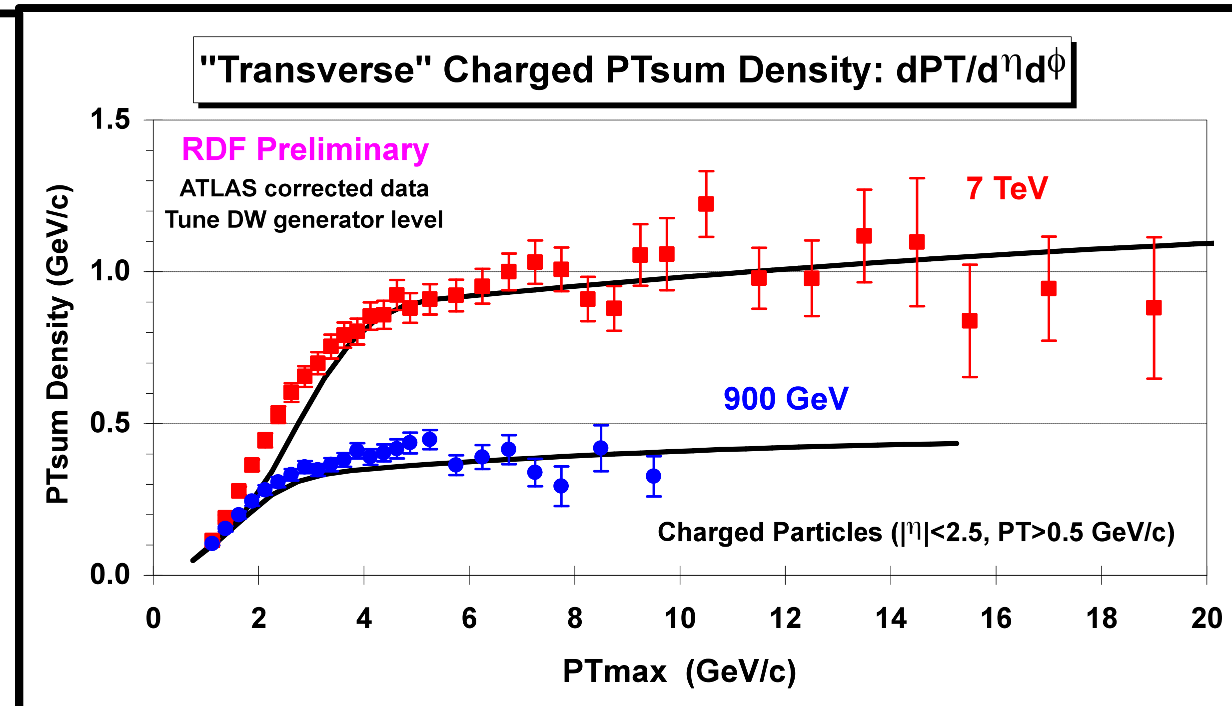

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R. Field: "See, I told you!" Y. Gehrstein: "they have to fudge it again"

The Pedestal
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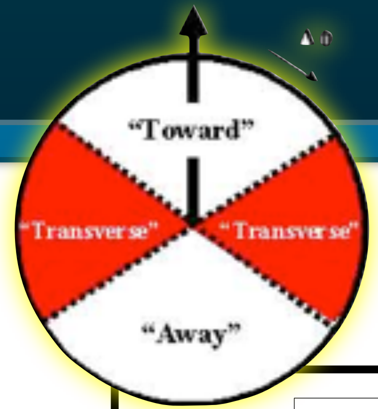

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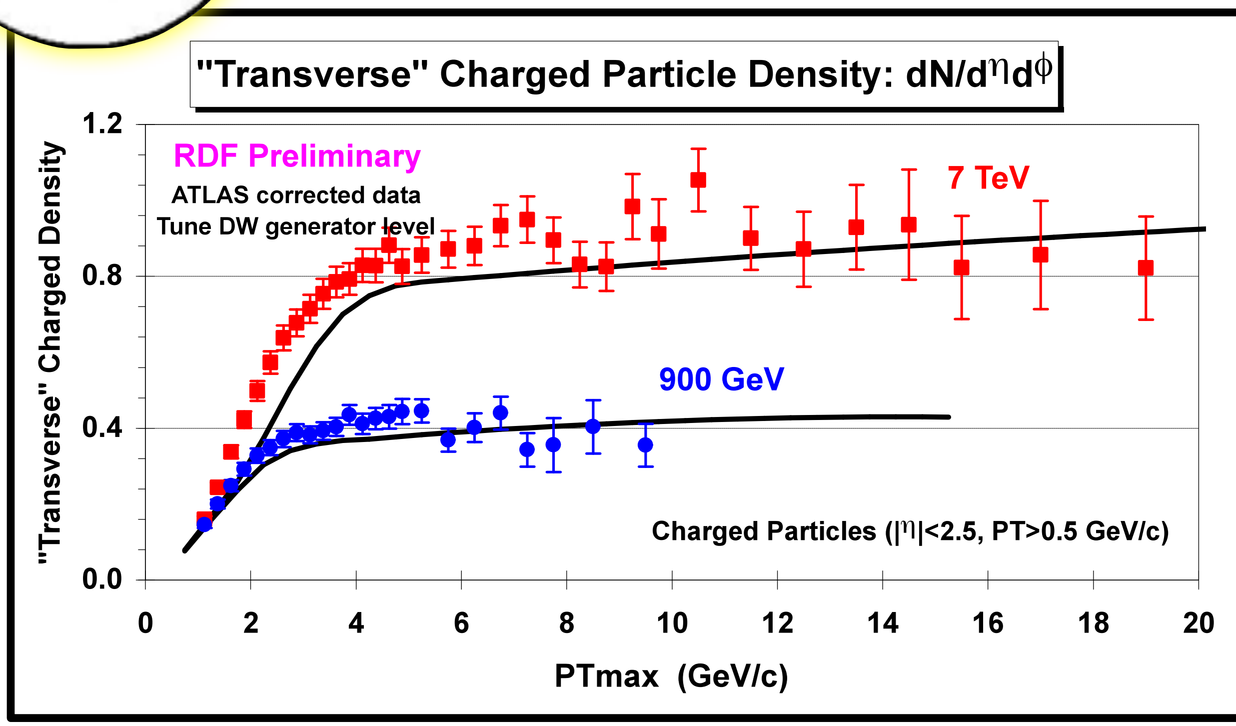

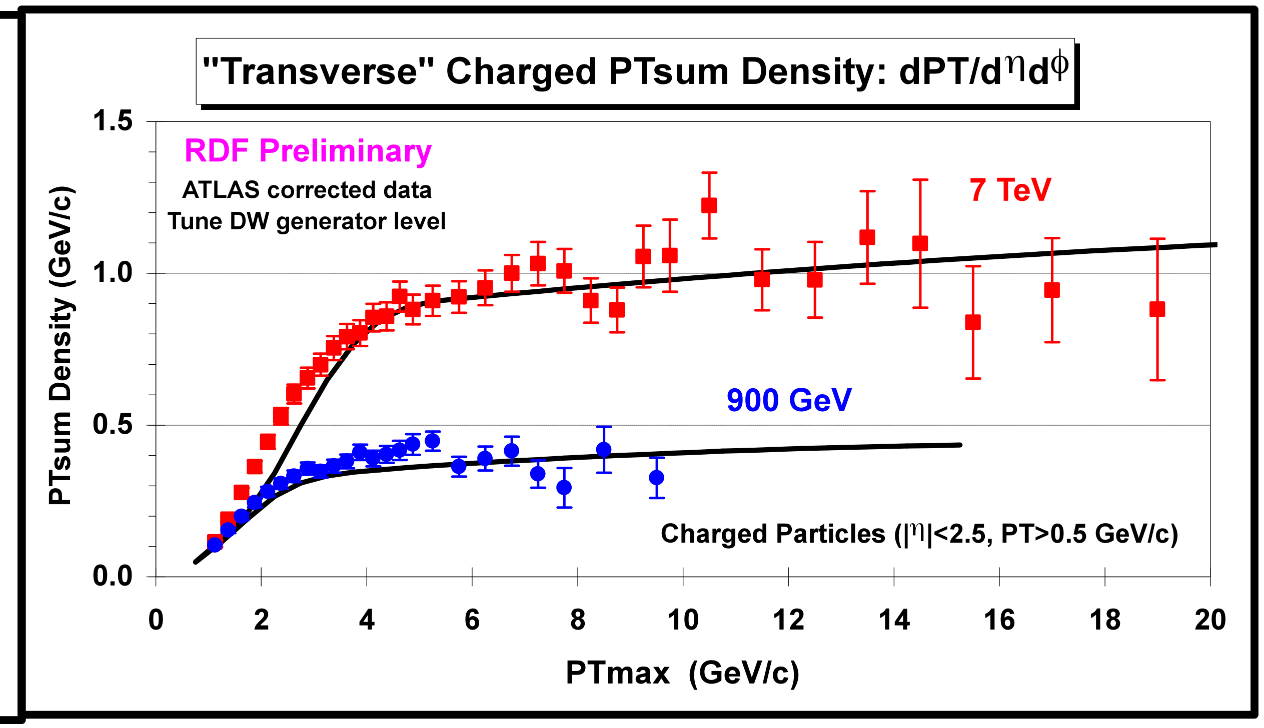

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Truth is in the eye of the beholder:

R. Field: "See, I told you!"

Y. Gehrstein: "they have to fudge it again"

From Hard to Soft

Main tools for high-p_T calculations

Factorization and IR safety

Corrections suppressed by powers of Λ_{QCD}/Q_{Hard}

Soft QCD / Min-Bias / Pileup

NO HARD SCALE

Typical Q scales ~ Λ_{QCD}

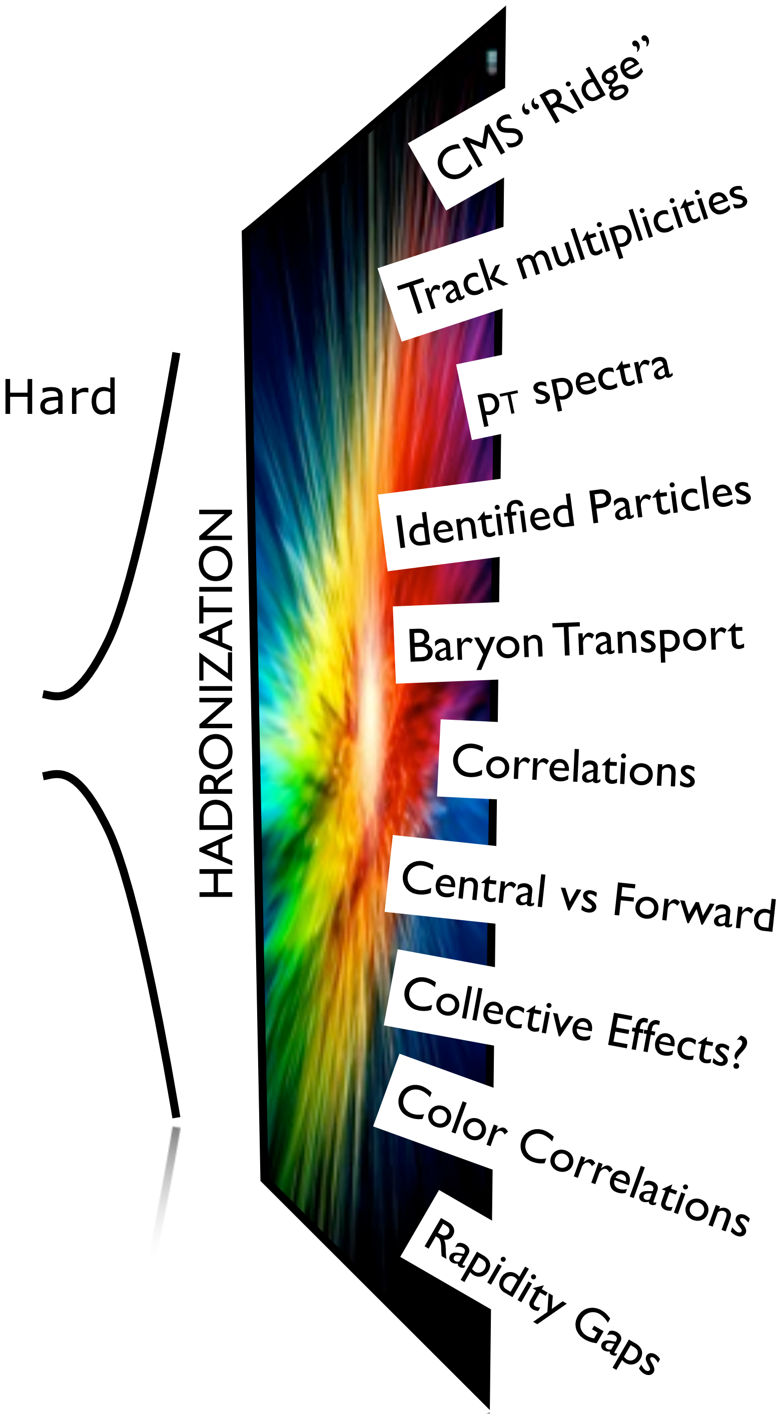
Extremely sensitive to IR effects

→ Excellent LAB for studying IR effects

~ ∞ statistics for min-bias

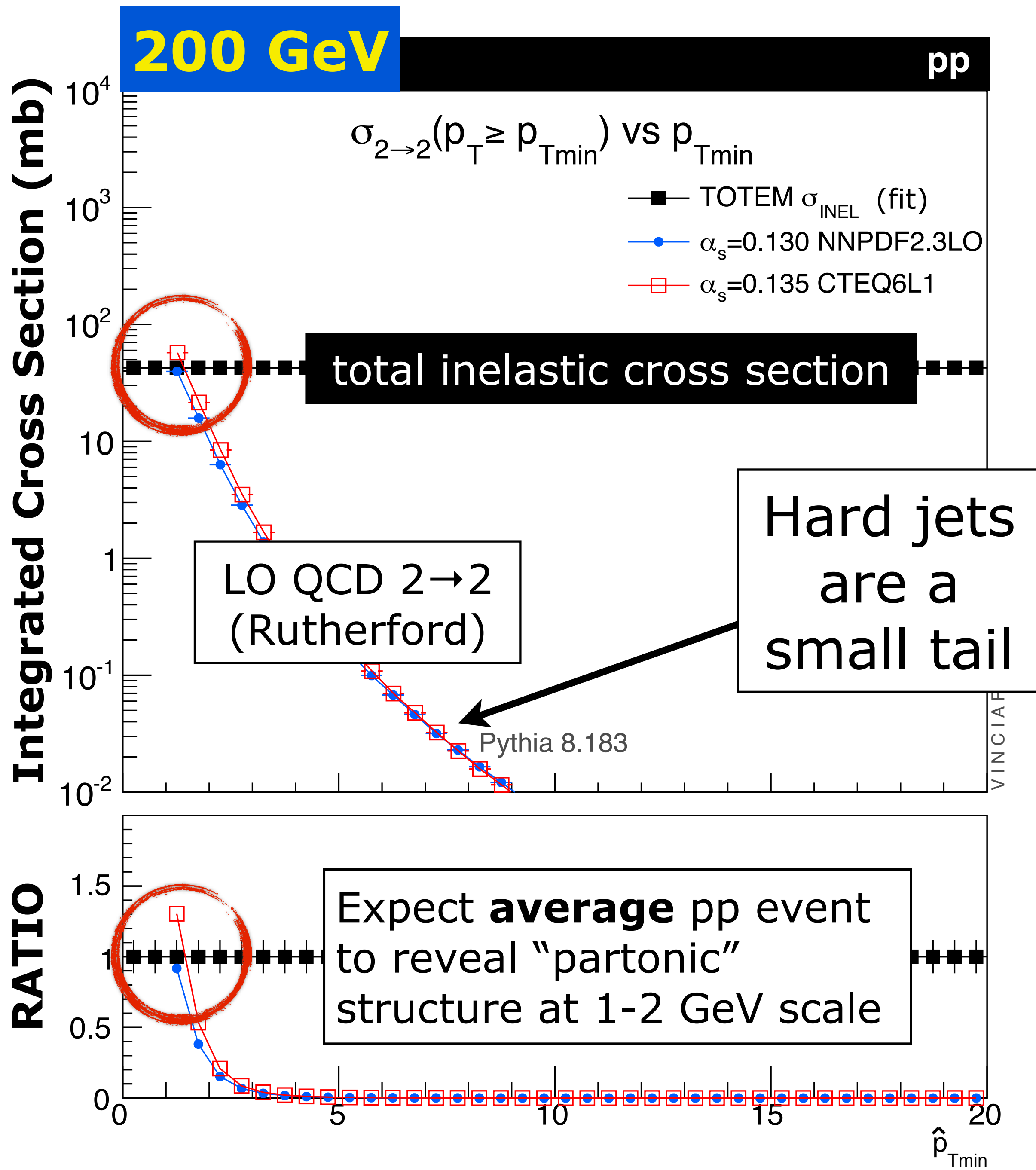
→ Access tails, limits

Universality: Recycling PU ↔ MB ↔ UE



Is there no hard scale?

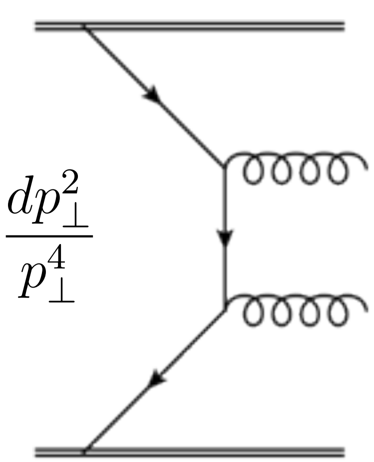
Compare total (inelastic) hadron-hadron cross section to calculated parton-parton (LO QCD 2→2) cross section



Leading-Order pQCD

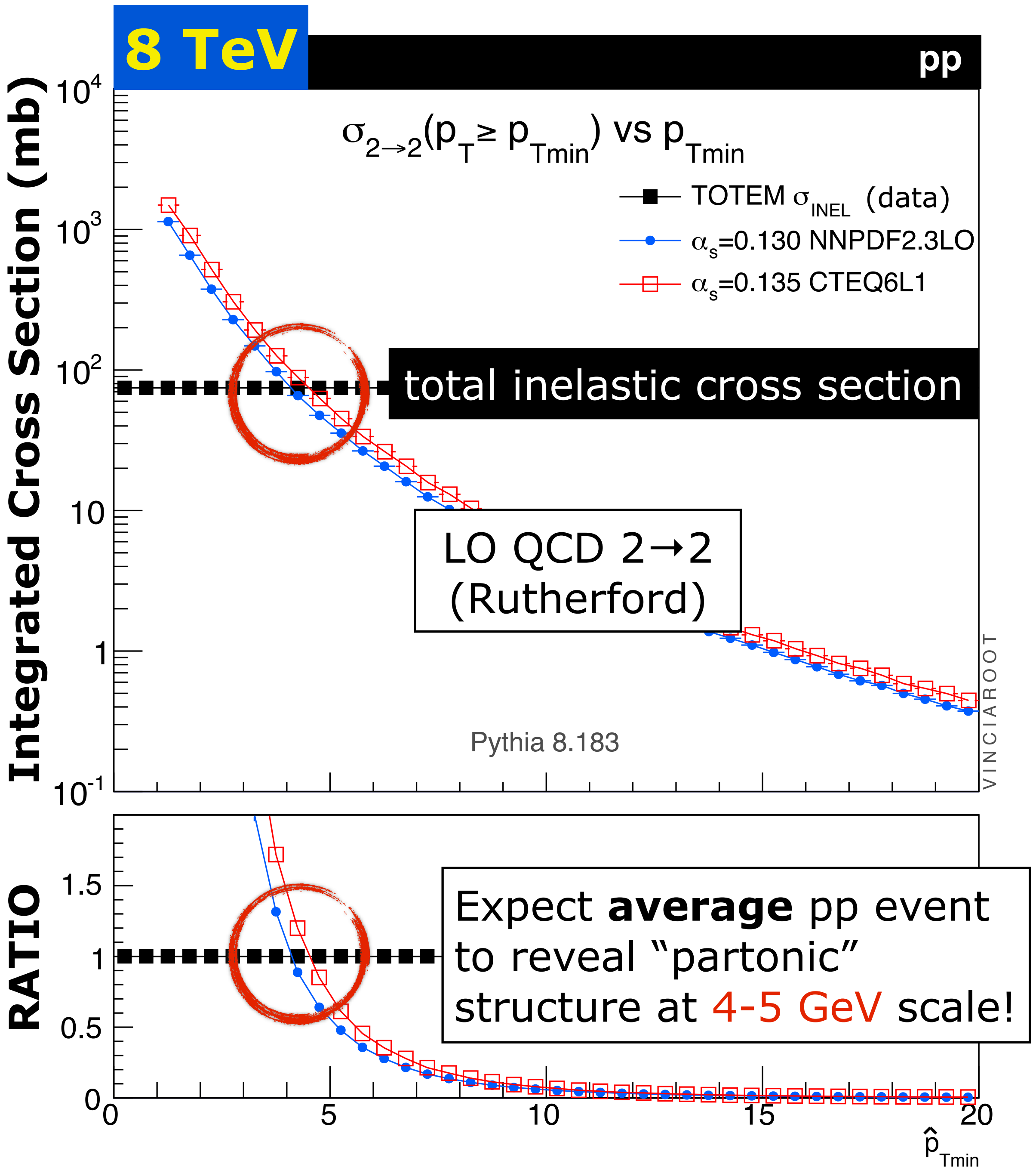
$$\int_{p^2_{\perp,\mathrm{min}}} dp^2_\perp \frac{d\sigma_{\mathrm{Dijet}}}{dp^2_\perp}$$

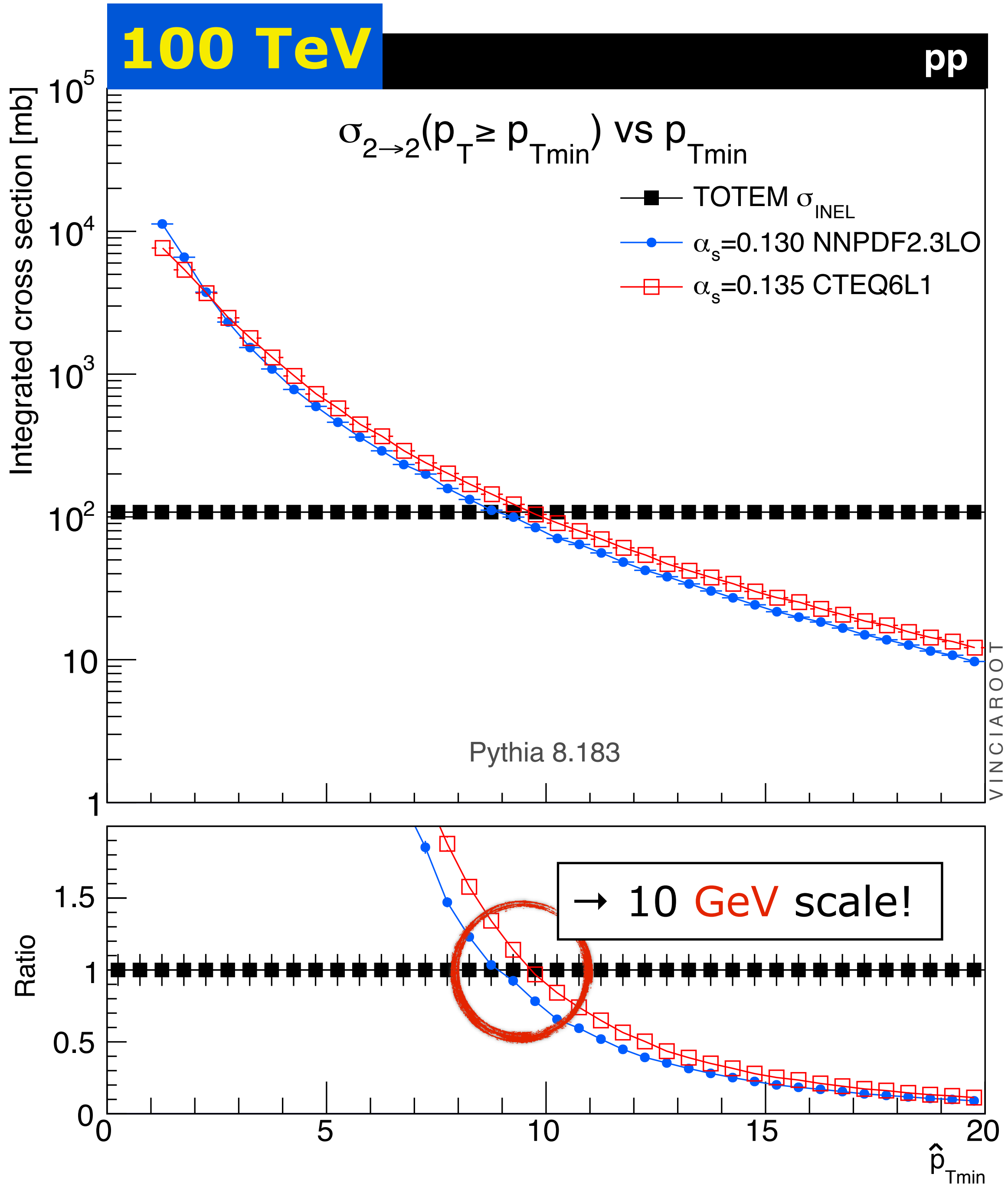
$$d\sigma_{2\to2} \propto \frac{dp^2_\perp}{p^4_\perp} \otimes \mathrm{PDFs}$$



→ 8 TeV → 100 Tev

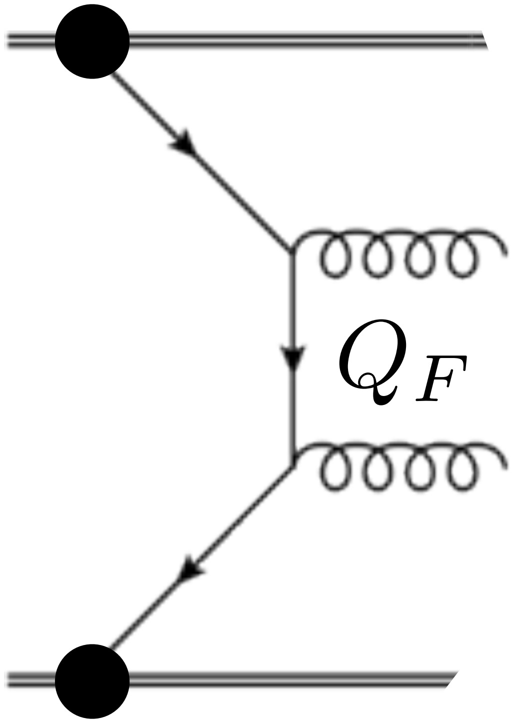
→ Trivial calculation indicates hard scales in min-bias





Physics of the Pedestal

Factorization: Subdivide Calculation



Multiple Parton Interactions go beyond existing theorems

→ perturbative short-distance physics in Underlying Event

→ Need to generalize factorization to MPI

Physics of the Pedestal

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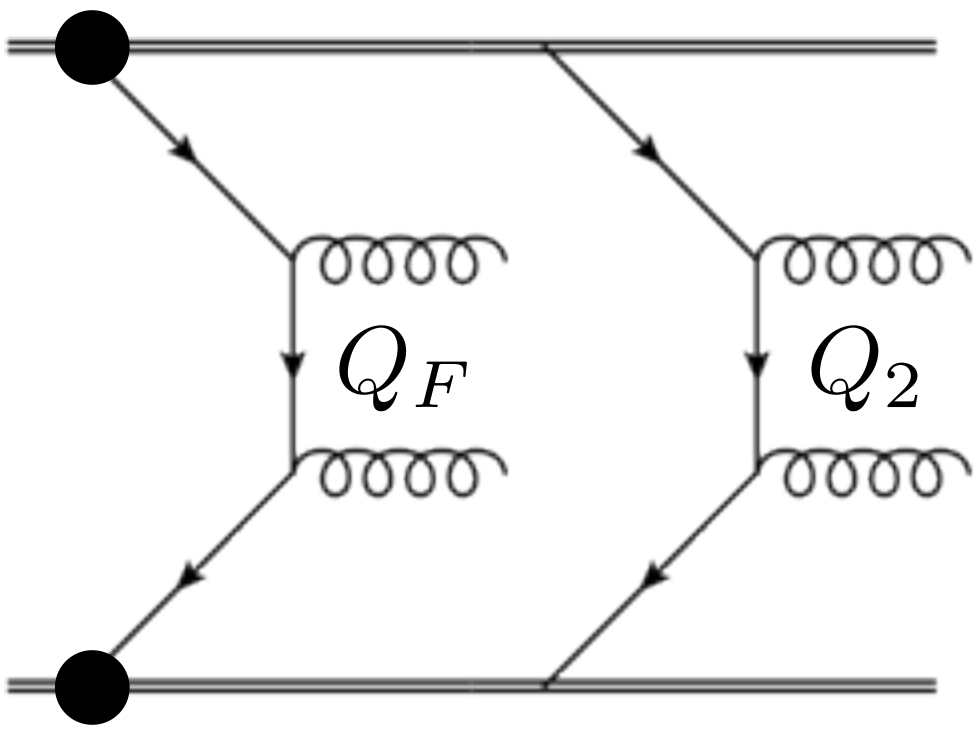

Multiple Parton Interactions go beyond existing theorems

→ perturbative short-distance physics in Underlying Event

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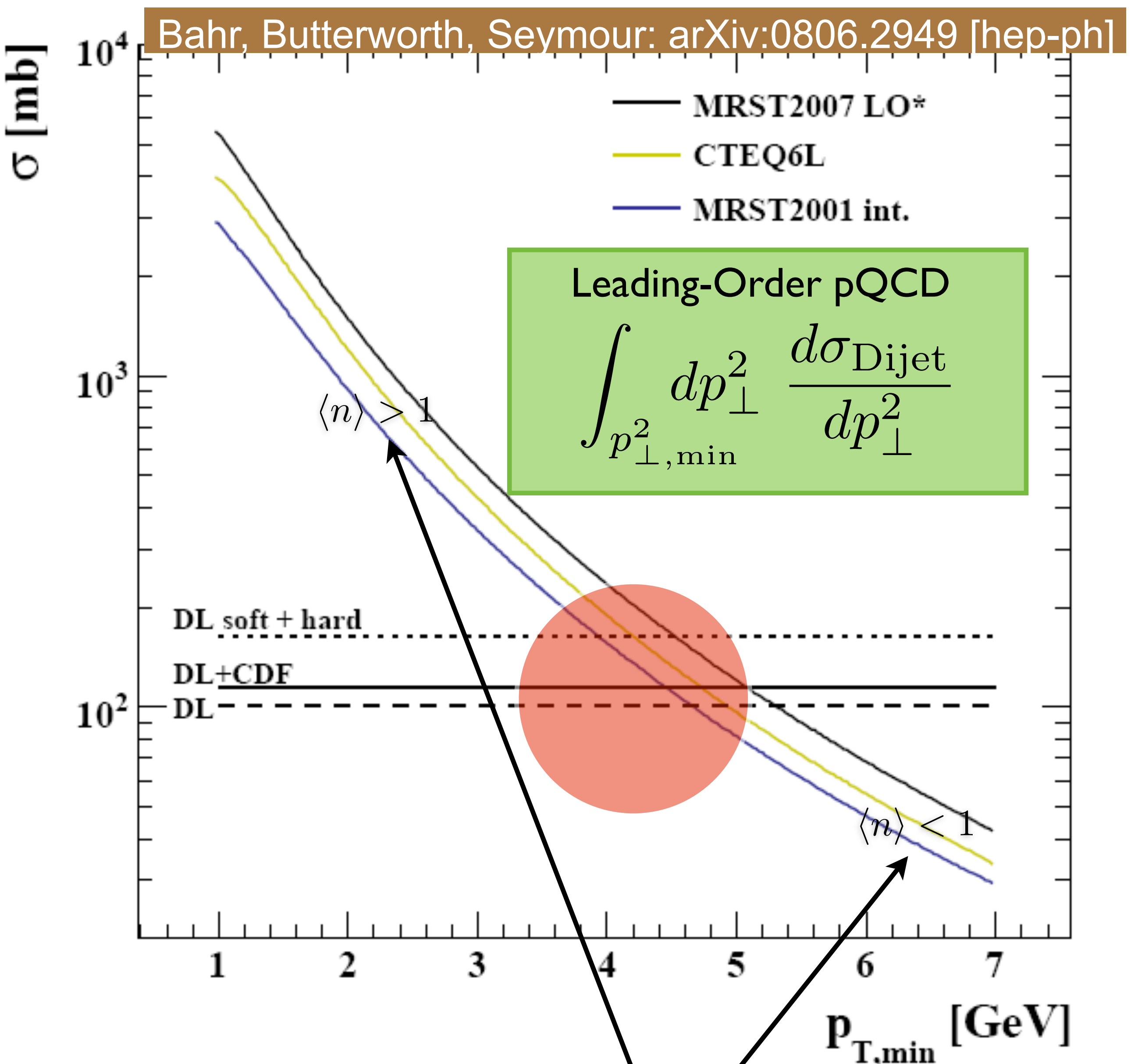
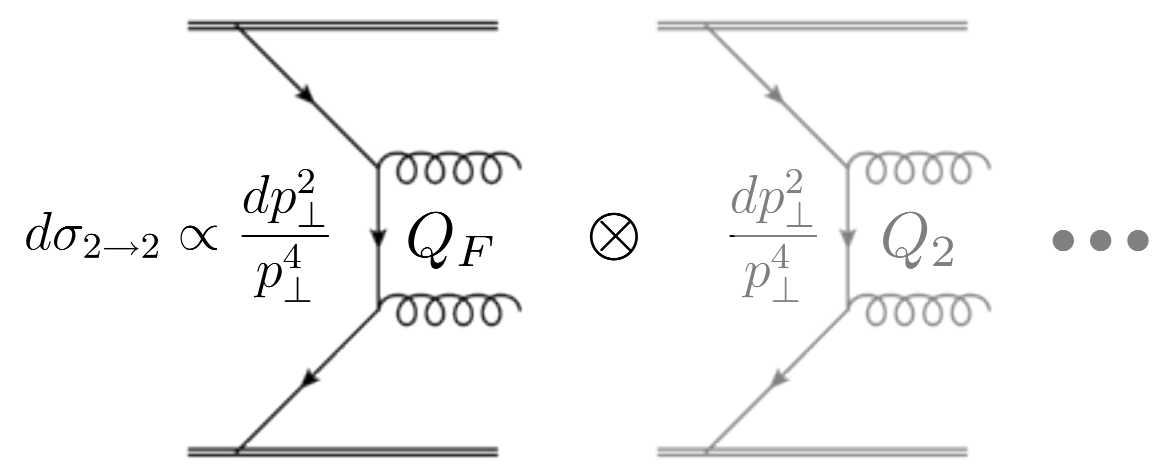
Multiple Parton Interactions

= Allow several parton-parton interactions per hadron-hadron collision. Requires extended factorization ansatz.



Earliest MC model ("old" PYTHIA 6 model)
Sjöstrand, van Zijl PRD36 (1987) 2019



Lesson from bremsstrahlung in pQCD: divergences → fixed-order breaks down

Perturbation theory still ok, with resummation (unitarity)

→ Resum dijets?
Yes → MPI!

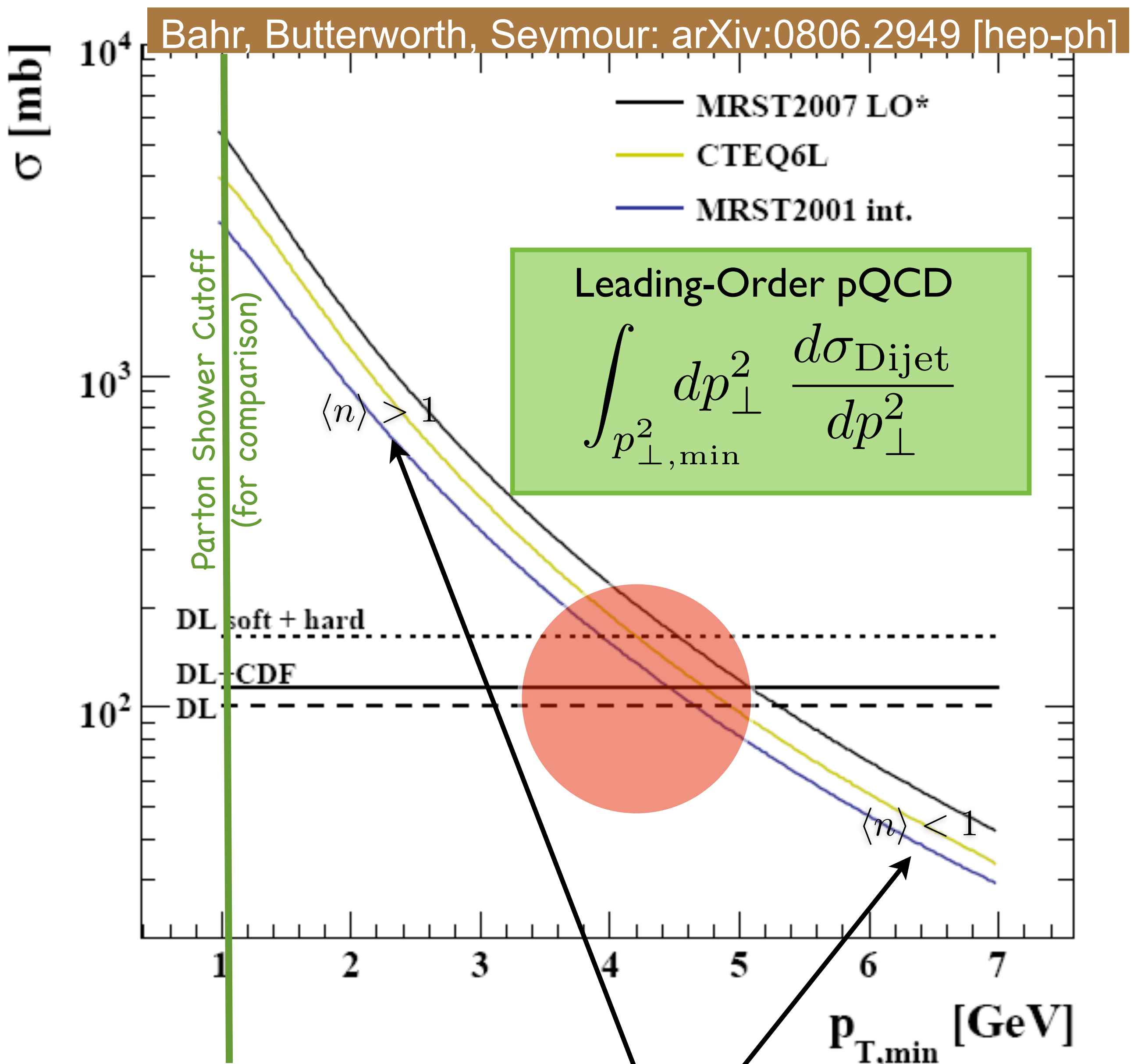
$$\sigma_{2\to 2}(p_{\perp\min}) = \langle n \rangle (p_{\perp\min})\ \sigma_{\rm tot}$$

Parton-Parton Cross Section

Hadron-Hadron Cross Section

Multiple Parton Interactions

= Allow several parton-parton interactions per hadron-hadron collision. Requires extended factorization ansatz.



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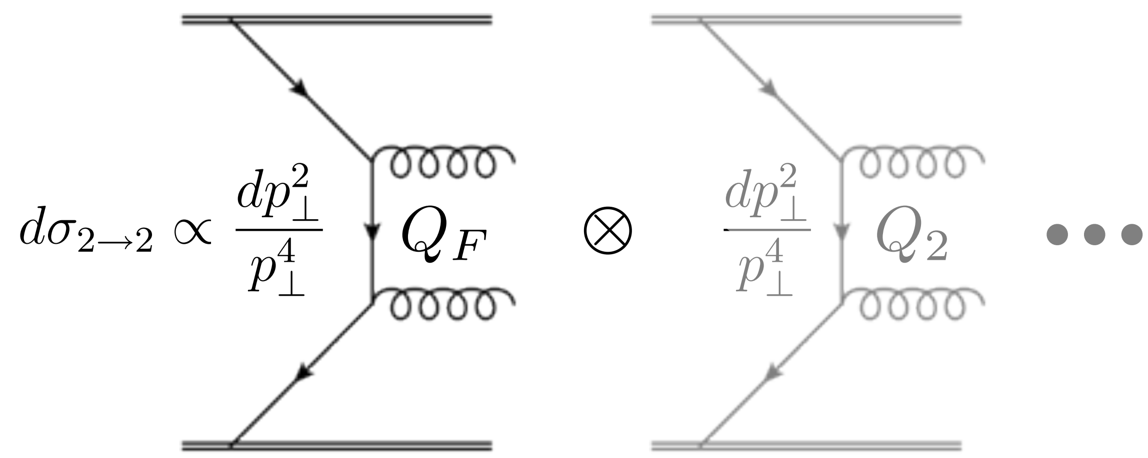

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Yes → MPI!

How many?

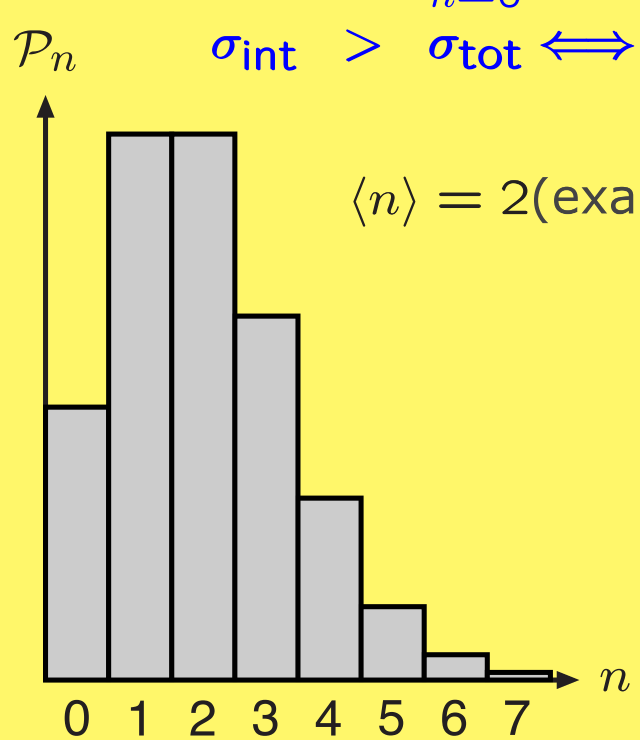
Naively $\langle n_{2\to2}(p_{\perp\min})\rangle = \frac{\sigma_{2\to2}(p_{\perp\min})}{\sigma_{\rm tot}}$

Interactions independent (naive factorization) → Poisson

$$\sigma_{\rm tot} = \sum_{n=0}^{\infty} \sigma_n$$

$$\sigma_{\rm int} = \sum_{n=0}^{\infty} n\,\sigma_n$$

$$\sigma_{\rm int} > \sigma_{\rm tot} \Longleftrightarrow \langle n \rangle > 1$$



$$\mathcal{P}_n = \frac{\langle n \rangle^n}{n!}\, e^{-\langle n \rangle}$$

Real Life

Color screening: $\sigma_{2\to2}\to0$ for $p_\perp\to0$

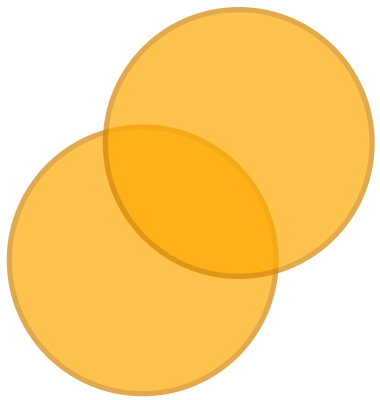
Momentum conservation suppresses high-n tail

Impact-parameter dependence

\+ physical correlations

→ not simple product

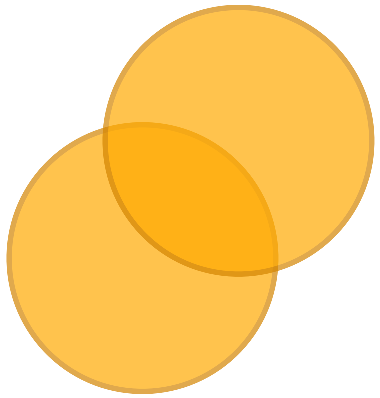
Impact Parameter



1. **Simple Geometry** (in impact-parameter plane)

Simplest idea: smear PDFs across a uniform disk of size πr_p^2
→ simple geometric overlap factor ≤ 1 in dijet cross section
Some collisions have the full overlap, others only partial
→ Poisson distribution with different mean <n> at each b

Impact Parameter



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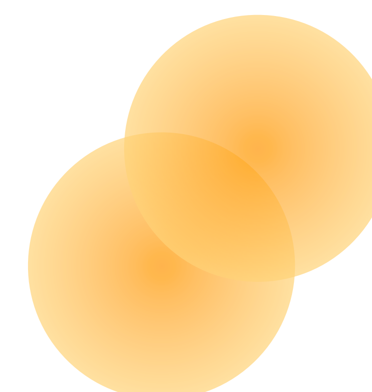
2. More realistic **Proton b-shape**



Smear PDFs across a non-uniform disk
MC models use Gaussians or **more**/less peaked
Overlap factor = convolution of two such distributions

→ Poisson distribution with different mean <n> at each b
"Lumpy Peaks" → large matter overlap enhancements, higher <n>

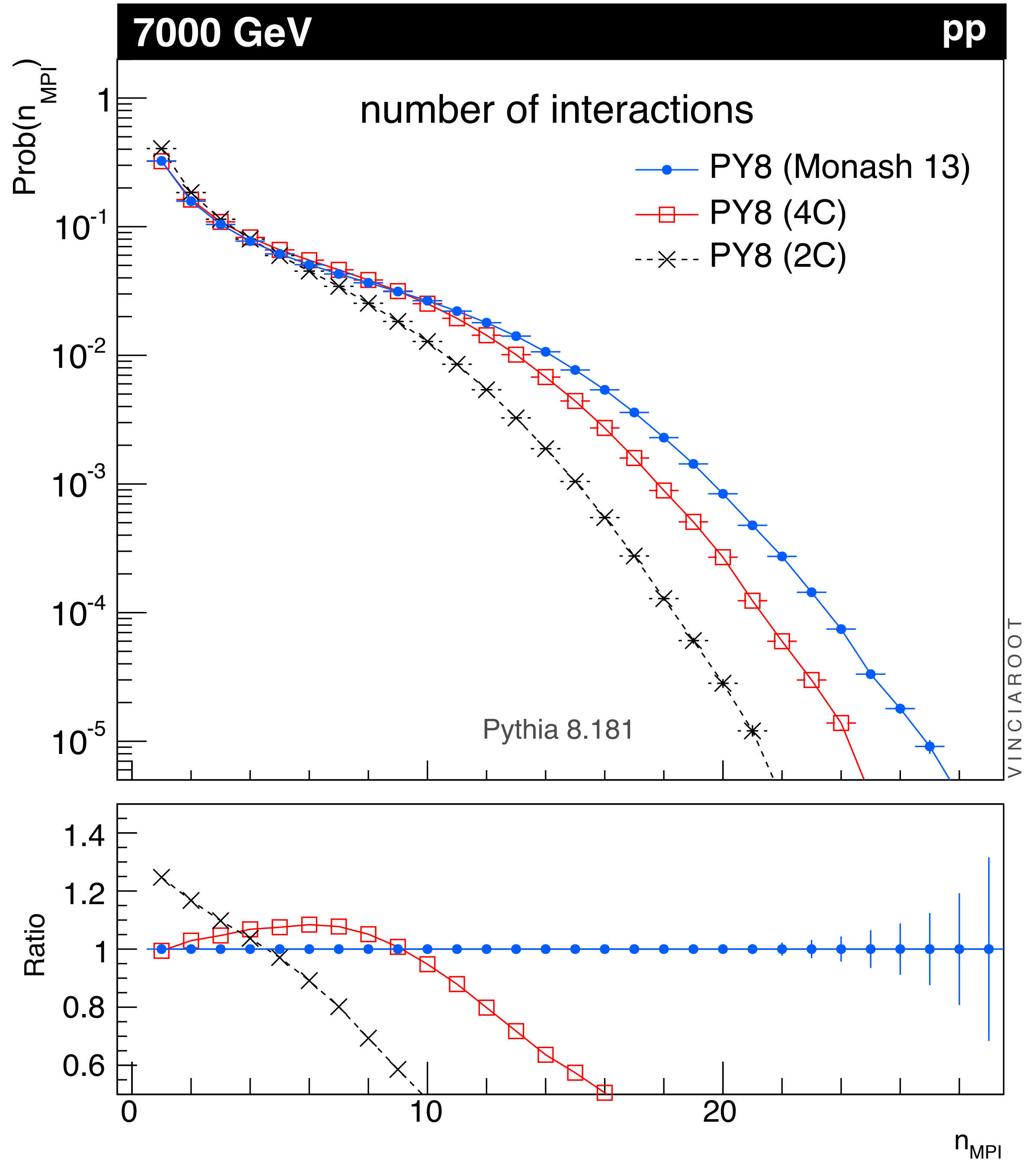
Note: this is an *effective* description. Not the actual proton mass density.
E.g., peak in overlap function (≫1) can represent unlikely configurations with huge overlap enhancement. Typically use total σ_{inel} as normalization.

Number of MPI *

Minimum-Bias pp collisions at 7 TeV

Averaged over all pp impact parameters

(Really: averaged over all pp overlap enhancement factors)

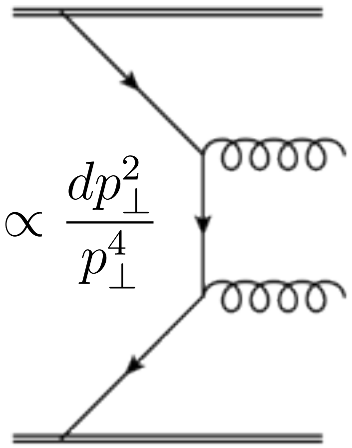


*note: can be arbitrarily soft

Caveats of MPI-Based Models

Main applications: Central Jets/EWK/top/ Higgs/New Physics

$$d\sigma_{2\to 2} \propto \frac{dp_\perp^2}{p_\perp^4}$$



$\otimes$ PDFs

High Q^2 and finite x

See also Connecting hard to soft: KMR, EPJ C71 (2011) 1617 + PYTHIA "Perugia Tunes": PS, PRD82 (2010) 074018 + arXiv:1308.2813

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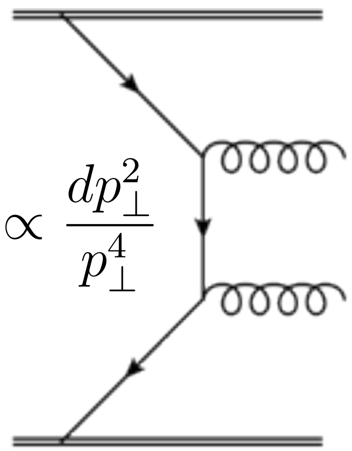
$$d\sigma_{2\to 2} \propto \frac{dp_\perp^2}{p_\perp^4}$$  $\otimes$ PDFs

High Q^2 and finite x

Extrapolation to soft scales delicate.
Impressive successes with MPI-based models but still far from a solved problem

- Form of PDFs at small x and Q^2
- Form and E_{cm} dependence of p_{T0} regulator
- Modeling of the diffractive component
- Proton transverse mass distribution
- Colour Reconnections, Collective Effects

Saturation

See talk on UE by W. Waalewijn

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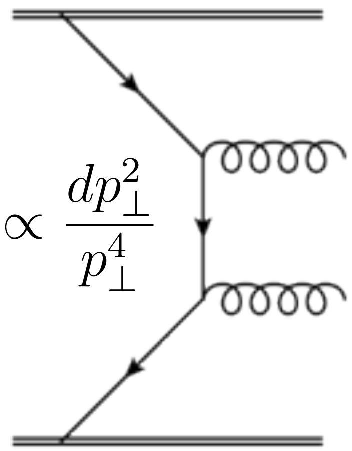

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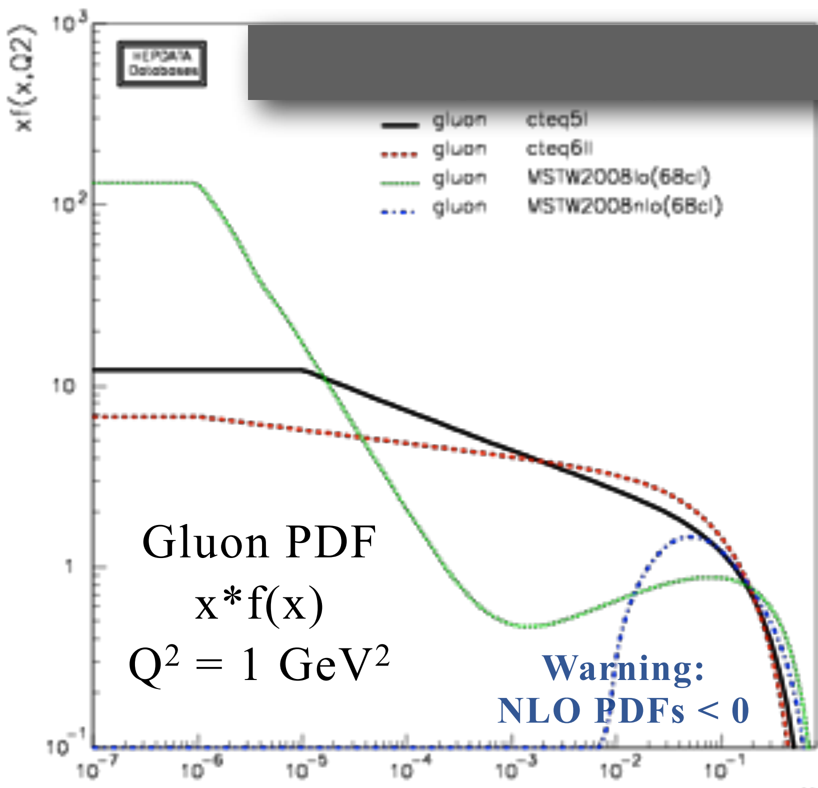
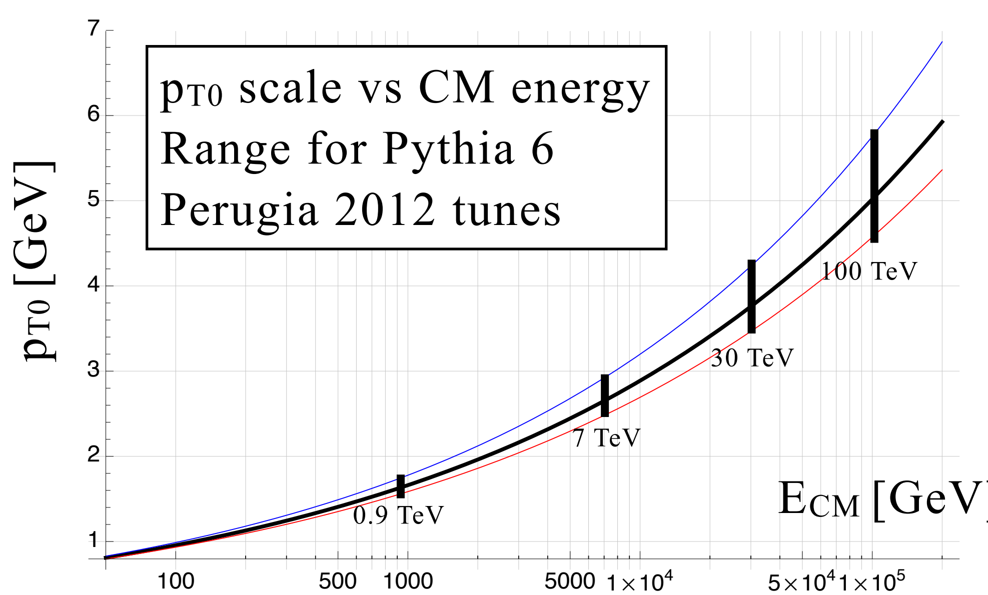
Poor Man's Saturation





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1: A Simple Model

The minimal model incorporating single-parton factorization, perturbative unitarity, and energy-and-momentum conservation

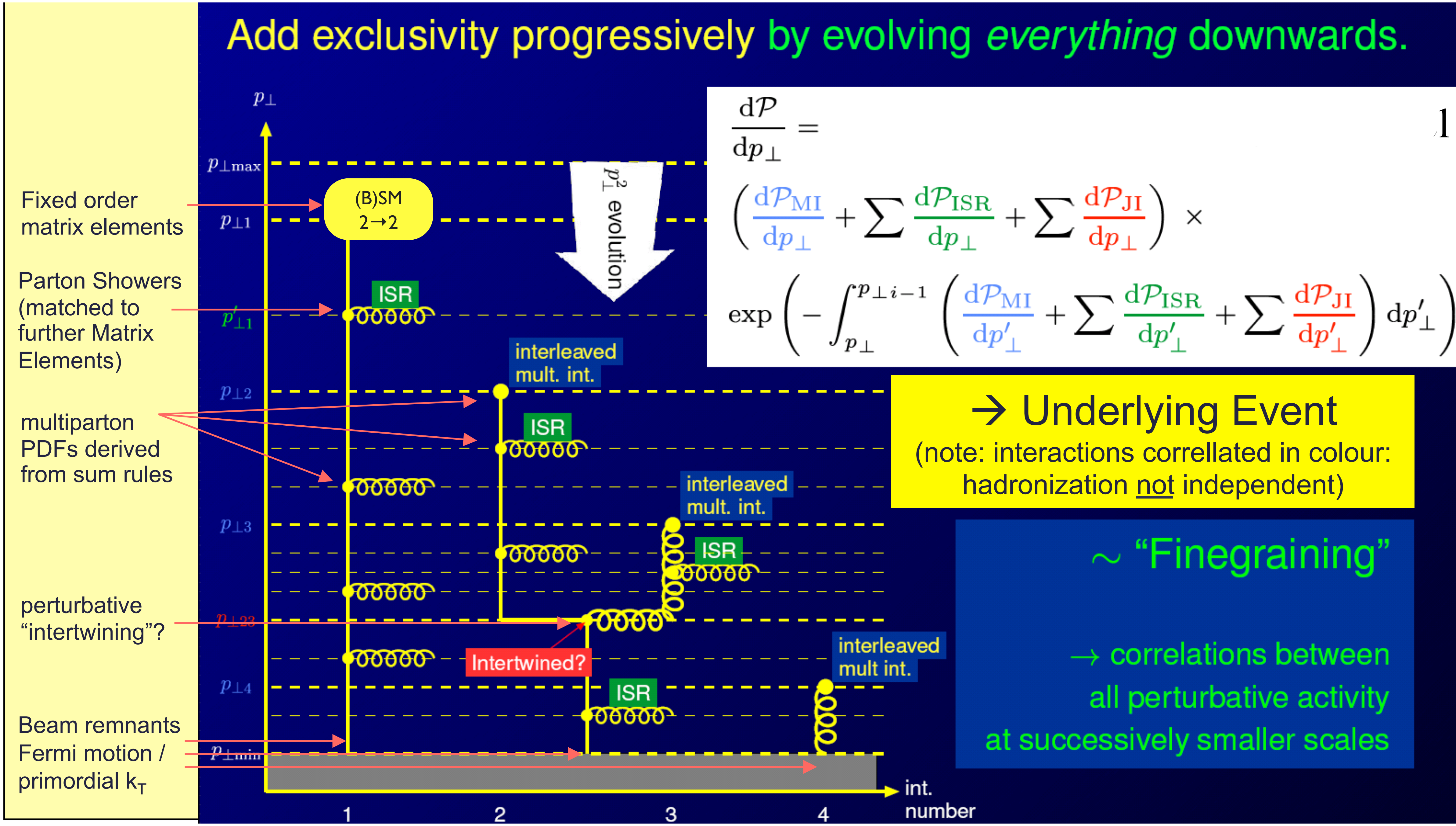
$$\sigma_{2\to 2}(p_{\perp\min}) = \langle n \rangle (p_{\perp\min})\ \sigma_{\mathrm{tot}}$$

Parton-Parton Cross Section   Hadron-Hadron Cross Section

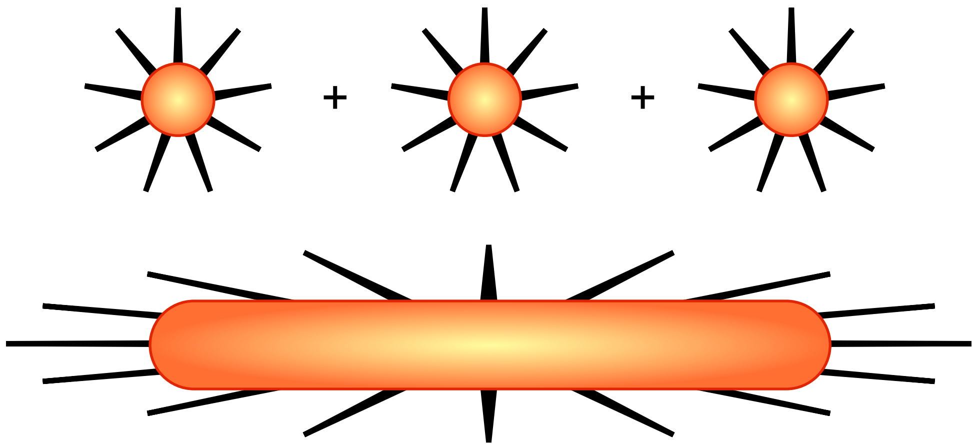
1. Choose $p_{T\min}$ cutoff
   = main tuning parameter
2. Interpret $<n>(p_{T\min})$ as mean of Poisson distribution
   Equivalent to assuming all parton-parton interactions equivalent and independent ~ each take an instantaneous "snapshot" of the proton
3. Generate n parton-parton interactions (pQCD 2→2)
   Veto if total beam momentum exceeded → overall (E,p) cons
4. Add impact-parameter dependence → $<n> = <n>(b)$
   Assume factorization of transverse and longitudinal d.o.f., → PDFs : $f(x,b) = f(x)g(b)$ (Ordinary CTEQ, MSTW, NNPDF, …)
   b distribution ∝ EM form factor → **JIMMY model** Butterworth, Forshaw, Seymour Z.Phys. C72 (1996) 637
   Constant of proportionality = second main tuning parameter
5. Add separate class of "soft" (zero-p_T) interactions representing interactions with $p_T < p_{T\min}$ and require $\sigma_{\mathrm{soft}} + \sigma_{\mathrm{hard}} = \sigma_{\mathrm{tot}}$
   → **Herwig++ model** Bähr et al, arXiv:0905.4671

2: Interleaved Evolution

Sjöstrand, P.S., JHEP 0403 (2004) 053; EPJ C39 (2005) 129

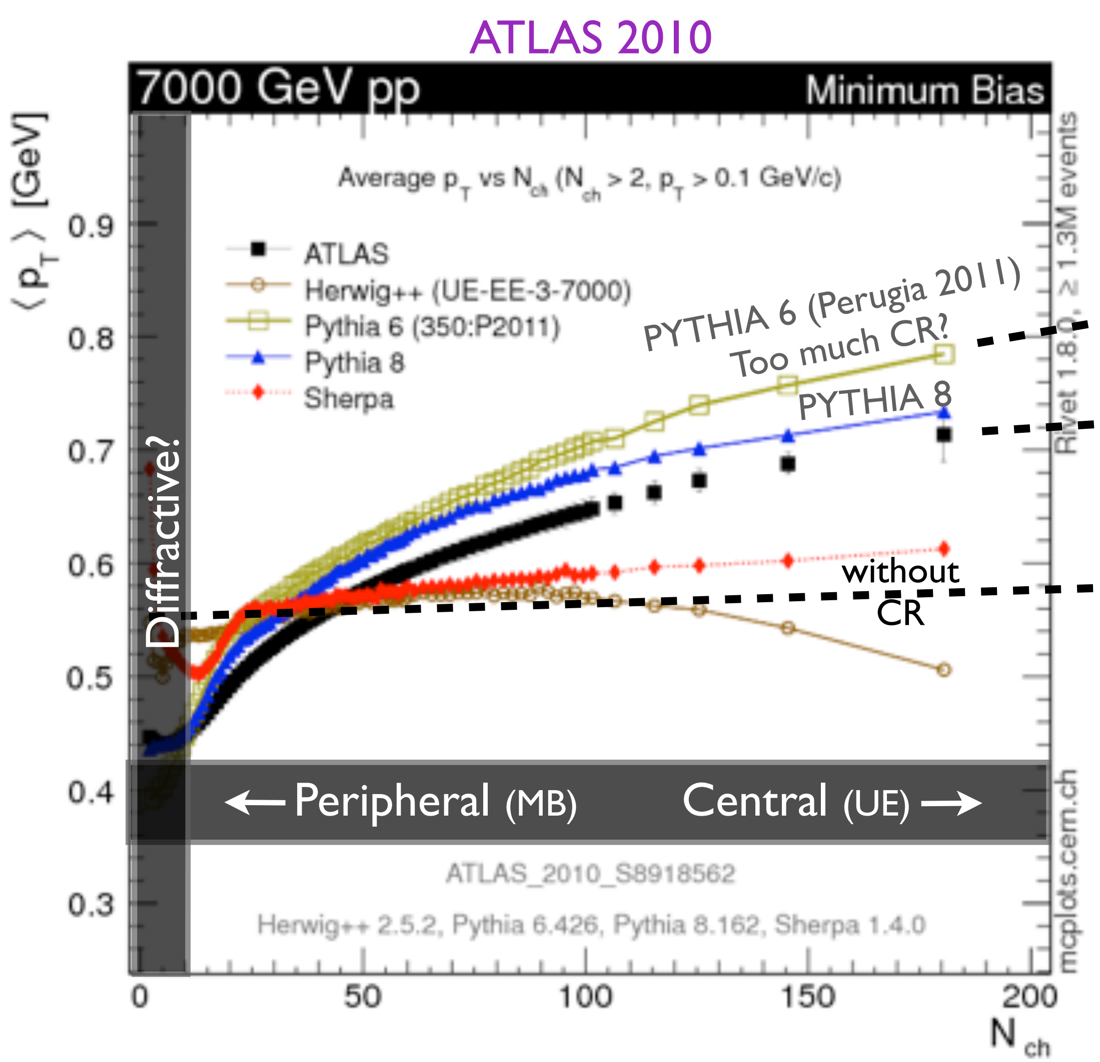


<p_T> vs N_{ch}



Independent Particle Production:
→ **averages stay the same**

Correlations / Collective effects:
→ **average rises**



Extrapolation to high multiplicity ~ UE

Average particles slightly too hard
→ Too much energy, or energy distributed on too few particles

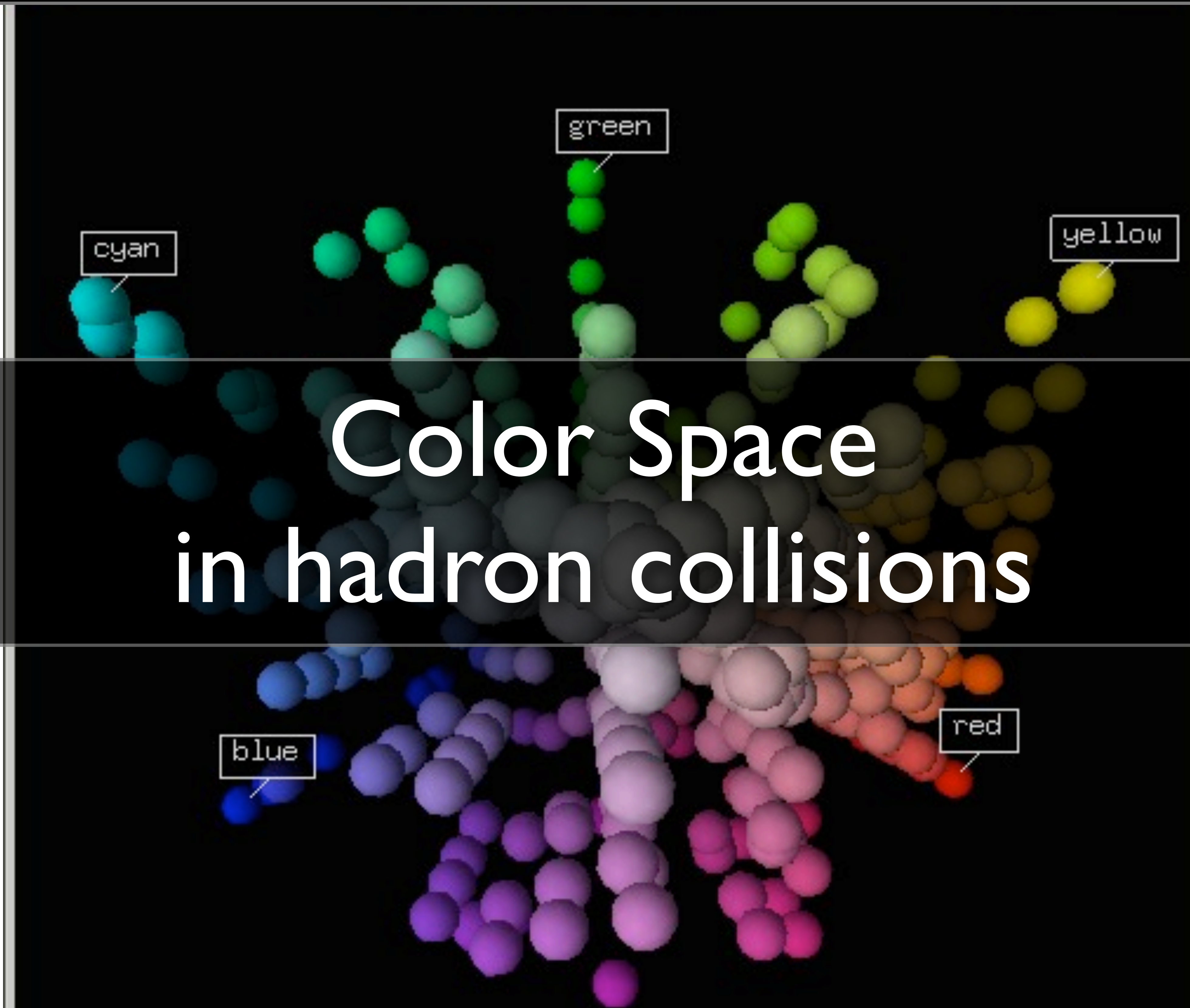
~ OK?

Average particles slightly too soft
→ Too little energy, or energy distributed on too many particles

Evolution of other distributions with N_{ch} also interesting: e.g., <p_T>(N_{ch}) for identified particles, strangeness & baryon ratios, 2P correlations, …

Plots from mcplots.cern.ch

Color Space in hadron collisions

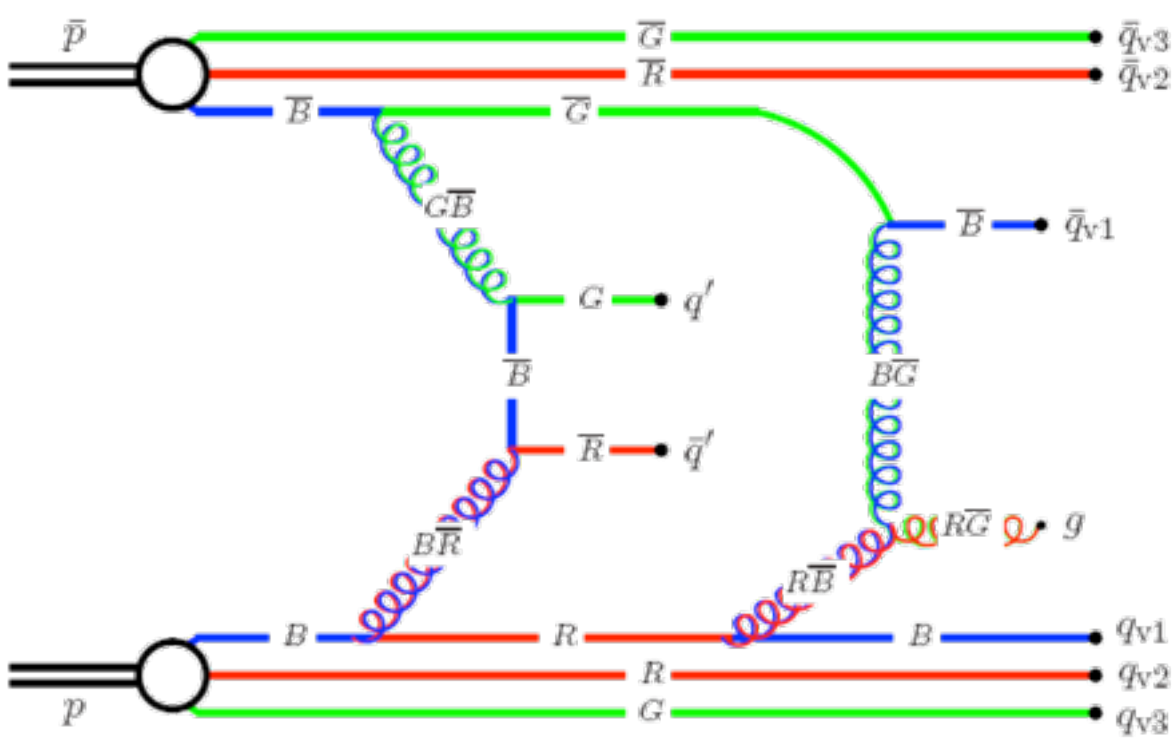


Color Correlations

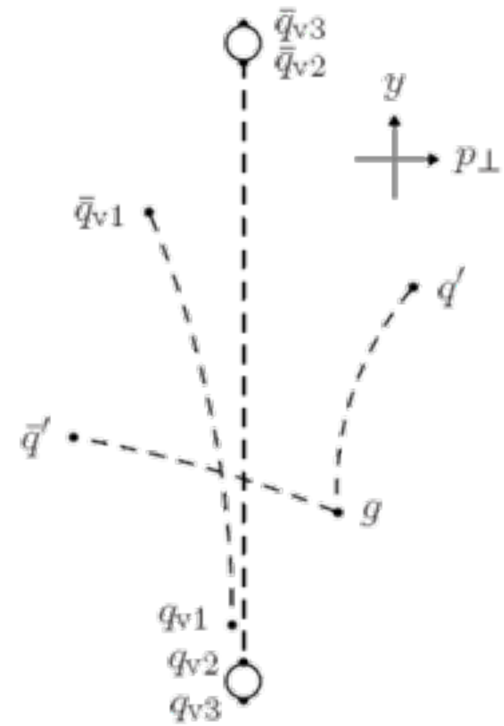
Each MPI (or cut Pomeron) exchanges color between the beams

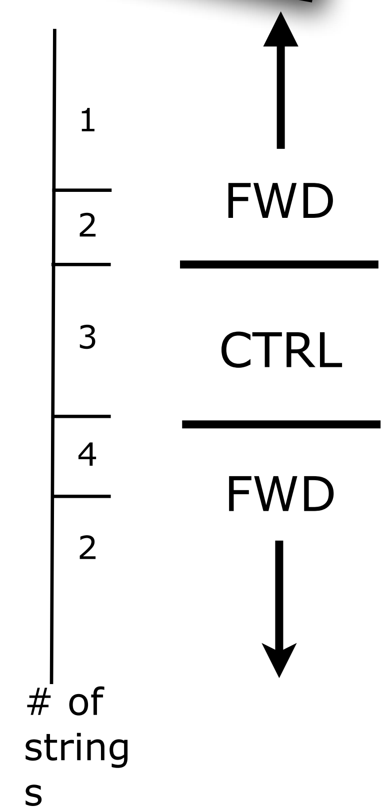
- ▶ The colour flow determines the hadronizing string topology
  - Each MPI, even when soft, is a color spark
  - Final distributions crucially depend on color space

Different models make different ansätze



Sjöstrand & PS, JHEP 03(2004)053





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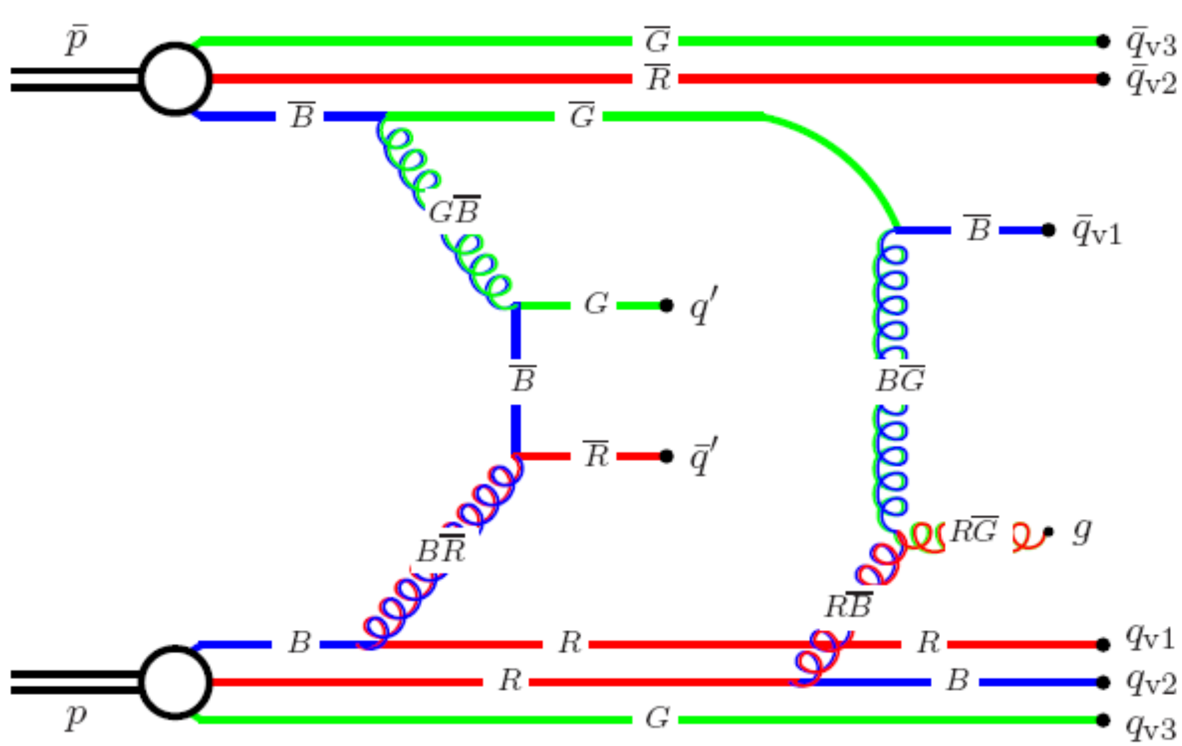

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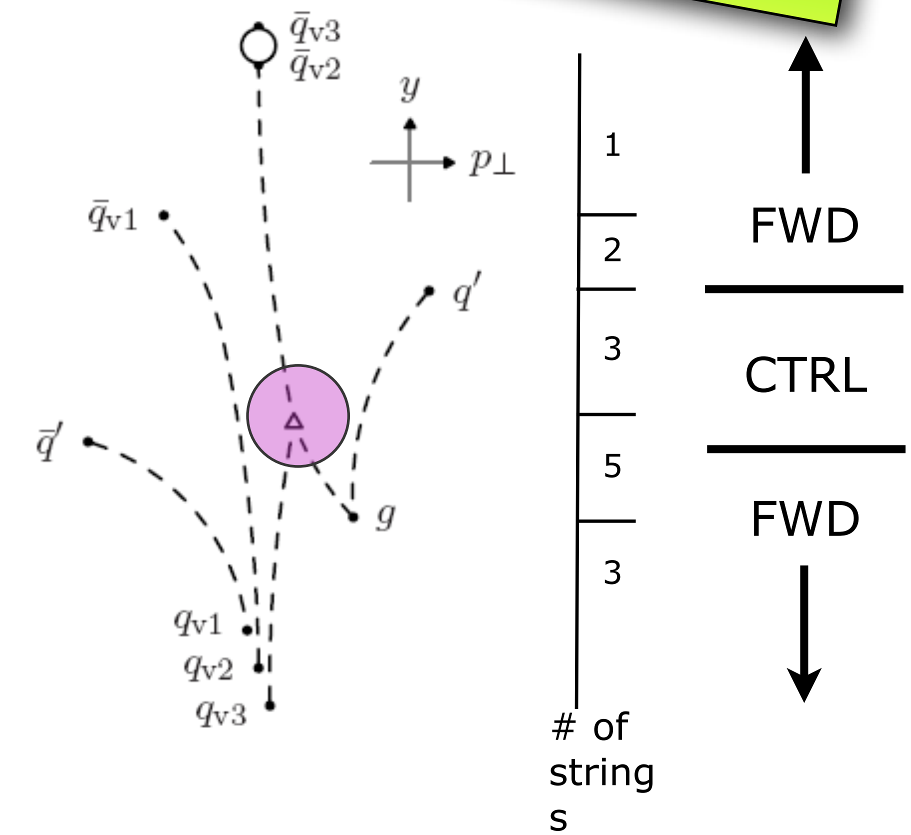
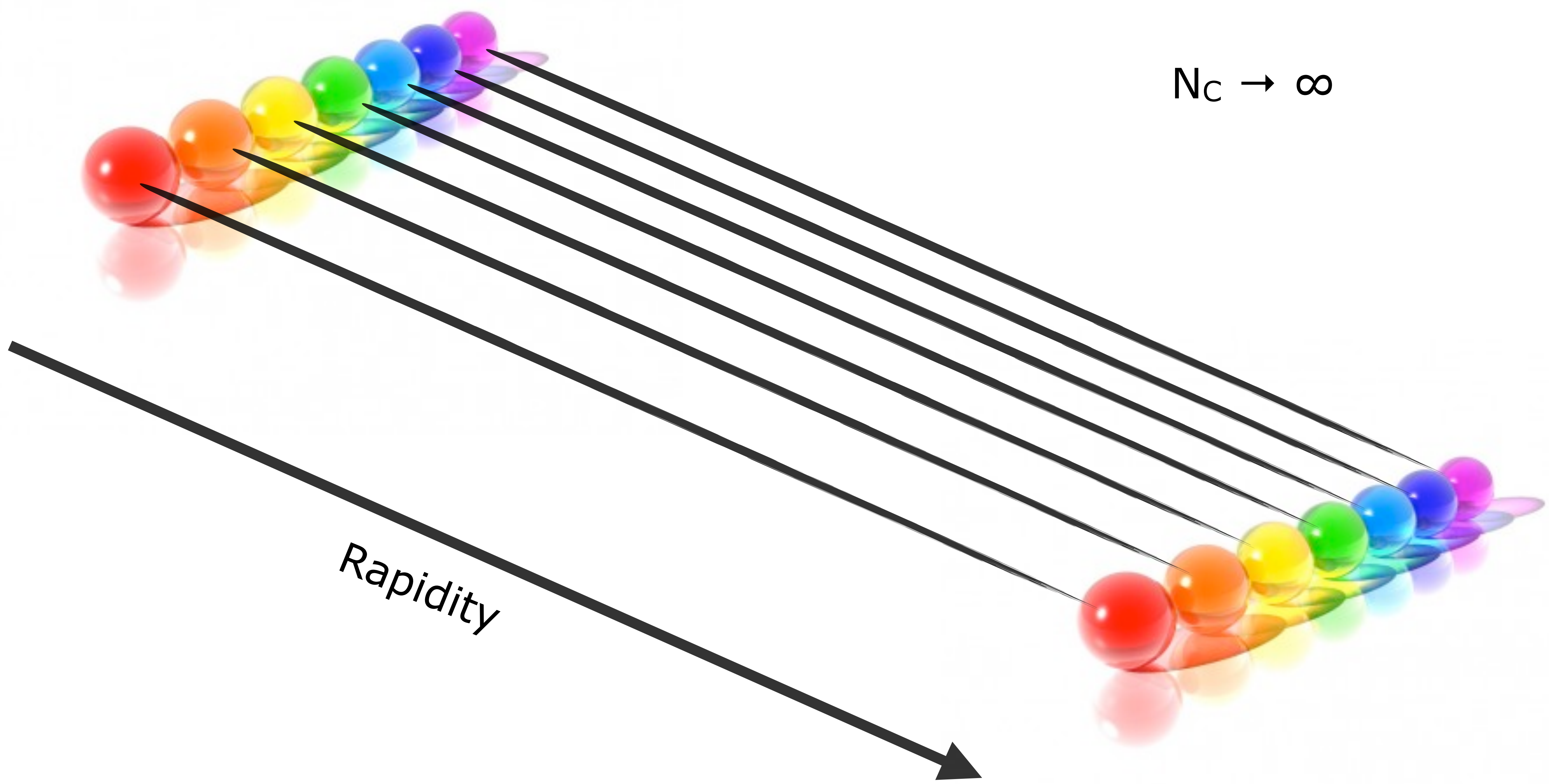


Color Connections

Better theory models needed

$N_C \to \infty$



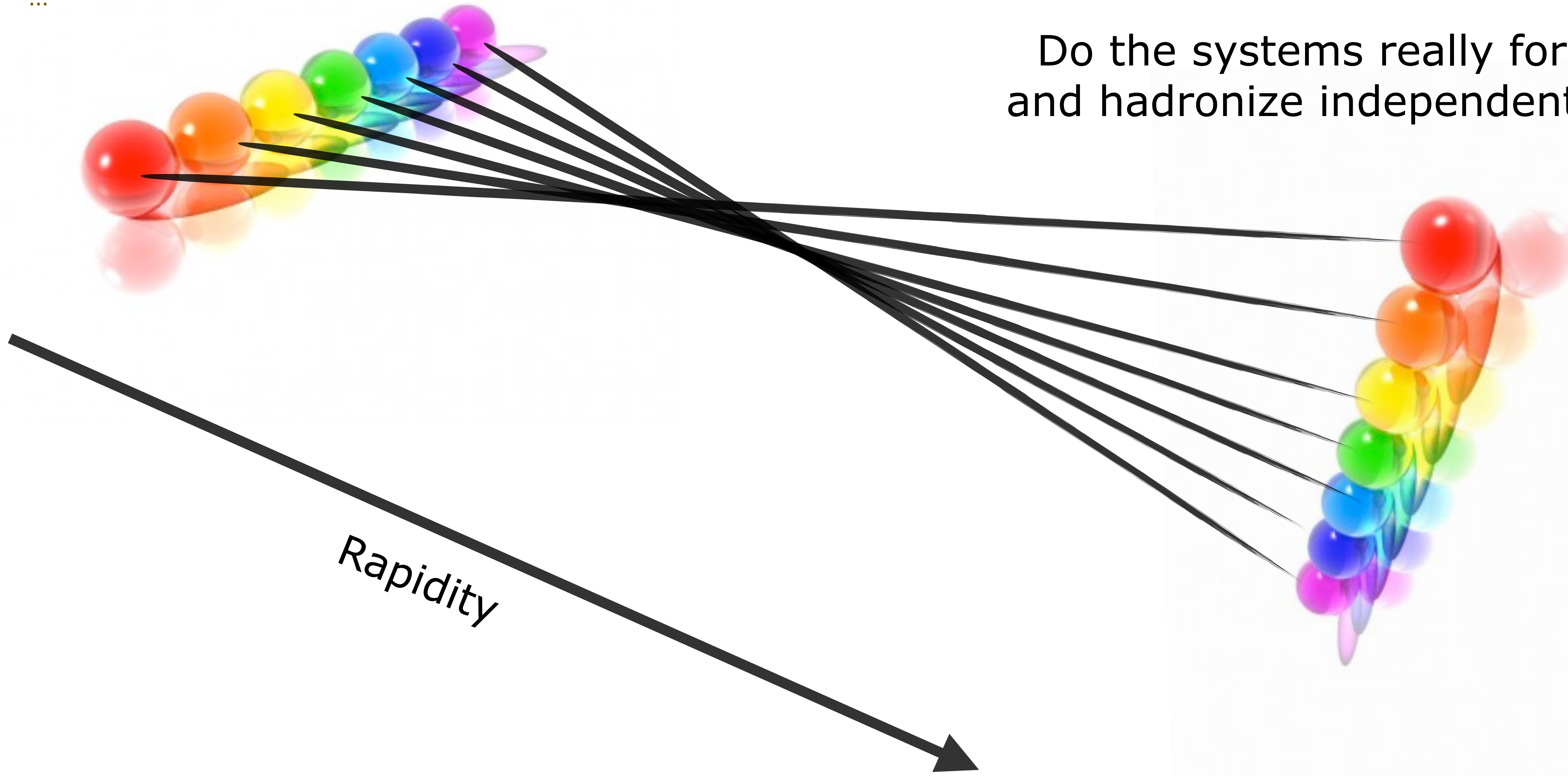
Multiplicity $\propto N_{MPI}$

Color Reconnections?

E.g.,
Generalized Area Law (Rathsman: Phys. Lett. B452 (1999) 364)
Color Annealing (P.S., Wicke: Eur. Phys. J. C52 (2007) 133)
…

Better theory models needed

Do the systems really form and hadronize independently?

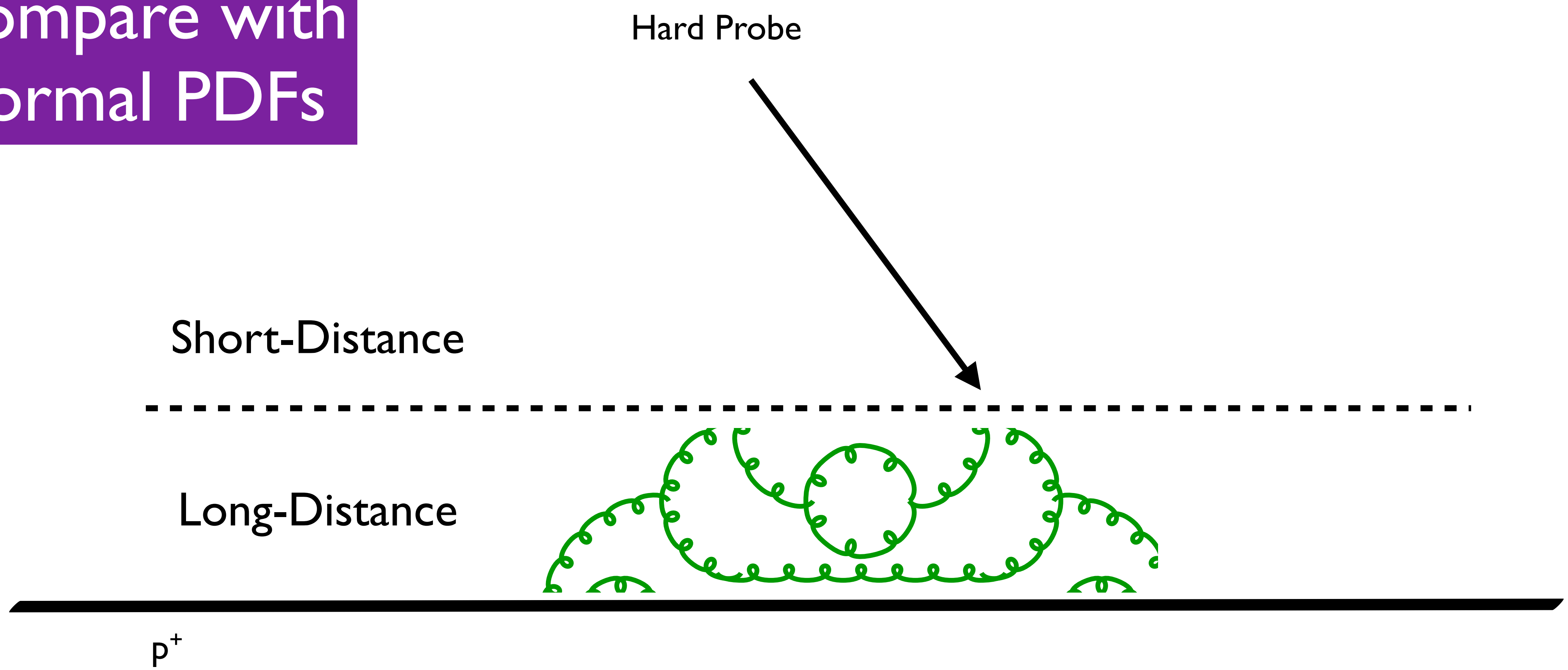


Multiplicity $\not\propto$ N_{MPI} ($<$)

(+ Diffraction)

"Intuitive picture"

Compare with normal PDFs



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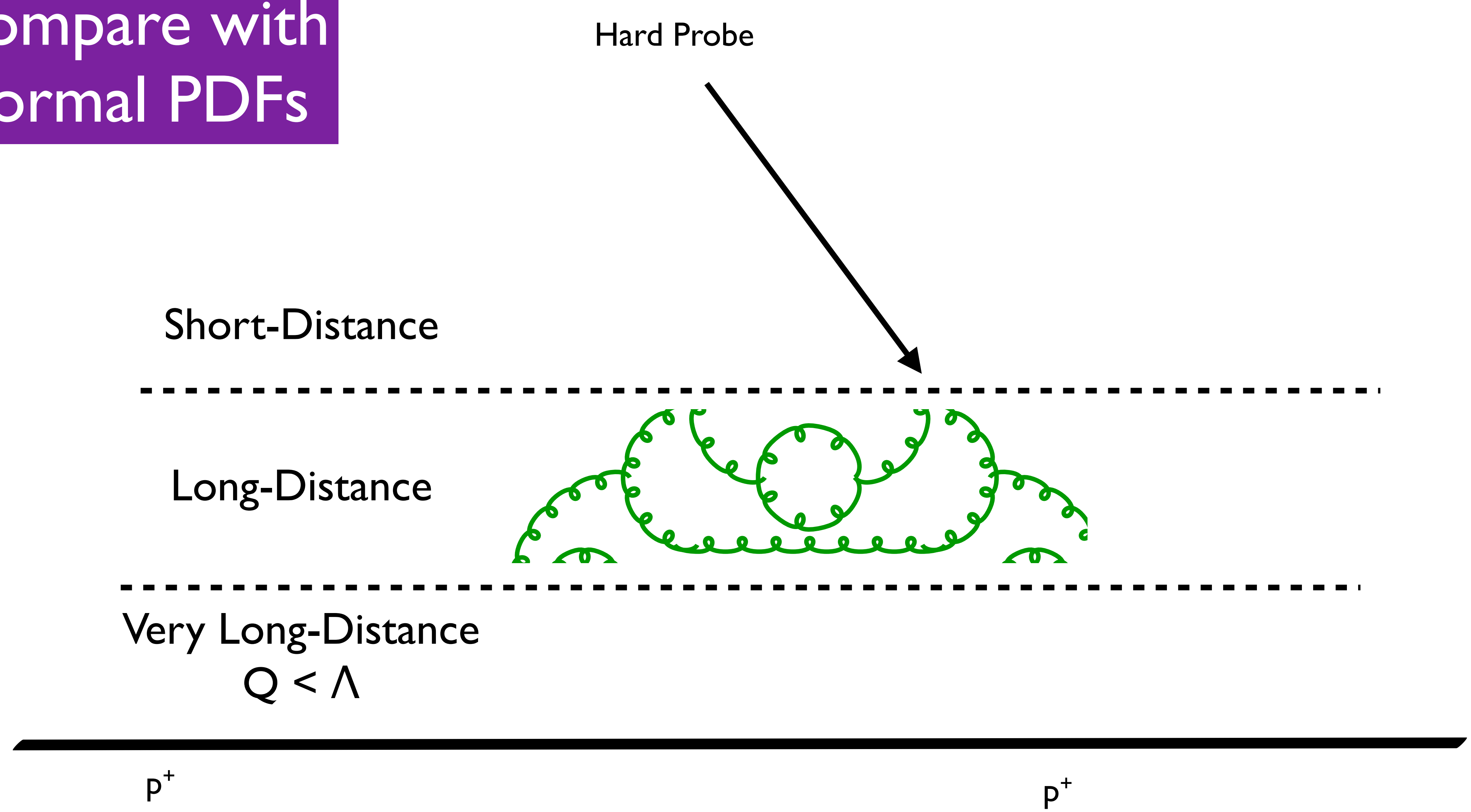

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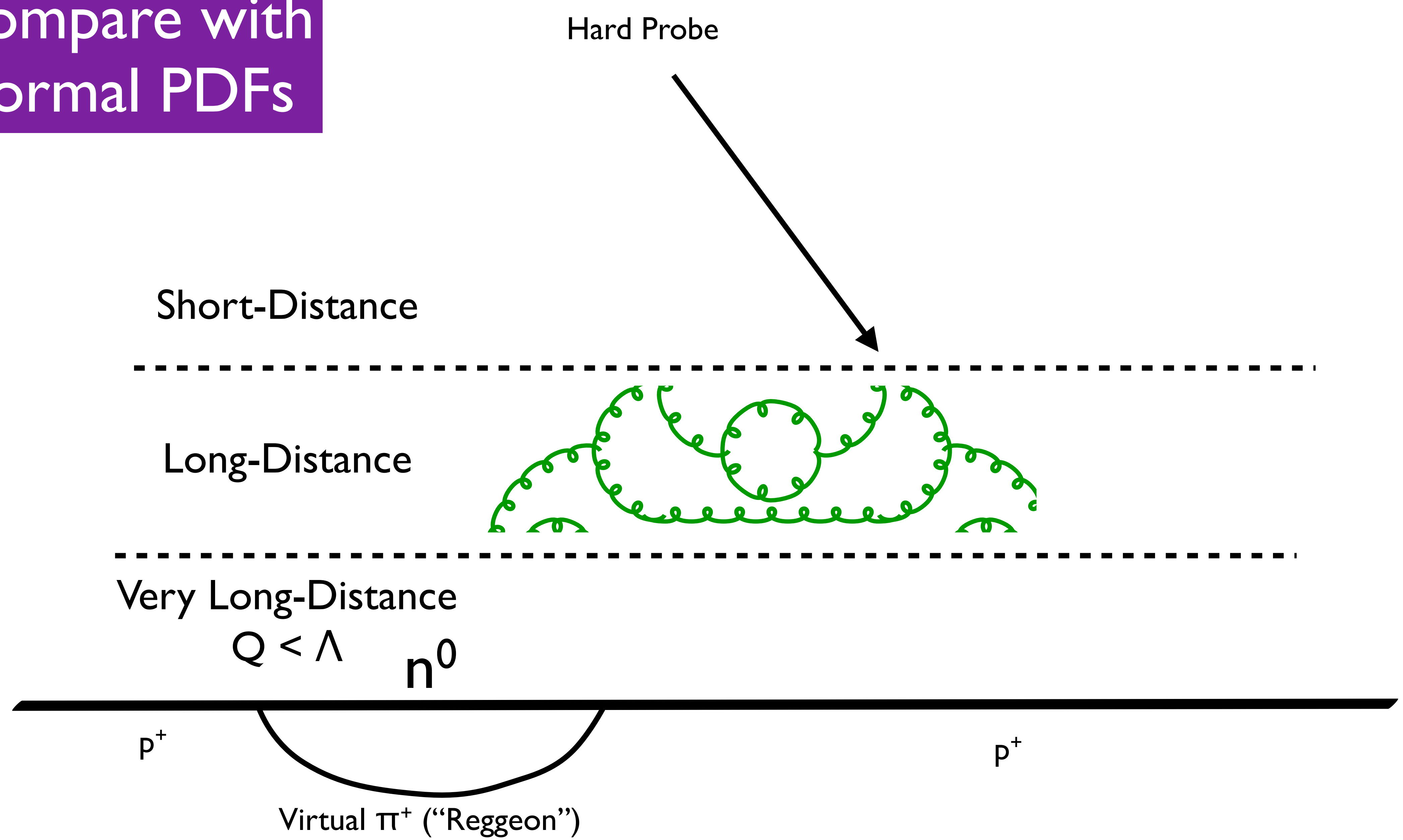

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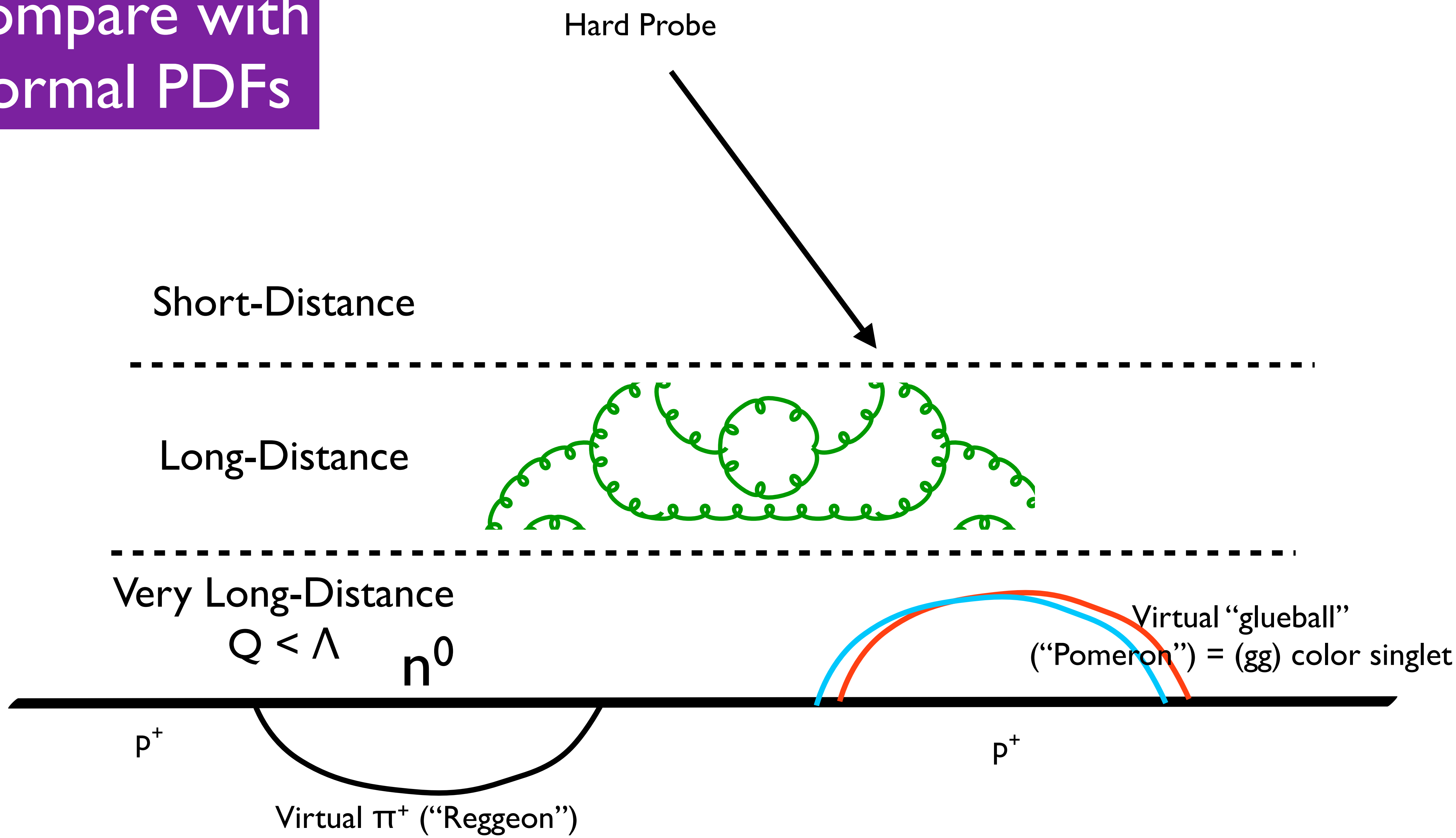

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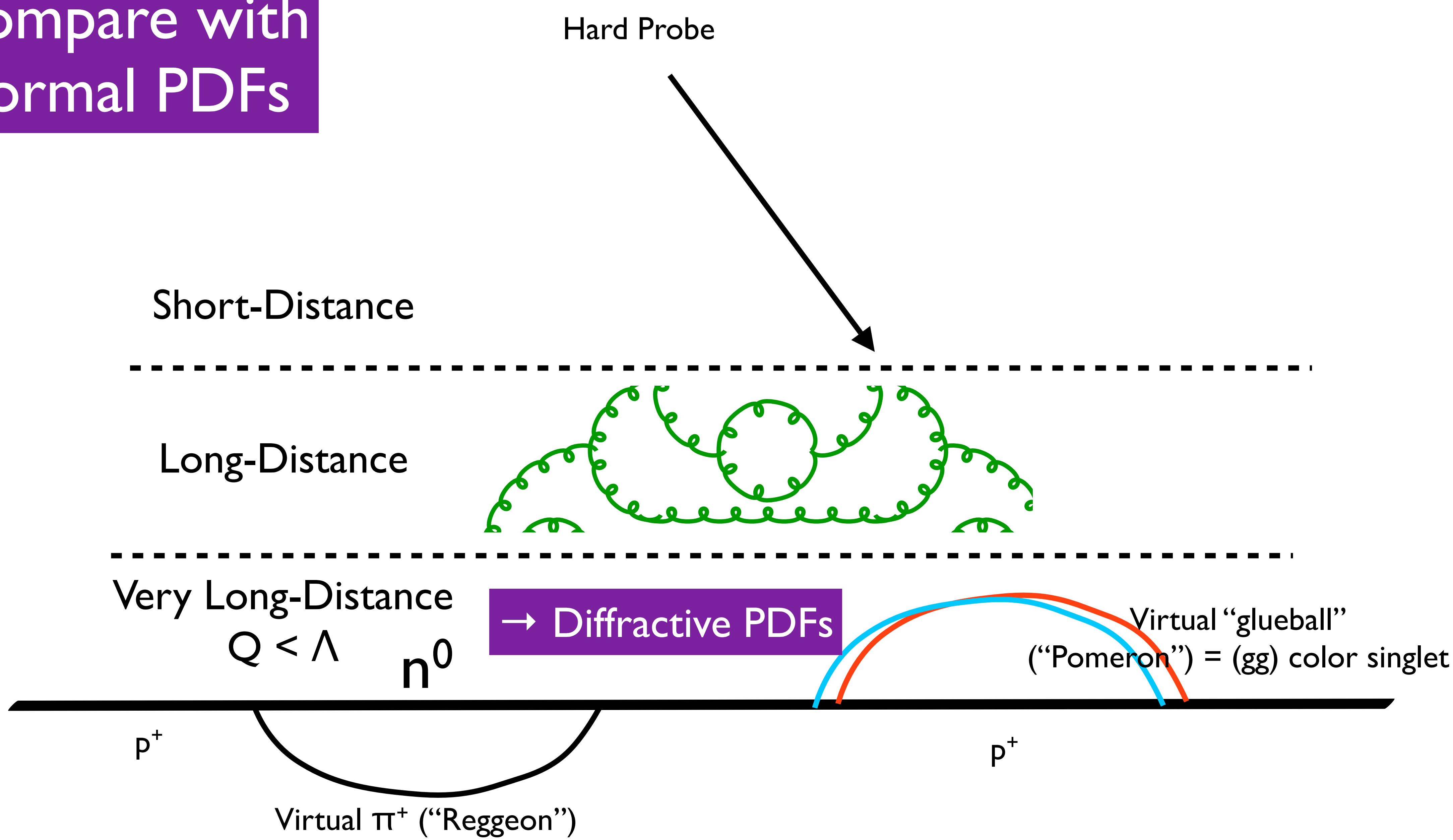

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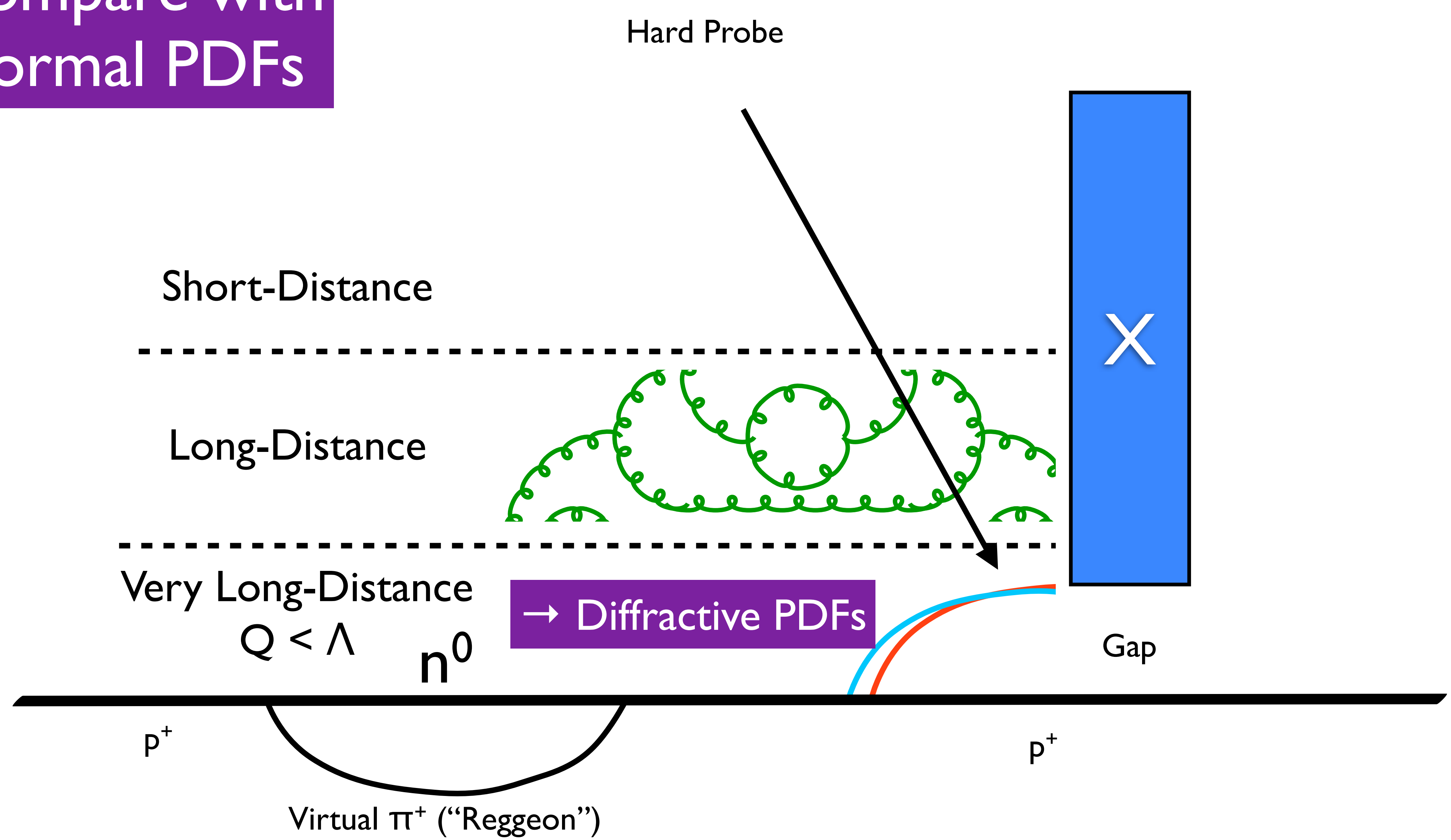

Tuning

means different things to different people